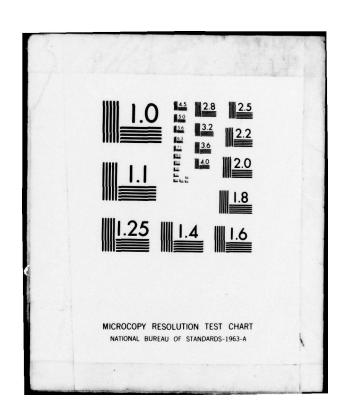
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A METHOD OF CHARACTERISTICS APPROACH TO THE PROBLEM OF SUPERSONIC FLOW PAST OSCILLATING CASCADES WITH FINITE BLADE THICKNESS

K. Vogeler

October 1978

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1. INTRODUCTION

The demand for increasing performance and efficiency of turbines and compressors in jet engines forces the industry to advanced designs. This means that it is no longer acceptable to use blades in turbomachines which are loaded considerably below their mechanical limits. There are two major ways to improve the engine thrust-to-weight and thrust-to-volume characteristics:

- Reduced weight and size
- Increased massflow, temperature and pressure.

It is obvious that only a compromise can be successful, as advances in the aerothermodynamics oppose those of the structure.

These trends require the use of slender and thin blades which are increasingly susceptible to flutter and vibration problems. The most important ones are:

- 1. Supersonic unstalled flutter
- 2. Forced response
- 3. Subsonic stall flutter
- 4. Choke flutter
- 5. Supersonic stall flutter

While the latter three are quite difficult to describe in a fluid mechanical model, the unstalled supersonic blade flutter is amenable to analysis with reasonable effort. What makes it even more interesting and at the same time highly important is the possibility of its occurrence at the design condition of the engine. Especially, modern fans with large diameters operate with the outer part of their blades in the transonic flow region (1. < M < 1.5). Their flutter susceptibility therefore makes the analysis of supersonic unstalled flutter increasingly important. The problem has been attacked not only in the United States but in all major industrialized countries, which shows that it is

a general one to be considered in the design of modern high-performance turbo-

During the past decade, several methods have been developed to predict supersonic blade flutter. All were based on the idealization of the actual flow by the planar flow through a staggered cascade of oscillating blades. However, the two-dimensional flow and the flat plate assumptions impose severe simplifications whose range of validity needs to be better understood.

Since the incorporation of three-dimensional flow effects is rather difficult, it seems logical to first explore the effect of blade thickness and shape on supersonic blade flutter while retaining the cascade concept. To this end, the nonlinear transonic small perturbation equation is adopted in this report as the governing equation. The use of this equation rather than the full potential equation or the Euler equations is suggested by Teipel's success to analyze the thickness effect of a single oscillating airfoil in low supersonic flow. Hence, the present work is an extension of Teipel's method to oscillating supersonic cascades for the purpose of determining the influence of steady non-uniform flow effects due to blade thickness, shape, camber or angle of attack on the oscillatory pressure distributions, forces and moments. The ultimate goal of this study is to replace the transonic small disturbance equation by the Euler equations so that the range of applicability of this simpler equation can be ascertained.

2. THE BASIC EQUATIONS

2.1 The Nonlinear Transonic Equation

Landahl presents in (1) the differential euqation which is valid for the transonic flow region

$$\left[M^{2}-1+\frac{U\left(\kappa+1\right)}{a^{2}}\frac{\partial\phi}{\partial X}\right]\frac{\partial^{2}\phi}{\partial X^{2}}-\frac{\partial^{2}\phi}{\partial Y^{2}}+\frac{1}{a^{2}}+\frac{\partial^{2}\phi}{\partial T^{2}}+\frac{2M}{a}\frac{\partial^{2}\phi}{\partial X\partial T}=0$$
(1)

with U - Free stream velocity

a - Local velocity of sound

M - Free stream Mach number

Φ - Velocity potential (X,Y,T)

K - Ratio of specific heats

Eq. (1) can be written non-dimensionally by using the terms

$$x = \frac{X}{c}$$
, $y = \frac{Y}{c}$, $t = \frac{TU}{c}$, $\phi = \frac{\phi}{cU}$,

where c is the chord.

Thus we get the new form

$$\left[M^{2}-1+M^{2}\left(\kappa+1\right)\frac{\partial\phi}{\partial\mathbf{x}}\right]\frac{\partial^{2}\phi}{\partial\mathbf{x}^{2}}-M^{2}\frac{\partial^{2}\phi}{\partial\mathbf{t}^{2}}+2M^{2}\frac{\partial^{2}\phi}{\partial\mathbf{x}\partial\mathbf{t}}=0$$
(2)

Following Teipel, who developed in (2) for single airfoils a method of characteristics using Eq. (1), in this work Eq. (2) is to be used. This has already successfully been done by Platzer, Chadwick and Strada (3,4,5,6,7) for a saingle oscillating airfoil and for oscillating cascades of wedges and thick blades with flat upper surfaces. Nevertheless, the basic steps, which lead to the solution of this problem shall be repeated in this report, to make it at the same time a summary of the work already done by the authors above.

It is an extension in so far as it introduces the thickness effect of the upper blade surface in an oscillating, staggered cascade and shows a method-of-characteristics-approach to the unsteady supersonic wake of not only a flat plate but also airfoils with small but finite thickness.

As we consider only small perturbations of the freestream flow values, the potential function ϕ of Eq. (2) can be split into a steady and an unsteady one. Furthermore, we assume only harmonic oscillations so that we can write

$$\phi(x,y,t) = \varphi(x,y) + \Psi(x,y) \cdot e^{ikt}$$
 (3)

where $k = \frac{\omega \cdot c}{U}$ is the reduced frequency.

Introducing (3) in (2), we can separate the unsteady from the steady problem and we obtain a set of two differential equations

$$[M^2 - 1 + (\kappa + 1) M^2 \varphi_{x}] \varphi_{xx} - \varphi_{yy} = 0$$
 (4)

and

$$[M^2 - 1 + (\kappa + 1) M^2 \varphi_x] \Psi_{xx} - \Psi_{yy} + [M^2 (\kappa + 1) \varphi_{xx} + 21k M^2] \Psi_x - M^2 k^2 \Psi = 0 \quad (5)$$

The boundary conditions for the flow over an oscillating airfoil can also be written as the sum of steady and unsteady influences:

$$h(x,t) = h_{1}(x) + h_{1}(x) \cdot e^{ikt}$$
 (6)

(h(x,t); x) is the true location of a surface point. Thus the boundary conditions for the steady and unsteady problem can be expressed

$$\varphi_y = \frac{\partial \mathbf{b}_0}{\partial \mathbf{x}} \equiv \text{slope of the surface}$$
 (7a)

$$\Psi_y = \frac{\partial h_1}{\partial x} + \frac{\partial h_1}{\partial t} =$$
the unsteady movement of a (7b)

Eq. (4) together with (7a) describes the transonic flow field over a fixed airfoil.

Following Sauer (8), we can attack the problem with the method of characteristics. The left- and right- running characteristics shall be indicated by α and β . We find for their slopes

$$\left(\frac{\partial y}{\partial x}\right)_{\alpha,\beta} = \pm \frac{1}{\sqrt{M^2 - 1 + (\kappa + 1) M^2 \varphi_x}}$$
 (8)

or with

$$\lambda = M^2 - 1 + (\kappa + 1) M^2 \varphi_{\mathbf{x}}$$
 (9)

$$\left(\frac{\partial y}{\partial x}\right)_{\alpha,\beta} = \pm \frac{1}{\sqrt{\lambda}} \tag{10}$$

The upper sign indicates the $\,\alpha$ - direction. Introducing a second substitution

$$\mu = \frac{3}{2} (\kappa + 1) M^2 \varphi_y$$
 (11)

The compatibility relation $\pm \sqrt{\chi} \varphi_{xx} + \varphi_{yx} = 0$

can now be written as
$$\left(\frac{\partial \lambda^{3/2}}{\partial \mathbf{x}}\right)_{\alpha,\beta} \mp \left(\frac{\partial \mu}{\partial \mathbf{x}}\right)_{\alpha,\beta} = 0$$
 (12)

That Eq. (12) holds, can easily be verified by resubstituting (9) and (11) in (12) and executing the differentation. The result will be Eq. (4).

In (12) we find only derivatives in the x-direction along the characteristics. Therefore we can integrate (12) easily and obtain

$$\lambda^{3/2} \mp \mu = \text{const} = C_{\alpha, \beta}$$
 (13)

We changed our variables from φ and φ to λ and μ . Consequently, we have to convert our boundary conditions for the steady problem:

y = 0:

$$\mu = \frac{3}{2} (\kappa + 1) M^2 \frac{\partial h_o}{\partial x}$$
 (14)

Eq. (13) says: as long as we move along one characteristic,

$$\left(\lambda^{3/2} \mp \mu\right)_{\alpha,\beta}$$

will not change. This makes (13) a tool to evaluate the original desired unknowns $\varphi_{_{\mathbf{X}}}$ and $\varphi_{_{\mathbf{Y}}}$ in the field.

In the free-stream field $\varphi_{\mathbf{x}}$ and $\varphi_{\mathbf{y}}$ are zero. Hence, there we have

$$\lambda_{\infty} = M^2 - 1$$

and

Therefore all the characteristics have here the slope

$$\left(\frac{\partial y}{\partial x}\right)_{\alpha,\beta} = \pm \frac{1}{\sqrt{M^2 - 1}}$$

and from (12) we obtain

$$C_{\infty} = \lambda_{\infty}^{3/2} = (M^2 - 1)^{3/2}$$
 (15)

Fig. 1 shows a β -characteristic of the free stream hitting the surface of the airfoil.

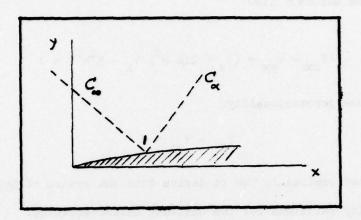


Fig. 1. Characteristics on the Airfoil Surface

In point 1 there has to be $C_{\infty} = C_{\alpha}$ with Eq. (9), (13) and (14) we now can find

$$\lambda_1 = (M^2 - 1)^{3/2} \pm \frac{3}{2} (\kappa + 1) M^2 \frac{\partial h_o}{\partial x}^{2/3}$$
 (16)

In the two-dimensional case there can be shown: for the characteristic leading away from the boundary there is not only

Ca = const

but also

μ_α = const and

 $\lambda_{\alpha} = const$

With Eq. (5) through (16) the steady problem (4) can be solved. We get a net of characteristics. On the gridpoints the values of $\varphi_{\mathbf{x}}$ and $\varphi_{\mathbf{y}}$ are known. The computational procedure is shown later.

With the knowledge of the characteristic net and the values for $\varphi_{\mathbf{x}}$, φ , λ and μ in each grid point we can now solve the unsteady part of the problem.

We put Eq. (9) and (11) into (5). The result is a new differential equation for the unsteady flow.

$$\lambda \Psi_{xx} - \Psi_{yy} + (\lambda_x + 2ik M^2) \Psi_x - M^2 k^2 \Psi = 0$$
 (17)

Further we assume irrotationality

$$\Psi_{xy} - \Psi_{yx} = 0 \tag{18}$$

Teipel shows explicitly how to derive from the system of Eq. (17) and (18) the compatibility relations for the unsteady characteristics.

$$\Psi_{xx} + \frac{1}{\sqrt{\lambda}} \Psi_{yx} + \frac{1}{\lambda} (\lambda_{x} + 2ik M^{2}) \Psi_{x} - \frac{1}{\lambda} k^{2}M^{2}\Psi = 0$$
 (19)

The geometry of the characteristic net is determined by the coefficients connected with the highest order terms. Thus we can see from Eq. (4) and (5) that the net remains the same for the unsteady flow problem.

To solve Eq. (17) via Eq. (19) by moving along the already known characteristics of the steady field, we need the unsteady boundary values along the airfoil and along the shock.

The first one is given with Eq. (7b).

$$\Psi_y = \frac{\partial h_1}{\partial x} + ikh_1$$

It reads for pitch - movement

$$\bar{\Psi}_{y} = - \left[1 + ik(x - b)\right]$$

and for plunge - movement

$$\Psi_{\mathbf{v}} = -i\mathbf{k}$$

where b is the normalized location of the pitching axis. These expressions are derived in section 2.2. The boundary conditions along the oscillating shock are much more difficult to obtain. This is done in section 2.3.

After obtaining the solution for \mathbf{u}_1 , \mathbf{v}_1 and Ψ the unsteady pressure coefficients can be computed from

$$c_{pl} = 2 \frac{p_{l}}{\rho_{\infty} U}$$

$$c_{pl} = -2(u_{l} + ik \Psi)$$
(20)

2.2 Boundary Conditions Along the Airfoil

The general expression of the location of an oscillating airfoil is given by Eq. (6)

$$h(x,t) = h_0(x) + h_1(x) \cdot e^{ikt}$$

 $h_Q(x)$ represents the surface and we assume for now that it is described by an analytical function for which the second derivative exists. With Eq. (7a) we can calculate φ_y on the airfoil:

$$\varphi_{y} = \frac{\partial h_{o}}{\partial x}$$

However, as we only consider slender bodies, we project the point down to the x-axis y=0, so that the boundary coordinates of the characteristic net are always (x,0) instead of (x,y). This has two reasons

- y = 0 makes the steady boundary step much less complicated, without introducing a considerable mistake;
- The oscillating movement of the airfoil can be reduced to the movement of a flat plate.

We examine two moving modes. The pitch- and plunge- mode.

PITCH:

Fig. 2 shows the deflected airfoil (flat plate) in a system of coordinates

As α is small and a harmonic motion, we can say

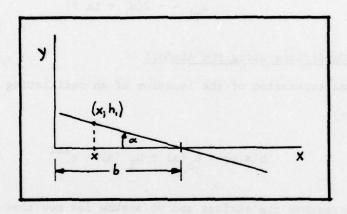


Fig. 2. Pitching Flat Plate

Thus we get

$$h_1 = (b - x) \alpha_0 \cdot e^{ikt}$$

The unsteady boundary condition is with Eq. (7b)

$$\Psi_y = [-\alpha_0 + \alpha_0 ik(b - x)] e^{ikt}$$

or

$$\Psi_{y} = -\alpha_{o} \left[1 + ik(x - b)\right]$$

where the exp (ikt) - term is omitted. This can be normalized by $|\alpha_0|$. So the final expression is the well known form

$$\Psi_{y} = -[1 + ik(x - b)]$$
 (21)

PLUNGE:

For the plunge mode the deflection h is no function of x . Again a harmonic motion is assumed.

Eq. (7b) gives us then the boundary value for

$$\Psi_y = -h_o \cdot ik e^{ikt}$$

To be consistent with the previous work (3 to 7), the downward deflection is defined as positive. Again the expression is normalized by $|\mathbf{h}_0|$ and the

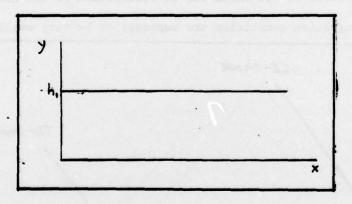


Fig. 3. Plunge Movement

exp(ikt) - term is omitted

$$\Psi_{y} = -ik \tag{22}$$

This approach to the boundary conditions is a rather physical one. A more rigorous derivation is shown by Bell in (9), including the derivation of Eq. (7b).

2.3 The Oscillating Shock Wave in an Oscillating Flow Field

The basic problem in using the method of characteristics is finding and introducing the proper boundary conditions. Section 2.2 gives us the influence of the moving airfoil into the field. The second boundary in the field between surface and shock are the unsteady flow properties immediately downstream of the shock. See Fig. 4.

It is

w - Velocity

u, v - x,y components of w

w.,wt - Normal and tangential components of

- Indicates properties behind the shock

 γ_{o} - Slope of the shock in the steady problem

 γ ' - Deflection of the shock due to oscillation of the airfoil

All the perturbation quantities are supposed to be very small.

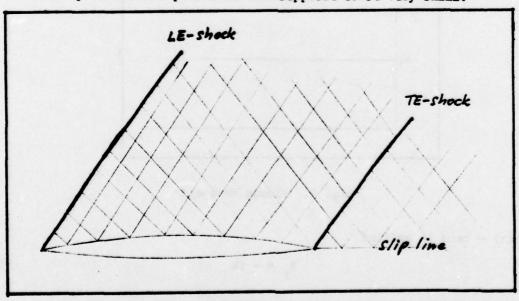


Fig. 4. Field of Characteristics Past an Airfoil

Teipel (10) found a way to obtain those properties for an airfoil in undisturbed supersonic flow. As the final object is to compute the flow in a staggered cascade, his work had to be extended. This was first done by Chadwick (4) who applied it to a cascade of wedges. Strada used Chadwick's equations in (5) to compute the inlet flow of a cascade with airfoils which have flat upper and curved lower surfaces.

Fig. 5 shows the velocities upstream and downstream of an arbitrary shock:

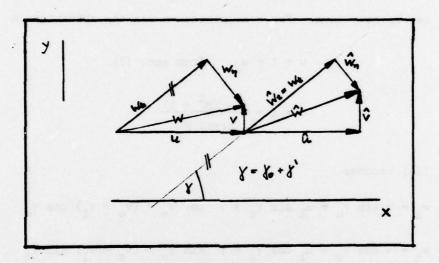


Fig. 5. Velocities at the Shock

The general case in the cascade is w # u . Thus we write

$$u = 1 + u_o + u_1$$

$$v = v_o + v_1$$

$$\hat{u} = 1 + \hat{u}_o + \hat{u}_1$$

$$\hat{v} = \hat{v}_o + \hat{v}_1$$
(23)

The only unknowns are $\hat{\mathbf{u}}_1$ and $\hat{\mathbf{v}}_1$, as the steady problem is already expected to be solved and the field in front of the shock is presumed to be well known. Then we can take from the geometry (Fig. 5).

$$w_n = u \cdot \sin (\gamma_0 + \gamma') - v \cos (\gamma_0 + \gamma')$$

$$w_t = u \cos (\gamma_0 + \gamma') + v \sin (\gamma_0 + \gamma')$$
(24)

If we apply Eq. (23) to (24) and neglect higher order terms like $u \gamma^1$, Eq. (24) can be rewritten. For convenience we use the abbreviation

$$v = 1 + u_0$$
 from eqn. (9)

$$v = 1 + \frac{\lambda - (M^2 + 1)}{(\kappa + 1) M^2}$$
(25)

Eq. (24) becomes

$$w_{n} = v \sin \gamma_{o} + u_{1} \sin \gamma_{o} + \gamma' \cos \gamma_{o} - (v_{o} + v_{1}) \cos \gamma_{o}$$

$$(26)$$

$$w_{t} = v \cos \gamma_{o} + u_{1} \cos \gamma_{o} - \gamma' \sin \gamma_{o} + (v_{o} + v_{1}) \sin \gamma_{o}$$

Now we take a look at the equation for normal moving shocks (11).

$$\hat{w}_{n} - w_{n} = \frac{2}{\kappa + 1} \cdot W(1 - \frac{a^{2}}{w^{2}}),$$
 (27)

where W is the relative velocity of the shock with respect to the fluid. To transform from the system moving with the fluid into an airfoil - fixed system of coordinates, we find the velocity of a point oscillating with the shock as

$$W^+ = W + w_n \tag{28}$$

See Fig. 6.

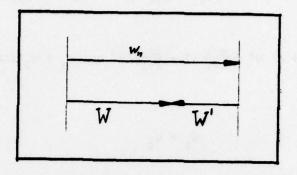


Fig. 6. Definition of Relative Velocities in a Shock-Point

Connecting Eq. (26) with Eq. (28) we obtain

$$W = W' - v \sin \gamma_0 - u_1 \sin \gamma_0 - \gamma' \cos \gamma_0 + (v_0 + v_1) \cos \gamma_0 \qquad (29)$$

Forming W^2 and again neglecting all terms of higher order, recasting Eq. (29) leads to

$$v^2 \sin^2 \gamma_o = W(W + 2v \sin \gamma_o)$$
 (30)

Reintroducing this into the shockpolar Eq. (27), we get

$$w_n - w_n = \frac{2}{\kappa + 1} W - \frac{2}{\kappa + 1} \cdot \frac{a^2}{v^2 \sin^2 \gamma_o} (W + 2 v \sin \gamma_o)$$
 (31)

By making use of Eq. (26), (29), (31) and the substitution

$$A = \frac{a_{\infty}^{2}}{v^{2} \sin^{2} \gamma_{o}} = \frac{1}{v^{2} M^{2} \sin^{2} \gamma_{o}} = \frac{1}{v^{2} M_{n}^{2}}$$
(32)

We obtain, again neglecting all higher order terms

$$\hat{w}_{n} = \frac{\kappa - 1}{\kappa + 1} \vee \sin \gamma_{o} \left(1 + \frac{2A}{\kappa - 1} \right) - \frac{\kappa - 1}{\kappa + 1} v_{o} \cos \gamma_{o} \left(1 - \frac{2A}{\kappa - 1} \right) +$$

$$+ \frac{2}{\kappa + 1} \left(1 + A \right) W^{+} + \frac{\kappa - 1}{\kappa + 1} \left(1 - \frac{2A}{\kappa - 1} \right) \left(\gamma^{+} \cos \gamma_{o} + u_{1} \sin \gamma_{o} - v_{1} \cos \gamma_{o} \right)$$
(33)

together with

$$\hat{\mathbf{w}}_{t} = \mathbf{w}_{t} \tag{34}$$

Out of Eq. (26) we have now expressions for the velocities behind the moving shock in terms of the known values and γ' and W'. In the next step we look at a point of the oscillating shock, Fig. 7.

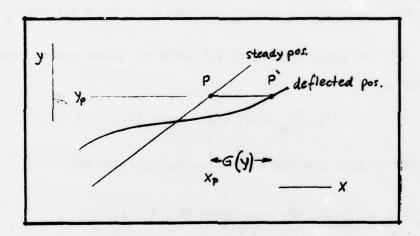


Fig. 7. Geometry of an Oscillating Shock

If we assume G to be a harmonic oscillation around the steady middleposition P, we can say

$$x = x_p + G(y) \cdot e^{ikt} , \qquad (35)$$

where G(y) is the amplitude of the shock vibration and k the same reduced frequency as given for the airfoil.

The normal motion of P is

$$W' = \left(\frac{\partial x}{\partial t}\right)_{y} \cdot \sin \left(\gamma_{o} + \gamma\right)$$

$$W' = ik \sin \gamma_{o} G(y)$$
(36)

again the exponential term is omitted.

As we know $\gamma_0(y)$ from the solution of the steady problem, we can find the x-coordinates of P:

$$x_p = \int_0^y ctg (\gamma_o (y)) dy$$

This completes Eq. (35) to

$$x = \int_{0}^{y} ctg \left(\gamma_{o} (y) \right) \cdot dy + G(y) \cdot e^{ikt}$$
 (37)

The total differential gives us

$$dx = \left(\frac{\partial x}{\partial t}\right)_y \cdot dt + \left(\frac{\partial x}{\partial y}\right)_t \cdot dy$$

$$dx = ik G(y) \cdot e^{ikt} \cdot dt + \left[etg \left(\gamma_{o}(y)\right) + \left(\frac{\partial G}{\partial y}\right)e^{ikt}\right] \cdot dy$$

with

$$\left(\frac{\partial \mathbf{x}}{\partial \mathbf{y}}\right)_{\mathbf{t}} = \operatorname{ctg}\left(\gamma_{0}(\mathbf{y}) + \gamma'(\mathbf{y})\right)$$

and now considering always a fixed y we can formulate

ctg
$$(\gamma_0 + \gamma') = ctg \gamma_0 + G_y \cdot e^{ikt}$$

Trigonometric relations and higher order terms going to zero give us

$$\gamma' = -\sin^2 \gamma_0 G_y \tag{38}$$

The transformation back to our Cartesian system of coordinates (Fig. 1) is done by

$$\hat{\mathbf{u}} = \hat{\mathbf{w}}_{t} \cos \left(\gamma_{o} + \gamma' \right) + \hat{\mathbf{w}}_{n} \sin \left(\gamma_{o} + \gamma' \right)$$

$$\hat{\mathbf{v}} = \hat{\mathbf{w}}_{t} \sin \left(\gamma_{o} + \gamma' \right) - \hat{\mathbf{w}}_{n} \cos \left(\gamma_{o} + \gamma' \right)$$
(39)

The following steps have to be executed on Eq. (39):

- Introducing Eq. (33), (34)
- Applying trigonometric relations
- Neglecting all higher order products
- Introducing Eq. (36) and (38)
- Separating the purely steady expressions from the rest of the equations

The result is a substitute for the shock polar (27) which gives us the unsteady velocities u_1 and v_1 behind the shock.

$$\hat{\mathbf{u}}_{1} - \mathbf{m}_{1} \quad \mathbf{G}_{y} - \mathbf{i} \quad \mathbf{m}_{2} \quad \mathbf{G} - \mathbf{m}_{3} \quad \mathbf{u}_{1} - \mathbf{m}_{4} \quad \mathbf{v}_{1} = 0$$

$$\hat{\mathbf{v}}_{1} - \mathbf{n}_{1} \quad \mathbf{G}_{y} - \mathbf{i} \quad \mathbf{n}_{2} \quad \mathbf{G} - \mathbf{n}_{3} \quad \mathbf{u}_{1} - \mathbf{n}_{4} \quad \mathbf{v}_{1} = 0$$
(40)

with A defined in Eq. (32) it is

$$m_{1} = \frac{2}{\kappa+1} \vee \sin^{2} \gamma_{0} \sin^{2} \gamma_{0}$$

$$m_{2} = \frac{2k}{\kappa+1} (1 + A) \sin^{2} \gamma_{0}$$

$$m_{3} = \sin^{2} \gamma_{0} \frac{\kappa-1}{\kappa+1} (1 - \frac{2A}{\kappa-1}) + \cos^{2} \gamma_{0}$$

$$m_{4} = \frac{\sin^{2} \gamma_{0}}{\kappa+1} (1 + A)$$
(41)

$$n_{1} = -\frac{2v}{\kappa+1} (\cos \gamma_{o} + A) \sin^{2} \gamma_{o}$$

$$n_{2} = -\frac{2k}{\kappa+1} (1 + A) \cos \gamma_{o} \sin \gamma_{o}$$

$$n_{3} = m_{4}$$

$$n_{4} = \sin^{2} \gamma_{o} + \cos^{2} \gamma_{o} \frac{\kappa-1}{\kappa+1} \left(1 - \frac{2A}{\kappa-1}\right)$$
(42)

These are exactly the coefficients Chadwick presented in (4). For the airfoil in undisturbed supersonic flow the perturbation quantities \mathbf{u}_1 and \mathbf{v}_1 in front of the shock are zero. With this in mind, Eq. (40) reduces to the shockpolar derived by Teipel in (10). It should be stated explicitly at this point that we followed Teipel very closely in this extension of his work and that Chadwick indicates the way in (4).

We have to do a last step, to make Eq. (40) a tool for computing the unsteady boundary values along the shock:

On the leading edge we know G=0 because the shock is always attached. $\hat{v}_1 \Big|_{y=0}$ is known from the boundary conditions on the airfoil, Eq. (21). Now we can isolate $\hat{u}_1 \Big|_{y=0}$ in Eq. (40):

$$G_{y|_{y=0}} = \frac{1}{n_{1}} \left[\hat{v}_{1|_{y=0}} - (n_{3} u_{1} + n_{4} v_{1}) \right]$$
 (43)

and

$$\hat{\mathbf{u}}_{1}\Big|_{\mathbf{y}=\mathbf{0}} = \frac{\mathbf{m}_{1}}{\mathbf{n}_{1}} \hat{\mathbf{v}}_{1}\Big|_{\mathbf{y}=\mathbf{0}} + \mathbf{n}_{1} \left(\mathbf{m}_{3} - \frac{\mathbf{m}_{1}}{\mathbf{n}_{1}} \cdot \mathbf{n}_{3}\right) + \mathbf{v}_{1} \left(\mathbf{m}_{4} - \frac{\mathbf{m}_{1}}{\mathbf{n}_{1}} \mathbf{n}_{4}\right)$$
(44)

After this initial step we develop finite differences along the steady shock to solve gradually for the unsteady values of \hat{u}_1 , \hat{v}_1 and G as shown later.

2.4 Connections Between the Linear and the Nonlinear System of Equations

In an earlier work Teipel developed in (12) a method of characteristics for an oscillating single flat plate. He derived an analytical solution for the unsteady boundary values along the shock. Using the perturbation velocity of sound rather than the perturbation potential, his concept was taken by Bell in (9) and by Platzer and associates in (3,13,14) to obtain results for a cascade of flat plates. They started from the Euler and continuity equations with the substitution

$$U(x,y) \cdot e^{ikt} = u, \qquad (45a)$$

$$V(x,y) \cdot e^{ikt} = \frac{1}{\sqrt{M^2 - 1}} \cdot v_1 \tag{45b}$$

$$C(x,y) \cdot e^{ikt} = \frac{2}{\kappa - 1} \frac{1}{M^2} \frac{a - a_{\infty}}{a_{\infty}}$$
 (45c)

where U, V and C are complex nondimensional amplitudes. The result is a set of differential equations which reads

Continuity:
$$\frac{\partial \mathbf{y}}{\partial \mathbf{x}} + \sqrt{\mathbf{M}^2 - 1} \cdot \frac{\partial \mathbf{y}}{\partial \mathbf{y}} + \mathbf{M}^2 \frac{\partial \mathbf{C}}{\partial \mathbf{x}} + i\mathbf{k} \mathbf{M}^2 \mathbf{C} = 0$$
 (46)

Euler:
$$\frac{\partial U}{\partial x} + \frac{\partial C}{\partial x} + ikU = 0$$
 (47)

Irrotationality:
$$\frac{\partial U}{\partial x} - \sqrt{M^2 - 1} \frac{\partial V}{\partial x} = 0$$
 (48)

Furthermore it is shown by Bell in (9) that the pressure in terms of Eq. (45) can be expressed as

$$p - p_{\infty} = \frac{2}{\kappa - 1} \rho_{\infty} a_{\infty} (a - a_{\infty})$$

or

As $(p - p_{\infty})$ is our unsteady pressure disturbance (the steady pressure disturbance is zero for a flat plate) we can say with Eq. (20)

$$2C = -2 (u_1 + ik \Psi)$$

$$C = - (u_1 + ik \Psi)$$
(50)

Now we take Eq. (45a), (45b) and (50) and substitute into Eq. (46) and we get

$$(M^2 - 1) \frac{\partial u_1}{\partial x} - \frac{\partial v_1}{\partial y} + 2 ik M^2 u_1 - k^2 M^2 \Psi = 0$$
 (51)

With $u_1 = \Psi_x$ and $v_1 = \Psi_y$ this is our Eq. (17), when we consider that for a flat plate the values for φ_x and λ_x are always zero.

$$\lambda = \sqrt{M^2 - 1}$$

Applying Eq. (50) on (47) gives us

or

$$u_1 = \frac{9x}{9A}$$

which is one of our basic equations. Obviously, Eq. (48) can be transformed into Eq. (18). Hence it is shown that the basic equations previously used by Teipel and Platzer for the flat plate are equivalent to the perturbation potential equations used in this work, when these are reduced to the linear case.

The next step shows how to obtain an anlytical expression for the oscillating shock generated by a single flat plate from the rather complicated differential equation (40).

As there are no perturbations in front of the shock, Eq. (40) reduces to the Teipel-form

$$\hat{u}_1 = m_1 G_y + i m_2 G$$
 $\hat{v}_1 = n_1 G_y + i n_2 G$

In addition the steady shock angle γ_0 can be expressed as

$$\sin \gamma_{o} = \frac{1}{M} \tag{52}$$

and with this A from Eq. (41), defined in Eq. (32), becomes

$$A = \frac{1}{v^2 M^2 \sin^2 \gamma_0} = 1$$

$$(v = 1 + u_0 = 1)$$
(53)

The shockpolar written down explicityly reads

$$\hat{\mathbf{u}}_{1} = \frac{2}{\kappa+1} \sin^{2} \gamma_{0} \sin^{2} \gamma_{0} G_{y} + i \frac{2k}{\kappa+1} 2 \sin^{2} \gamma_{0} G$$

$$\hat{v}_1 = -\frac{2}{\kappa+1} (\cos 2\gamma_0 + 1) \sin^2 \gamma_0 G_y - i \frac{2k}{\kappa+1} 2 \cos \gamma_0 \sin \gamma_0 G$$

Using the trigonometrical relations and Eq. (52) we obtain after some adding and subtracting

$$\hat{\mathbf{u}}_{1} = -\frac{1}{\sqrt{M^{2} - 1}} \hat{\mathbf{v}}_{1} \,. \tag{54}$$

for the velocities immediately downstream of the shock. We now go into the compatibility relation Eq. (19). Linearized with $\lambda_{\mathbf{x}} = 0$ it is for the upper side

$$\frac{\partial u_1}{\partial x} - \sqrt{\frac{1}{M^2 - 1}} \frac{\partial v_1}{\partial x} + \frac{2\pm k}{M^2 - 1} u_1 - \frac{1}{M^2 - 1} k^2 M^2 \Psi = 0$$

Eq. (54) shows that

$$\frac{\partial \mathbf{u_1}}{\partial \mathbf{x}} - \frac{1}{\sqrt{\mathbf{M}^2 - 1}} \frac{\partial \mathbf{v_1}}{\partial \mathbf{x}} = 2 \frac{\partial \mathbf{u_1}}{\partial \mathbf{x}}$$

Landahl (1) shows that $(\Psi + \varphi)$ should be continuous across a shock. As all the steady components are zero in this case his result applies here to the unsteady potential alone. On the other hand we consider the shock thickness as infinitely small. Assumming $\hat{d\Psi}$ to be nonzero,

would become infinitely large across the shock. Thus $\widehat{\Psi}$ can only be zero along the unsteady shock of a single flat plate, or

for a cascade.

The remaining steps are fairly straight forward. With the conditions shown above, Eq. (19) becomes

$$2\frac{\partial \mathbf{u_1}}{\partial \mathbf{x}} + \frac{2\mathbf{i}\mathbf{k}}{\mathbf{M}^2} \mathbf{u_1} = 0$$

or

$$\frac{\partial u_1}{\partial x} = -ik \frac{M^2}{M^2 - 1} u_1$$

this can be integrated

$$u_1 = const \cdot e - i \frac{k M^2}{M^2 - 1} x$$

for x = 0 we get the perturbation for the leading edge

$$\hat{\mathbf{u}}_{LE} = -\frac{1}{\sqrt{M^2 - 1}} \cdot \hat{\mathbf{v}}_{LE}$$

can be obtained from the unsteady boundary conditions. The final LE expression for the unsteady velocities behind the shock attached to an oscillating flat plate is now

$$\hat{u}_{1}(x) = \hat{u}_{1_{LE}} \cdot e^{-i\frac{kM^{2}}{M^{2}-1}x}$$
 (55)

This result was also obtained by Bell (9) and considering the different systems of coordinates, it is equal to the solution found by Teipel in (12).

2.5 The Wake

Fig. 8 shows the characteristic net for the airfoil and wake regions. The fluid in the fields 2 and 11 has to adjust via the two trailing edge shocks in such a way, that the wake condition on the slip-line in the fields 4 and 13 is not violated. This condition requires flow tangency, continuity of pressure and normal velocity across the slip-line, i.e.,

$$\theta_4 = \theta_{13}$$
 and $cp_4 = cp_{13}$ (56) $w_{n_4} = w_{n_{13}}$

w is the velocity component normal to the slip-line. Again this problem is split into steady and unsteady parts.

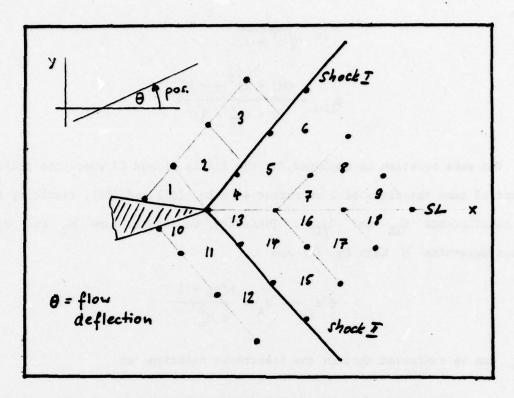


Fig. 8. The Net of Characteristics for the Steady Wake

Shapiro (11) gives a series solution for the pressure difference connected to a direction-change $\Delta\theta$ = θ_4 - θ_2 for supersonic flow.

$$2 \frac{p_4 - p_2}{\kappa p_4 M_2^2} = \pm c_{Iu} \cdot (\theta_4 - \theta_2) + c_{IIu} \cdot (\theta_4 - \theta_2)^2$$
 (57)

- + left running Mach lines
- right running Mach lines

with

$$c_{Iu} = \frac{2}{\sqrt{M^2 - 1}}$$

$$c_{IIu} = \frac{(M_2^2 - 2)^2 + \kappa M_2^4}{2 \cdot (M_2^2 - 1)^2}$$
(58)

The same equation is employed for the fields 11 and 13 where the indices 11 and 13 take the place of 2 or rather 4 in Eq. (57) and (58), resulting in the coefficients C_{IL} and C_{IIL} . Therefore, we need to know M_2 and M_{11} . We can determine M with Eq. (9) and

$$\frac{1}{2}c_{p} = -\varphi_{x} = \frac{p/p_{\infty} - 1}{\kappa M_{\infty}^{2}}$$

 p/p_{∞} can be expressed through the isentropic relations as

$$p/p_{\infty} = f (\kappa, M_{\infty}, M)$$

with this we obtain

$$M^{2} = \frac{2}{\kappa - 1} \left[\frac{1 + \frac{\kappa - 1}{2} M_{\infty}^{2}}{\left[1 - \frac{\kappa}{\kappa + 1} (\lambda - M_{\infty}^{2} + 1) \right]^{\frac{\kappa - 1}{\kappa}}} - 1 \right]$$

For a given point in which λ is known we can find the flow direction for $\theta_m = 0$ as parallel to the x-axis.

$$tg \theta = \frac{v_o}{1 + u_o}$$

from Eq. (11)

$$v_0 = \frac{2}{3} \frac{\mu}{(\kappa+1) M_{max}^2}$$

and (13)

$$\mu = \mp \left[(M_{\infty}^2 - 1)^{3/2} - \lambda^{3/2} \right]$$

we arrive through Eq. (25) at

$$tg \Theta = \mp \frac{2}{3} \cdot \frac{(M_{\infty}^2 - 1)^{3/2} - \lambda^{3/2}}{\lambda + M_{\infty}^2 \kappa + 1}$$
 (60)

The upper sign indicates the left-running (- lower surface) and the lower sign the right running characteristic (= upper surface).

Now we have the coefficients C_{Iu} , C_{IIu} , C_{IL} and C_{IIL} . In addition it must be

and

Thus we can say

$$\frac{P_4}{P_{\infty}} = \frac{1}{2} \times M_{\infty}^2 \left[c_{1u} \cdot (\Theta_4 - \Theta_2) + c_{1Iu} \cdot (\Theta_4 - \Theta_2)^2 \right] + 1$$

$$\frac{P_{13}}{P_{\infty}} = \frac{1}{2} \times M_{\infty}^2 \left[-c_{1L} \cdot (\Theta_4 - \Theta_{11}) + c_{1IL} \cdot (\Theta_4 - \Theta_{11})^2 \right] + 1$$
(61)

The equality of these two equations gives us a way to solve for θ_4 . As we took only second order terms of $\Delta\theta_1$ the result is $\theta_4 = \theta_{13} = 0$ for zero angle of attack. This is correct, because we assume isentropic flow across the shocks. The reason why it is done in this rather difficult way is that the upwash of the nonisentropic flow can be easily added to the program by simply adding the third coefficient C_{III} given by Shapiro. Actually, this is already done in the program. It has been reversed through setting $C_{III} = 0$ in order to be consistent with the assumption of isentropic flow across the weak shocks.

After we have determined the flow in field 4 and 13 we compute the trailing edge shocks and the rest of the wake field. The step from the slip-line is repeated from the fields 5 and 14 to 7 and 16 as described, etc. All other steps in the wake field are general characteristic steps like those over the airfoil.

It is now possible to solve for the unsteady flow field of the wake. This is again done on the same grid locations which are known from the steady solution.

The conditions for this part of the solution are

- No Pressure-Jump Across the Slip-Line

$$u_{12u} + ik \Psi_{2u} = u_{12L} + ik \Psi_{2L}$$
 (62)

(see Fig. 10).

- Tangential Velocity on the Slip-Line

$$\frac{\mathbf{v_o} + \mathbf{v_1}}{1 + \mathbf{u_o} + \mathbf{u_1}} = \mathsf{tg} \; (\varepsilon_o + \varepsilon') \tag{63}$$

(see Fig. 9)

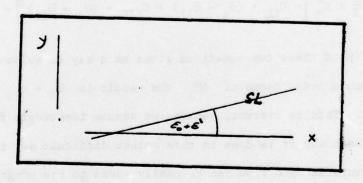


Fig. 9 Slip-Line Geometry

With trigonometric relations Eq. (63) becomes

$$\varepsilon' = v_0 - (1 + u_0) \operatorname{tg} \varepsilon_0 + v_1 - u_1 \operatorname{tg} \varepsilon_0$$

from the steady solution we know

$$v_0 - (1 + u_0) tg \epsilon_0 = 0$$

hence we obtain for the unsteady deflection of the slip-line

$$\varepsilon' = v_1 - u_1 \operatorname{tg} \varepsilon_0$$

as ε' has to be equal for the upper and the lower side of the flow field,

$$(v_1 - u_1 \cdot tg \varepsilon_0)_{2u} = (v_1 - u_1 tg \varepsilon_0)_{2L}$$
 (64)

is the first equation in the system, which we want to solve for $(u_1, v_1)_{2u}$ and $(u_1, v_1)_{2L}$. Indices refer to Fig. 10.

The next step uses Eq. (64). Fig. 10 shows the simultaneous step from the upper and lower wake field to the slip-line

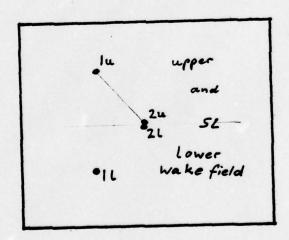


Fig. 10. Slip-Line Step

with

$$\Psi_2 = \Psi_1 + \frac{\mathbf{x}_2 - \mathbf{x}_1}{2} \left(\mathbf{u}_{12} + \mathbf{u}_{11} \pm 2 \frac{\mathbf{v}_{12} + \mathbf{v}_{11}}{\sqrt{\lambda_1} + \sqrt{\lambda_2}} \right)$$
 (65)

We can substitute ψ_2 by the known ψ_1 and the desired unknowns. This gives us the second equation of the system when properly recast in left- and right- hand sides.

The two missing statements which make the system solvable comes from the compatibility relation (19). Applied to the upper and lower slip-line steps, respectively, they complete the system.

The computational procedure is explained in more detail in chapter 3.

3. THE SINGLE OSCILLATING AIRFOIL: WING

3.1 The Program Organization Over the Airfoil

The name of the program for the single oscillating airfoil is W I N G. It was designed to use the subroutines as often as possible, that is over the surface as well as in the wake field. For this purpose a special notation was introduced which is shown in Fig. 11 and Fig. 12.

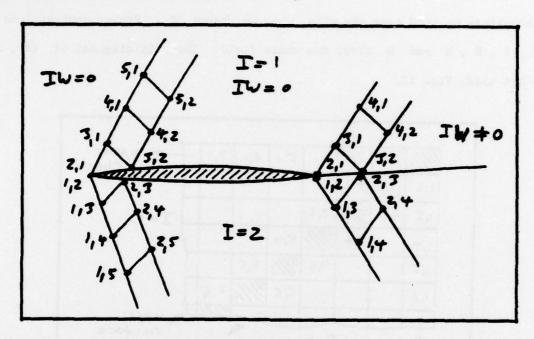


Fig. 11. Notation of the Fields in Wing

The arrays which contain properties of the flow field have the dimension (K, M, N). M and N give the location of the gridpoint: Point M on the N-th characteristic for the upper field and vice versa for the lower field.

K and IW indicate where we are:

IW = 0 Not yet in the wake

IW = 0 Computation of the wake

K = 1 Field already known

K = 2 Field computed now

The upper and lower side is indicated by I=1 and I=2 respectively. Thus we can switch the changing sign of the characteristics with

$$(-1)^{(I+1)}$$

This notation has the advantage that we can compute the field on both sides of the wing (IW = 0) by switching I. For the wake we still can use the subroutines we used over the wing, when we change K. Proper combinations from IW, I, K, M and N cover the whole field. The main diagonal of (M, N) is not used, Fig. 12.

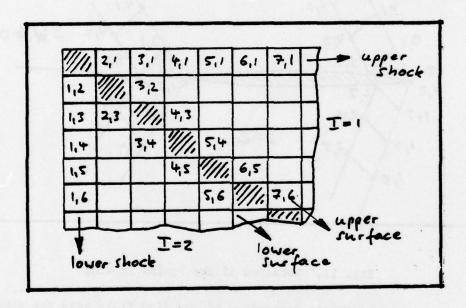


Fig. 12. Array Organization in Wing

Originally for the first shock we need only the Teipel - form of Eq. (2.40). But we have to use the whole expression in our step from the airfoil into the wake because the unsteady velocities over the airfoil are nonzero. So the sub-routine R A N D S includes the extended shockpolar. To be consistent in our treatment of the shock at the leading- and trailing edge, an initial field in

front of the airfoil is generated by the characteristics with the slope

$$\frac{dy}{dx} = \pm \frac{1}{\sqrt{M^2 - 1}}$$

all perturbation properties are set to zero. During the computation of the flow field over the airfoil (IW = 0) the index K has the meaning

K = 1 Initial field

K = 2 Field over the airfoil

After this calculation step we do not need the initial field any longer and we copy (2, M, N) to (1, M, N). Now the K=2 - array is free for the wake field results and K indicates

K = 1 Field over the airfoil

K = 2 Wake field

In the shown listing of W I N G the arrays have the size (2, 25, 25) to keep the storage area reasonable. However, with 16 points on the airfoil the wake is only computed to a distance of 1.7 times the chordlength from the leading edge. If this is not enough, the program can be blown up by increasing the arrays to (2, KV, KV).

The value of KV has to correspond with the first statement of the program. All the DO-loops are then dimensioned correctly. Another change has to be done in subroutine L I F T. Here all the arrays stay the same besides A. This is only a dummy-array to define the space between X, U, Q, P and PX, PS, PU. A must have the size

$$A(22 \cdot KV^2 + 4 \cdot KV - 1000)$$

The next page shows an approximate block diagram of the program. On the sides are the names of the subroutines involved in the step of the respective block. This diagram is only a summary of what is really done. A lot of

details have been skipped in order to keep it clear. One of the most important routines does not even appear. FIND is an orientation subprogram that works only on the field (1, M, N). If the arbitrary coordinates (x,y) are the input of FIND, it comes out with the indices (M2, N2) which give us the mesh of the field in which (x,y) is positioned; see Fig. 13.

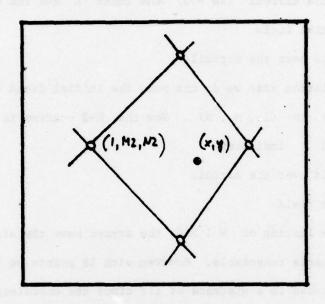
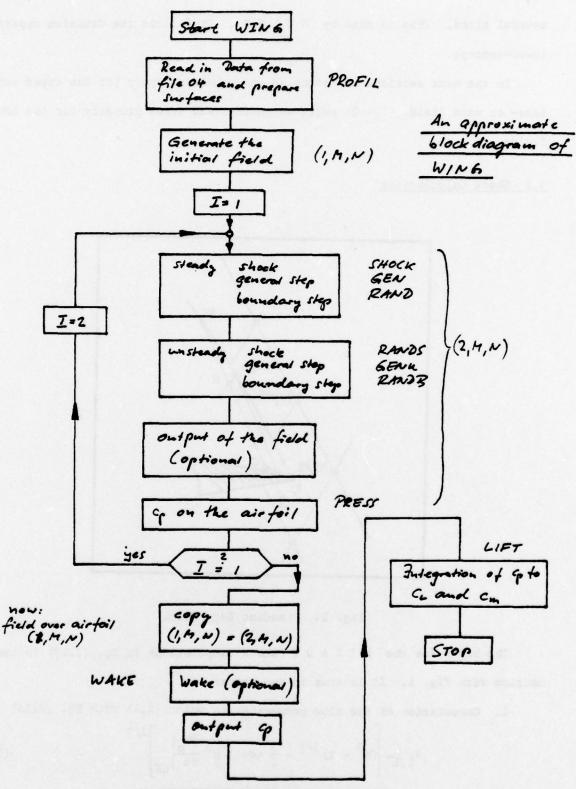


Fig. 13. A Characteristic Mesh

Again I is the switch for upper and lower sides. If (1, M, N) does not contain (x,y), FIND indicates this by setting IE = 1, which is otherwise zero. Besides this there is a message in the output. IE = 1 causes for the LE-shock to assume free stream values in front of it. For the TE-shock it terminates further computations, because there would be no sense in it. If (1, M, N) contains (x,y) and (1, M2, N2) can not be found, this will be in the output also. Then we know, something went wrong. But this should not happen. It is only a precautionary feature.

Another important background - subroutine is SOLVE. During the run through WING, complex systems of linear equations have to be solved



AN APPROXIMATE BLOCK DIAGRAM OF WING

several times. This is done by SOLVE. It follows the Gaussian upper/lower-concept.

In the next sections all examples and equations apply for the upper surface- or wake field. I = 2 switches indices and signs properly for the lower fields.

3.2 Shock Calculations

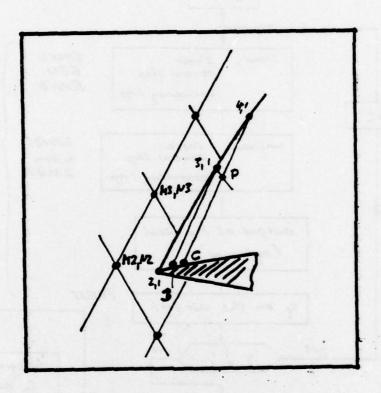


Fig. 14. Leading Edge Shock

The basis for the STEADY shock computations is Eq. (2.13) in connection with Fig. 1. It is done in two steps:

1. Computation of the flow properties in point (2,1) with Eq. (2.16)

$$\lambda_{2,1} = \left[(M^2 - 1)^{3/2} - \frac{3}{2} (\kappa + 1) M^2 \frac{\partial h_o}{\partial x} \Big|_{LE} \right]^{2/3}$$
 (1)

which is all we want to know for a point in the steady field. The slope of the shock at the LE is taken as the average of the characteristic slopes upstream and downstream of the shock:

$$tg \gamma_{o_{2,1}} = \frac{2}{\sqrt{\lambda_{2,1} + \sqrt{\lambda_{M3, N2}}}}$$
 (2)

2. All the other points of the shock are found by the same step done repeatedly:

From the LE we make a step DX(I) to find B (Fig. 14). Employing $\frac{\partial h_o}{\partial x}|_B$ in Eq. (2.16) gives us

$$\lambda_{3,1} = \lambda_{B}$$

and so the slope of the characteristics starting in B

$$tg \alpha_{B} = \frac{1}{\sqrt{\lambda_{B}}}$$

The intersection of the shock with the slope tg γ and the B - $^{\circ}2,1$ characteristic is the point (3,1). Here the slope of the shock is changed to

tg
$$\gamma_{o_{3,1}} = \frac{2}{\sqrt{\lambda_{3,1} + \sqrt{\lambda_{M3, N2}}}}$$

we apply the same procedure for point C respectively (4,1); etc. for the whole shock.

The points (M2,N2) or rather (M3,N3) are part of the initial field for the LE-shock. In case of the TE-shock they are part of the field over the airfoil. They are spotted in FIND.

All this is done in SHOCK.

The unsteady shock can be determined also with two basic steps:

1. The initial step was already described with Eq. (2.43) and (2.44). After this we have $\widehat{\mathbf{u}}_1$ and $\widehat{\mathbf{v}}_1$ at the origin of the shock

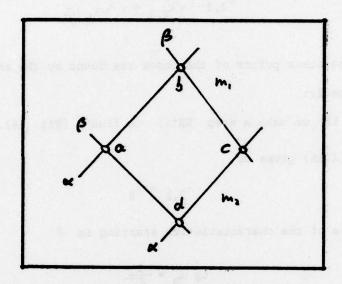


Fig. 15. Computation of λ_{x}

2. For the remaining points we take the two Eq. (2.40) (Shockpolar) and Eq. (2.19) (Compatibility Relations on the Characteristics Between (3,1) and (4,1); see fig. 14).

This gives us a complex system of three linear equations with the solution \hat{u}_1 , \hat{v}_1 and G in each point. The system is solved by subroutine SOLVE. The coefficients of Eq. (2.40) are computed in COEFF1.

However, to express Eq. (2.19) in finite difference form, we have some difficulties with $\lambda_{\mathbf{x}}$. Consider Fig. 15.

If we are moving from a to b , $\Delta\lambda$ is zero because we are on a left-running characteristic.

 λ is in point a

$$\lambda_{\mathbf{x}|_{\alpha}} = \frac{\lambda_{\mathbf{c}} - \lambda_{\mathbf{a}}}{\mathbf{x}_{\mathbf{c}} - \mathbf{x}_{\mathbf{a}}}$$

through $\lambda_c = \lambda_\alpha$ and approximately $x_c - x_a = 2 (x_c - x_d)$, we can say

$$\lambda_{\mathbf{x}}\Big|_{\alpha} = \frac{1}{2} \frac{\lambda_{\mathbf{d}} - \lambda_{\mathbf{a}}}{\mathbf{x}_{\mathbf{d}} - \mathbf{x}_{\mathbf{a}}}$$

or

$$\lambda_{\mathbf{x}}|_{\alpha} = \frac{1}{2} \lambda_{\mathbf{x}}|_{\beta}$$

We go now back to Fig. 14 and Eq. (2.19). Assuming that setting $\lambda_{\mathbf{X}}(3,1) = \lambda_{\mathbf{X}}(P)$ causes only a small error because both points are very close together, we can express $\lambda_{\mathbf{X}}(P)$ as

$$\lambda_{\mathbf{x}}(\mathbf{P}) = \frac{1}{2} \frac{\lambda_{4,1} - \lambda_{3,1}}{\mathbf{x}_{\mathbf{p}} - \mathbf{x}_{3,1}}$$

With this the finite difference form of Eq. (2.19) along the α - characteristic from P to (4,1) is no problem. It is written without the second index. Thus 3,1 becomes 3:

$$\frac{\hat{\mathbf{u}}_{14} - \hat{\mathbf{u}}_{13}}{\mathbf{x}_4 - \mathbf{x}_p} - \frac{\hat{\mathbf{v}}_{14} - \hat{\mathbf{v}}_{13}}{\sqrt{\lambda_4}} + \frac{\hat{\mathbf{u}}_{14} + \hat{\mathbf{u}}_{13}}{\sqrt{\lambda_4} + \sqrt{\lambda_3}} \cdot \left[\frac{1}{2} \frac{\lambda_4 - \lambda_3}{\mathbf{x}_p - \mathbf{x}_3} + 2ik \, \mathbf{M}^2 \right] - k^2 \mathbf{M}^2 \frac{\hat{\mathbf{v}}_4 + \hat{\mathbf{v}}_3}{\sqrt{\lambda_4} + \sqrt{\lambda_3}} = 0$$
(4)

With Eq. (2.64) substituted for $\hat{\Psi}_4$ and the \hat{u}_3 , \hat{v}_3 and $\hat{\Psi}_3$ already known, this is the third equation of the system after recasting into right and left hand sides.

After obtaining the results with SOLVE, $\hat{\Psi}_4$ is calculated separately. For the initial step at the shock origin we set

$$\hat{\Psi}_2 = \Psi_{M2,N2}$$

The procedure for the trailing edge shock is exactly the same. In the boundary values tg ϵ_0 is substituted for

The unsteady boundary values of the field behind the shock are computed in subroutine $R\ A\ N\ D\ S$.

3.3 The Boundary Step

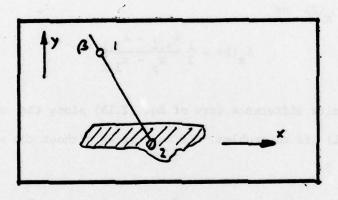


Fig. 16. Boundary Step

Fig. 16 shows the step from the field to the surface. It is done in R A N D . On the β - characteristic from 1 to 2 we have to satisfy two equations: Eq. (2.16)

$$\lambda_2 = \left[\lambda_{\infty}^{3/2} - \frac{3}{2} \left(\kappa + 1 \right) M^2 \frac{\partial h_o}{\partial x} \right|_2^{2/3}$$
 (5)

and

$$\frac{y_2 - y_1}{x_2 - x_1} = -\frac{2}{\sqrt{\lambda_2 + \lambda_1}}$$

To make it easier, we set y2 to zero and get

$$\lambda_2 = \left(\frac{2(x_2 - x_1)}{\gamma_1} - \lambda_1\right)^2 \tag{6}$$

Setting Eq. (6) equal to (5), we can find by iteration the \mathbf{x}_2 which matches both of them. With \mathbf{x}_2 we then find λ_2 . Here the steady problem is solved. The unsteady part is very straightforward. The known k gives us Ψ_1 this point via the unsteady boundary conditions Eq. (2.21) or (2.22). Now the only unknown left is Ψ_1 . It can be separated from the compatibility form. Again Ψ_2 is obtained separately afterwards from Eq. (2.64). The unsteady step is computed in RANDB.

3.4 The General Step

 $\lambda_{\rm b}$ and $\lambda_{\rm d}$ are known (Fig. 15). From the α - characteristic we know: $\lambda_{\rm c} = \lambda_{\rm d}$

Finding the point c for the steady case is only a geometrical problem.

It is the intersection of the two lines with the slopes

$$m_1 = -\frac{2}{\sqrt{\lambda_b} + \sqrt{\lambda_c}}$$
 and

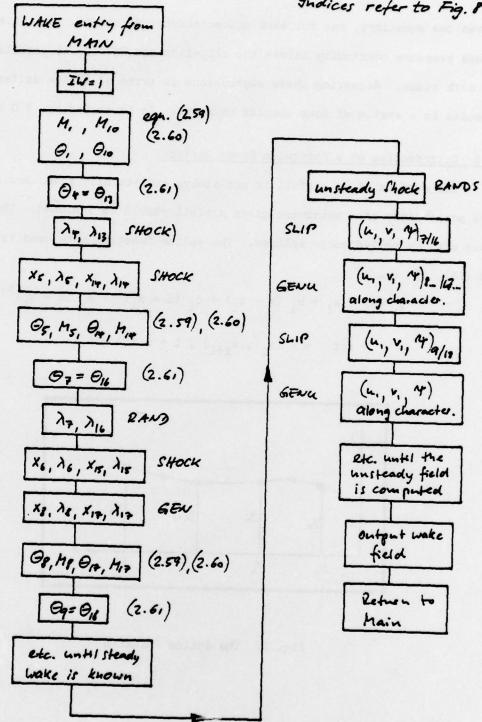
$$m_2 = \frac{1}{\sqrt{\lambda_c}}$$

they run through the points b and d. This is done in GEN.

3.5 Computation of the Wake

For the wake computation W I N G leaves the main program completely and works in a new organization program, called W A K E . This was mainly done to separate the entirely different computation sequence from the organization from the field above the surface. In section 2.5 it was already indicated how to solve for the slip-line and then for the wake field. The wake field computation is shown in the following flow diagram. The indices refer to Fig. 8. Attached to the blocks are the names of the subroutines or the equations involved in the step.

Once the steady field is known, the calculation becomes straightforward again. The conditions directly downstream of the TE are defined by the unsteady shockpolar (2.40). Applied simultaneously for the upper and lower airfoil-side, they give the conditions for pressure continuity (2.62) and parallel flow on the slip-line (2.64). With these four equations we can solve for $(\hat{u}, \hat{\tau})_{up}$ and $(\hat{u}, \hat{\tau})_{lo}$. This is the slightly more difficult initial step for the TE-shock. After it is done, the computation runs along the shock or rather along the characteristics with the same subroutines as over the airfoil. The basic difference lies in the step to the slip-line. Here RANDB can not be used any longer. SLIP does the simultaneous step from the upper and the lower wake field to the slip-line. The compatibility relation (2.19)



An Approximate Block Diagram for the Wake Computation

gives two equations, one for each characteristic direction. Eq. (2.62) demands pressure continuity across the slip-line and Eq. (2.64) parallel flow on both sides. Recasting these expressions in terms of finite differences results in a system of four complex unknowns. It is solved by SOLVE.

3.6 Introduction of a Pointwise Given Surface

The surface of an airfoil is not always analytically given and therefore the possibility of a pointwise given airfoil should be included. This can be done with so called cubic splines. The spline function here used is described in (16). It reads

$$S_{i}(x) \equiv a_{i} + b_{i}(x - x_{i}) + c_{i}(x - x_{i})^{2} + d_{i}(x - x_{i})^{3}$$

for $x \in [x_{i}, x_{i+1}]$, $i = 0$ (1) $n - 1$

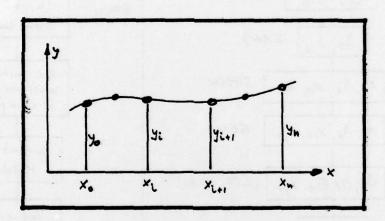


Fig. 17 The Spline Function

The demand for

$$S_{i} (x_{i}) = y_{i} \qquad i = 0 (1) n$$

$$S_{i} (x_{i}) = S_{i-1} (x_{i}) \qquad i = 1 (1) n$$

$$S'_{i} (x_{i}) = S'_{i-1} (x_{i})$$

$$S''_{i} (x_{i}) = S''_{i-1}$$

$$i = 1 (1) n-1$$

leads to a system of equations which give for each interval (x_i, x_{i+1}) the coefficients of the spline $S_i(x)$. With

The coefficients can be expressed for known c_i as

$$a_{i} = y_{i}$$

$$b_{i} = \frac{1}{h_{i}} (a_{i+1} - a_{i}) - \frac{h_{i}}{3} (c_{i+1} + 2 c_{i})$$

$$d_{i} = \frac{1}{3h_{i}} (c_{i+1} - c_{i})$$
(7)

The c_i are the solution of the linear algebraic system

$$h_{i-1} \cdot c_{i-1} + 2 c_{i} (h_{i-1} + h_{i}) + h_{i} c_{i+1} =$$

$$\frac{3}{h_{i}} (a_{i+1} - a_{i}) - \frac{3}{h_{i-1}} (a_{i} - a_{i-1}) \qquad i = (1) n-1$$
(8)

where c and c are presumed to be zero:

Once the coefficients a_i through d_i are known, we have a set of polynomials which are excellent for interpolation between x_i and x_{i+1} . As in each x_i not only the value but also the slope and the curve of the two neighboring polynomials are identical, $S_i(x)$ is also applicable to determine $S_i(x)$ and $S_i'(x)$ of a pointwise given function between those points:

$$S_{i}'(x) \equiv b_{i} + 2 c_{i} (x - x_{i}) + 3 d_{i} (x - x_{i})^{2}$$

 $S_{i}''(x) \equiv 2 c_{i} + 6 d_{i} (x - x_{i})$

In WING the subroutine PROFIL follows two options:

LO4 = 1 Airfoil Analytical Given

LO4 = 2 Pointwise Given

For LO4 = 2 system (8) is solved. With (7) the coefficients can be recreated at any time. Thus it is possible to read in the geometry of any reasonable airfoil and obtain good approximations for position, slope and curve at any station of it. These values are needed for the steady boundary conditions in B O U N D called from R A N D and S H O C K.

3.7 RESULTS

In this chapter WING - results are compared with (2), Teipel's linear and nonlinear airfoil results, using the method of characteristic and with Verdon's (18) analytical results for the velocity and pressure distribution over a flat plate.

Figs. 18 through 23 provide a comparison of the Teipel - with the W I N G results. Although it can be said that the linear solutions ($\tau = 0$) generally agree quite well, this is not true for higher Mach numbers, as can be seen in

Fig. 20 for M=1.4. For this case the real part of C_p shows considerable differences distributed toward the TE. For $\tau \neq 0$ there are differences disbuted over all cases. This was already noted by Chadwick (4) and Strada (5) who argued that the deviations could be caused by different difference equations, averaging procedures and number of grid points. However, this could not be proved because Teipel does not expose his complete set of finite difference equations.

This work does not show a comparison with Chadwick's results (4) obtained for airfoils $\tau \neq 0$ and wedges. Nevertheless this has been done. It shows excellent agreement in all cases no matter whether airfoil or wedge. Thus it can be said that W I N G is at least equivalent to Chadwick's solution over the airfoil.

Further W I N G - results are compared with Verdon's work (18), which provides analytically derived results over the flat plate and downstream of it into the wake.

First it should be noted that the C_p - distribution for $0 \le x \le 1$ is very good. This can be seen in Fig. 24 and 25. The disagreement for the plunge motion in Fig. 25 looks appreciable only because of the extended scale. The original Verdon plots in this area are rather hard to read because of his desire to show the wake pressure distributions rather than those over the airfoil for the zero upwash. One of his results is that the assumption of $v_1 = 0$ on the slip-line is wrong.

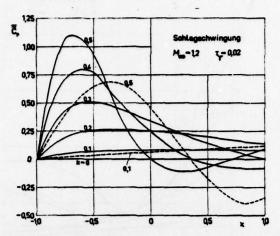
It has to be $C_p = 0$. With this assumption he gets a v_1 - distribution on the slip-line shown in Fig. 26 and 27.

The system of equations used in W I N G does not presume $cp_u = 0$ and $cp_L = 0$ but only $cp_u = cp_L$. It turns out that $|c_p|$ on the slip-line is of the order 0.000001. The upwash behind the airfoil obtained by W I N G is shown in Fig. 26 and 27 for $1 \le x \le 4.6$. For this range the agreement is

considered to be excellent for pitch and plunge. The reason for not showing results in the area x > 4.6 is storage difficulties in the time sharing system of the used IBM 360/67.

If we proceed from the slip-line into the wake field, Verdon results versus W I N G - results are shown in Fig. 28 and 29. Again the agreement is very good. The total unsteady pressure along characteristics with origin x = 1.08 and x = 2.0 from the LE is shown. In this way the pressure field over the wake can be represented. Fig. 28 and 29 are very important, as the single airfoil is only the first step to the cascade.

The good agreement with Verdon's results justifies the method of characteristic approach for the problem of the oscillating cascade. The advantage of this approach is the flexibility with which characteristic computational procedures can be designed.



 $\tau_{\rm T} = 2\tau$

Abb. 8 a. — Druckverteilung \tilde{C}_{ρ} (Realteil) für ein Parabelbogenprofil mit dem Dickenverhältnis $\tau = 0.02$ (Parameter : reduzierte Frequenz k).

Fig. 18A. From (2)

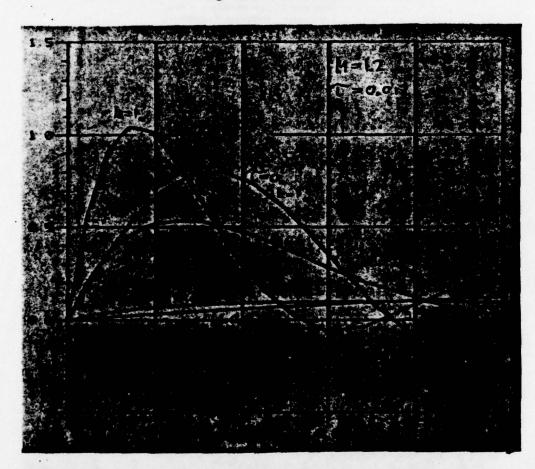


Fig. 18B. Pressure Distribution (Real Part) from W I N G , Plunge Motion

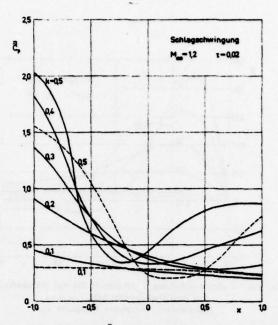


Abb. 8 b. — Druckverteilung \tilde{C}_p (Imaginärteil) für ein Parabelbogenprofil mit dem Dickenverhältnis $\tau = 0.02$ (Parameter : reduzierte Frequenz k).

Fig. 19A. From (2)

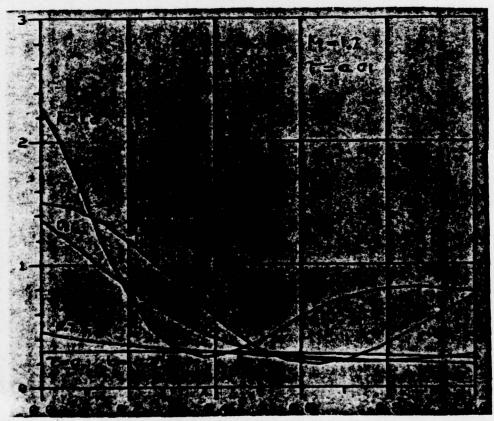


Fig. 19B. Pressure Distribution (Imaginary Part) from W I N G , Plunge Motion

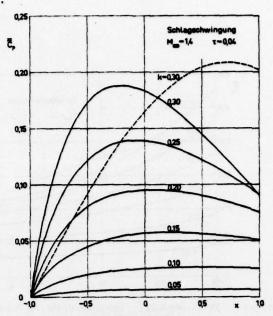


Abb. 9 a. — Druckverteilung \tilde{C}_p (Realteil) für ein Parabelbogenprofil mit dem Dickenverhältnis := 0,04 (Parameter : reduzierte Frequenz k).

Fig. 20A. From (2)

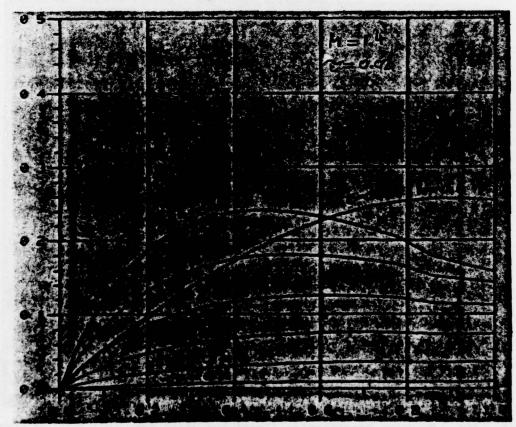


Fig. 20B. Pressure Distribution (Real Part) from W I N G , Plunge Motion

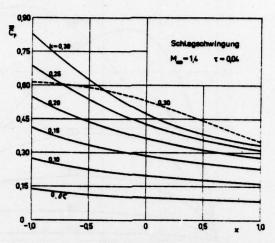


Abb. 9 b. -- Druckverteilung C, (Imaginärteil) für ein Parabelbogenprofil mit dem Dickenverhältnis \tau = 0,04

(Parameter : reduzierte Frequenz k).

Fig. 21A. From (2)

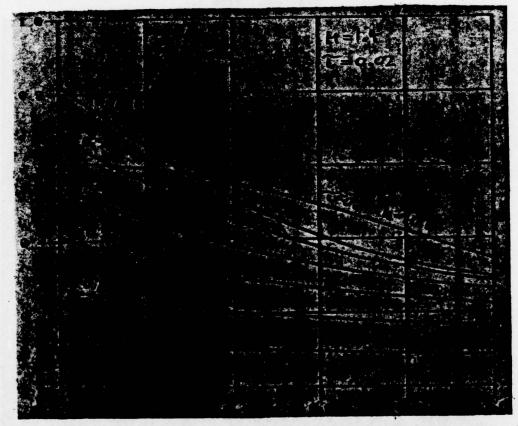


Fig. 21B. Pressure Distribution (Imaginary Part) from W I N G , Plunge Motion

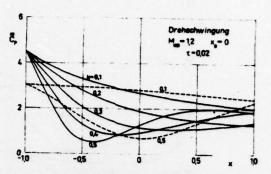


Abb. 13 a. — Druckverteilung Č, (Realteil) für ein Parabelbogenprofil mit dem Dickenverhältnis : = 0,02 (Parameter : reduzierte Frequenz k).

Fig. 22A. From (2)

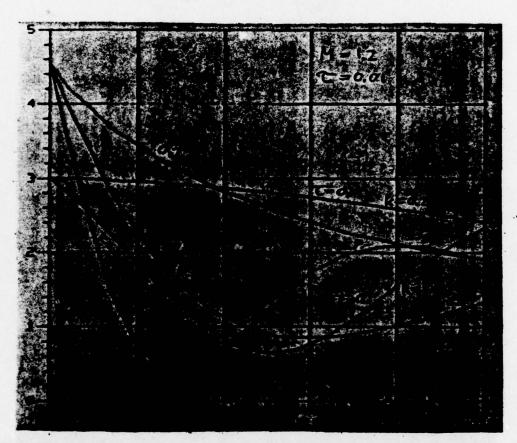
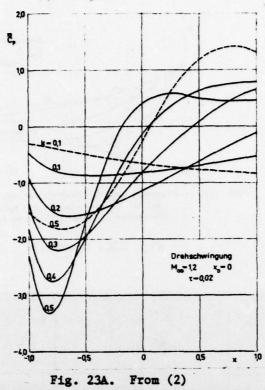


Fig. 22B. Pressure Distribution (Real Part) from W I N G , Pitching Motion



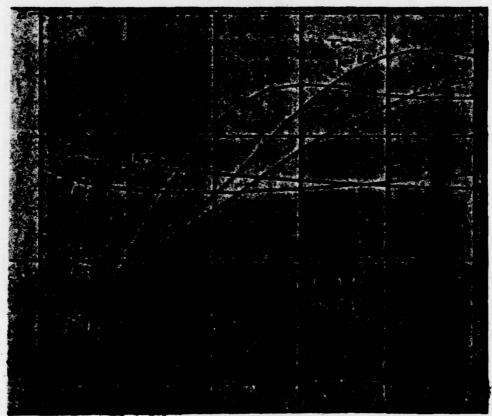


Fig. 23B. Pressure Distribution (Imaginary Part) from W I N G , Pitching Motion

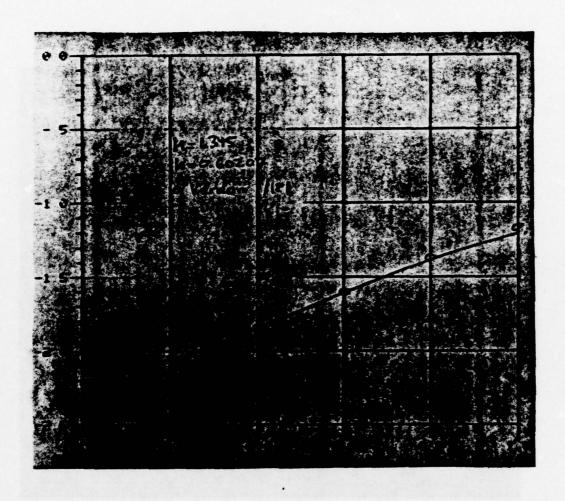


Fig. 24A. Pressure Distribution on Airfoil (Real Part) for Pitching Motion

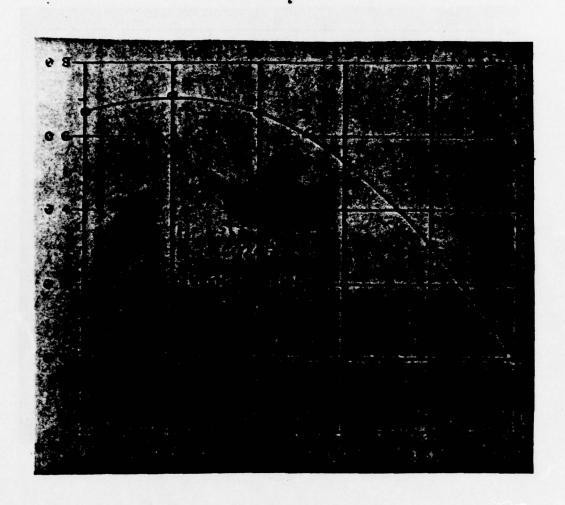


Fig. 24B. Pressure Distribution on Airfoil (Imaginary Part) for Pitching Motion

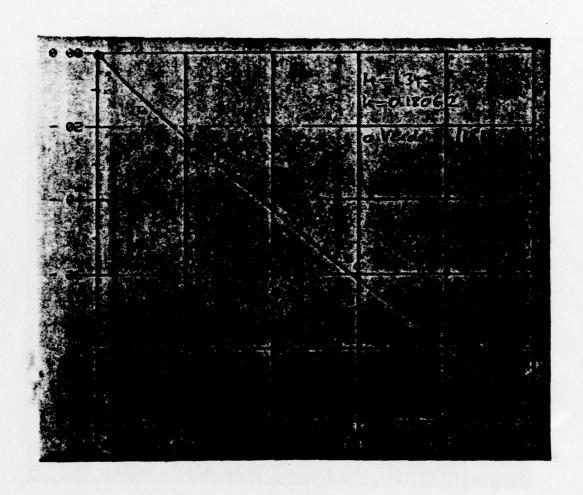


Fig. 25A. Pressure Distribution on Airfoil (Real Part) for Plunge Motion

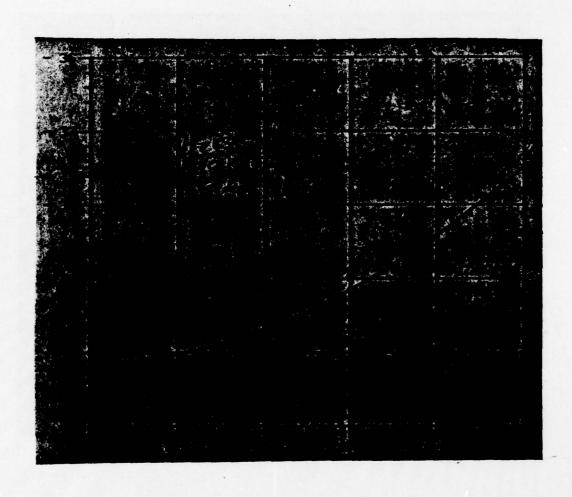


Fig. 25B. Pressure Distribution on Airfoil (Imaginary Part) for Plunge Motion

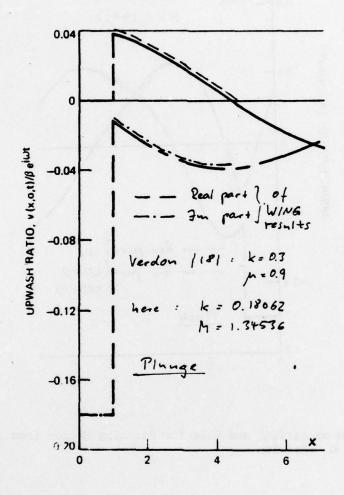


Fig. 26. Upwash on Airfoil and Wake for Plunge Motion from (18) Compared with W I N G -Results

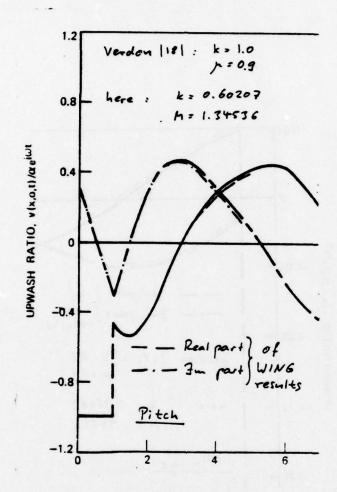


Fig. 27. Upwash on Airfoil and Wake for Pitching Motion from (18) Compared with W I N G -Results

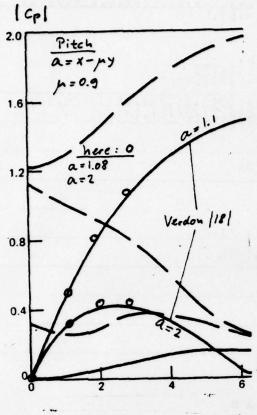


Fig. 28. Pressure Distribution Along Wake Characteristics for Pitching Motion from (18) Compared With W I N G - Results

DISTANCE ALONG CHARACTERISTIC, My

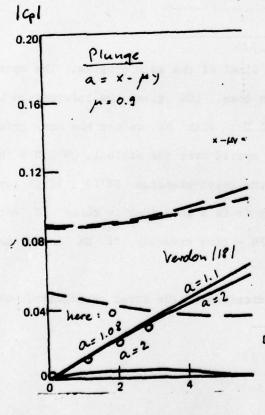


Fig. 29. Pressure Distribution Along Wake Characteristics for Plunge Motion From (18) Compared With W I N G -Results

DISTANCE ALONG CHARACTERISTIC, My

	3.8	Manual, Sample Data, Listing and Output of W I N G
		XXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
		SCILLATING SINGLE AIRFOIL FOR MACH.GT.1
_	c	CMFLFX#8 L, V, AI, PSI, G, DYDXU, PU
	C	GMMCN/EA/ ALC,ALC,AM,C1,AK,I,IE,IW
-	- 5	OMPCN/EC/ T.B.CYCX,CZYCXZ,OYOXU,AI,11,T3
		OMMEN/EC/ T.B.CYEX,C2YEX2,OYDXU,AI,11,T3 OMMEN/SP/ XS(2,50),YS(2,50),A(50,4),R(2,50) CMMEN/SCL/ EL(4.4).FI(4).ES(4)
	FG	GMMCN V(2,25,25),X(2,25,25),P(2,25),U(2,25,25),PSI(2,25,25), (2,25,25),AL(2,25,25),Y(2,25,25),Q(2,3),DX(2),IN(2,25),
	FP	X(2,25),PS(2,20),PU(2,20),TET(2),TEE(20)
č_		PTICNS:
		ET = C : COMPLETE OUTPUT =1 : CNLY PRESSURE DISTRIBUTION =2 : STOP
		C2 = C : PITCH =1 : PLUNGE
		C3 =C : NC WAKE =1 : WAKE SHALL BE COMPUTED C4 =1 : SURFACES ARE ANALYTICALLY GIVEN
		=2 : WEDGE (NO WAKE) =0 : SURFACES ARE POINTWISE GIVEN A = NUMBER OF REINTS ON THE SURFACE
C	k	V=25
10.	2 F	I=(MPLX(3.,1.) EAO(4,999) LO1,LO2,LO3,LO4,MA F(LC1.E6.2) GOTC 101
	I	F([C4.8C.2) LO3 = C Ax=MA ALL E3CEIL (b5.LC4.T1.T2.NX)
10	CR	EAD(4,1000) AK,AM,C,8,DX(1),DX(2) F(AM.EC.0.) GOTC 1C2

MAIN PROGRAM

GOTO 100

1C1 K=C
WRITE(1+1016) K
C ENC CF MAIN FREGRAM

Above is shown the beginning and the final of the main program. The options LO1 through LO3 determine what the program does. LO4 gives the information how the surface is given, needed in PROFIL. With MA we set how many gridpoints we want to distribute approximately equal spaced over the airfoil. WING has a built-in stepsize control, which changes the input-stepsize DX(I), if it turns out to be too large or too small. However it is a good idea to chose DX correctly because the iteration may increase the CPU - time severely, if DX is far away from the proper value.

For that purpose Fig. 30 gives an indication of the order of magnitude of DX for 14 points.

The next read - statements are in PROFIL:

```
SLEFCUTINE FRCFIL(WE,LC4,T1,T2,N)
CCMMCA/SF, XS/2,50),XS(2,50),A(50,4),P(2,50)

C FREPARATION OF THE PROFIL - SUFFACES

REAC(4,1000) T1,T2,NT,IKP
IF(LC4.E6.C) REAC(4,1001) SP
IF(L04.E6.2) REAC(4,1001) WE

DO E J=1.2

N=NT
T=T1
IF(J.E6.2) T-T2
IF(J.E6.2)
```

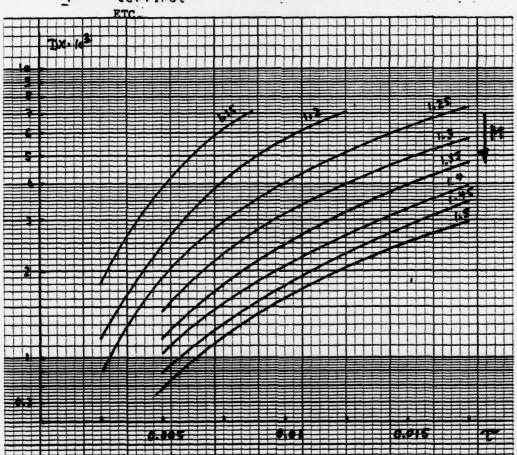


Fig. 30. Stepsize for N = 14 Points on the Airfoil

T1 and T2 represent the thickness of the airfoil. As W I N G is presented here, each surface is given by the parabola

$$y = 4\tau x (1 - x)$$

with

T = Tl for the upper

and

T = T2 for the lower side

For LO4 = 0 NT determines the number of points on the blade and IKP sets a flag to print out the input data.

IKP = 0 NO OUTPUT OF (X,Y)

- 1 OUTPUT

SP is a scaling factor to make the input data non-dimensional. If the chord is already 1, then SP = 1. Otherwise SP is the characteristic length, normally the chord. LO4 = 2 changes the airfoil into a symmetric wedge with the slope WE in degrees and unit chord, Fig. 31.

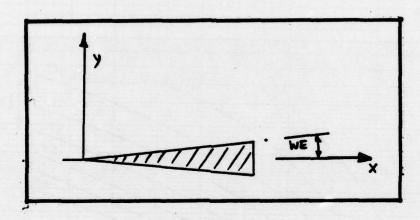


Fig. 31. Wedge Geometry

For LO4 = 0 the x,y-coordinates of the airfoil are read in accordance to the format statement 1001. The airfoil shape can be examined prior to entering W I N G in the test program T E S T as shown in chapter 5. T E S T is a preparing program for the airfoil.

If the airfoil is given by another function than Eq. (9), one has to change the statements in BOUND which compute y' = DYDX and y'' = D2YDX2.

After leaving PROFIL there is another read for

AK = k

AM = M

C = ae

B = x = b

DX(1) = STEPSIZE ON THE UPPER SIDE

DX(2) = STEPSIZE ON THE LOWER SIDE

If the run is finished, W I N G jumps back to this line and expects changed data for this airfoil. For M = 0 it jumps to the start of the main program for mode and/or airfoil changes. However for LO1 = 2 W I N G is finally terminated. This way a whole series of airfoil shapes and aerodynamic conditions can be examined in a single run. It should be noted that LO1 = 0 produces a complete field - output which is normally not necessary and rather long.

For LO3 = 1 the wake field output can not be suppressed as it is the only information which is printed about the wake. The whole wake may be skipped with LO3 = 0.

Besides the normal output file 06 W I N G has a second one: File 01.

Only the unsteady pressure distribution over the surfaces is written here. This file is used after the run by the plot program P L O T 1 to make plots of the pressure distribution on the Tektronix terminal in the computer center of the NPS. (S. Ch. 5).

W I N G contains still about 150 troubleshooting statements. All important subroutines start with

IKK = 0

The next statement is usually an if-statement converted into a comment card.

If this is activated, one can choose here the conditions for a complete tracing of this particular subprogram. Thus

If (IW.NE.O) IKK = 1

would trace this program during the whole wake - computation. One can imagine that a general IKK = 1 will result in irresponsible amounts of paper.

SAMPEL DATA SETS :

AIRFOIL ANALYTICALLY GIVEN:

FILE: FILE	FTO	4F001 P1		NAVAL	POSTGRADUATE
1101 16	0.01	20.0	mild to the second	Gerry Mids 36	
0.60	1.20	1.4	C.5	0.467E-02	0.4675-02
1101 16	6:55	1.4	0.5	0.467E-02	0.467E-02
1.00	1.20	20 C	0.5	0 • 467E-02	0.4676-02
2222 22	c.5	5.0	ö.5	0.0	3.0
OPTIONS:			- Only Pressu	re Distributi	
(Refer to WING)			- Plunge		(1)
			- No Wake	State DOT A "	(0)
			- Analytical - 16 Points		(1)
AIRFOIL DATA			T1 = T2 = 0	0.01	
		For	M = 1.2		
			- 1.4		
			b = 0.5		
		DX (1)	- Dx: (2) - 0.0	00467	

The program runs through k=1.0, 0.6 and 0.2. After M=0 we want with the same options the linear case TL=T2=0 for k=1.0 and 0.2. Finally we stop WING.

For the option-set 1001 the same data would be computed for the pitch mode. The sample output is the result for the first run of this data. Figs. 18b and 19b show the plots of the pressure distribution which we obtain with the data above. Figs. 22B and 23B show the same for the pitching mode.

AIRFOIL POINTWISE GIVEN :

2 17 1		343	
1 11.53	17.84	24-14	
o he manufacture of the second second	49.35	55.65	
5 74.55 95 0.08	8).88 0.0516	87-21	1000
0 100.	100.	100.	
5 44.18	=7.17	25.21 -2.49 56.83	
3 -2.0C 7 75.78	-2.97 82.06	-2.84	
9 -1.54	- C.C	-1.17	
5 1.4	0.5	0.5446-03	0.4546-02
	56 0.074 5 74.55 95 0.08 6 100 9 12.59 6 12.59 6 1.65 5 44.18 3 -3.60 7 75.78 9 -1.64 9 100	55 43.05 49.35 56 0.074 C.CE7 574.55 8).88 95 0.08 9.0516 0 1CC 1C0. 9 12.59 18.90 12.59 18.90 12.59 18.90 12.59 18.90 12.59 18.90 12.59 18.90 13.70 75.78 82.06 9 -1.54 -1.57 0 1CC 1CO	55 43.C5 49.35 55.65 56 0.074 C.CE7 0.097 574.55 8).88 87.21 95 C.OE 0.0516 0.016 C.C. 100 10C. 9 12.59 18.90 25.21 6 1.65 -2.12 56.83 3 -3.0C -2.97 -2.84 7 75.78 82.06 88.36 9 -1.54 -1.57 -1.17 0 1CC 100 100 100 100 100 100 100 100 10

OPTIONS:

IN PROFIL:

- ONLY PRESSURE DISTRIBUTION

- NO PITCH

- NO WAKE

- POINTWISE GIVEN SURFACES

- 16 POINTS ON EACH SIDE

- T1 = 0.01

NOT IMPORTANT

- T2 = 0.03

- N = 17 POINTS GIVEN FOR EACH SIDE

- IKP = 0 NO READ-BACK OF DATA

- SP = CHORD = 100

This example shows the different input-type for L04 = 0. The profil thickness T1 and T2 is not important any longer. But the variables have to be defined because they appear in the output. They can be used for identifying purposes.

Here x-location and y-value are given in percent of chord, therefore SP = 100.

b*C/U= 1.03C, N= 1.20, K= 1.40, E/C= 0.50, EX= 0.4675-02, T/C= 0.3109

FRESSURE-DISTRIBUTION UPPER SURFACE:

	PCINT	×	CPS	RC PU	ICPU	
	2, 1	0.0	0.143E 0C 0.128 00	-0.111E C1	-0.455E -0.402E	01 01
	6. 5	(.14C C.193	0.1135 dd C.968 = 01 C.6076 = 01	-0.1815 C1 -0.2095 C1 -0.2035 C1	-0.235E -0.155E	01 01
	i, 6 6, 7	C.251 C.312	0.443/-J1 0.475E-01 0.304E-01	-0.1235 01 -0.1235 01 -0.7255 00	-0.6215	33 00 33
	10, 5 11,10 12,11	C.448 C.522	0.1288-01	-0.332E CC	-0.604E	00 00
	13,12 14,13	0.685	-0.239E-11 -0.429E-31 -0.625E-01	0.7878-C1 -0.1125 CC	-0.116E -0.141E -C.158E	01 01
	15.14	0.869	-0.8262-01 -0.103E 00	-0.2978 00 -C.5015 00	-0.163F -0.157E	01 01
-	1/116	1.00€	-6-116E 0C	-0.5648 60	-0 ·155E	¥ŧ

w*C/U= 1.000, M= 1.20, K= 1.40, B/C= 0.50, EX= 0.4672-02, T/C= 0.010C

FRESSURE-CISTRIBLTION LOWER SURFACE:

	PCINT	X	CPS	RCPU	ICPU	
	1. 2	C.C C.043	0.143E 30 0.128E 30	0.7 2.1115 01	0.455E 01 0.492E 01	
	3, 4	0.090	0.113E 00 0.968E-01	0.181E C1 3.209E G1	0.326 E 01 0.235 E 01	
	5, 6 6, 7 7, 8	0.251	0.8072-01 0.643E-01 0.475E-01	0.2032 C1 0.1695 C1 0.1235 C1	0.1598 01 0.957E JJ 0.621E DD	
	8, 9	0.378	0.304E-01 0.128E-01	0.7255 00 2.332F 00	0.4925 00	
t	10.11 11.12 12.13	0.601 C.685	-0.531F-02 -0.239F-01 -0.429F-01	-0.1335 CO -0.7875-C1	0.116E 01 0.141E 01	
	13,14	C.774 C.869 C.969	-0.6256-01 -0.8266-01 -0.1036-00	0.1128 00 0.2978 00 0.5018 CC	0.158F 01 0.163F 01 C.157F 01	
	16,17	1.000	-0.110E 00	0.564E 00	0.155F 31	

MCMENTUM- AND LIFT - COEFFICIENTS FOR A SINGLE AIRFOIL WITH A SURFACE DESCRIBED ABOVE:

UNSTEADY -3.0585 1.3515 0.2039 -0.2964

4. The Oscillating Finite Cascade: CASCADE

4.1 The Physical Difference of the Two Problems

The name of the program, developed to calculate the inlet flow of a staggered finite cascade with thick airfoils is CASCADE.

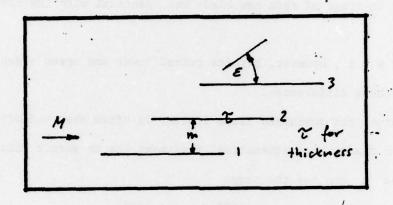


Fig. 32. Staggered Cascade of Airfoils

Nearly all basic steps to simulate this problem are already done in W I N G and therefore described in Chapter 3. However, there are two major differences which have to be considered in its mathematical treatment.

1. Only the first blade is exposed to the free stream flow without any disturbances. This means for the cascade of flat plates already that the perturbations of (n-1) blades hit the n-th blade and have influence on the development of the flow and the shocks at this airfoil. For the cascade of curved blades there is an additional consequence:

The constant value

$$C_{\infty} = (M_{\infty}^2 - 1)^{3/2}$$

from Eq. (2.15) can not be used any longer, as in general the steady flow will not have the free stream velocity.

Similarly, the wake slip line will not be parallel to the x-axis but it will rather have the direction of the field in front of that particular blade. This is the reason why the wake slope in W I N G was not simply set to zero. Because of these difficulties the first approach to the cascade including a thickness effect was made with airfoils whose upper surface were flat. This way the steady field in front of each new blade was identical with the free-stream field (see /5/).

C A S C A D E , however, permits curved lower and upper sides and is able to consider these differences.

2. As in actual turbomachines blade flutter is often observed with a phase lag from blade to blade, the mathematical treatment has to permit this. The phase lag is called μ and has the range

which covers all cases.

The introduction of the phase lag can be done relatively easily and is shown in Section 4.3.

4.2 The Constant Value Along the Characteristics

We consider Eq. (2.13):

$$\lambda^{3/2} = \mu = C_{\alpha,\beta}$$

(where in this case $\,\mu\,$ is not the phase lag) as long as the incoming characteristics originate in the freestream, $\,\mu\,$ is zero and therefore

$$c_{\alpha,\beta} = (M^2 - 1)^{3/2}$$

(Fig. 1)

This still holds in a cascade of flat plates, because here the steady field is identical with the freestream field. This is also true for a cascade with

blades which have flat upper surfaces.

It is not valid any longer when we permit curved surfaces on both sides, as the steady inlet flow field for the blade (n \geq 2) will in general be deflected. Therefore μ on the characteristics is not zero and has influence on $^{\text{C}}_{\alpha,\beta}$.

Fig. 33 shows the geometry of an incoming characteristic which is reflected into the cascade.

In A it has to be

$$C_{\infty} = \lambda_{\infty}^{3/2} = (\lambda_{A}^{3/2} + \mu_{A}) B$$
 (1)

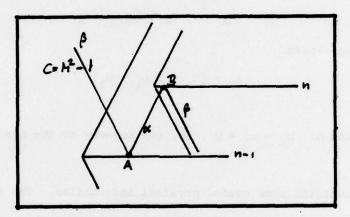


Fig. 33, Reflection of Characteristics

Thus we get λ_A by solving for it, as μ_A is known from the steady boundary conditions.

The next step provides the new constant along the a- characteristic.

$$C_{A\alpha} = \lambda_A^{3/2} - \mu_A \tag{2}$$

This is now the value which has to be applied in B in order to solve for λ_{B} .

$$\lambda_{A}^{3/2} - \mu_{A} = \lambda_{B}^{3/2} = \mu_{B}$$

$$\lambda_{\rm B} = \left[\lambda_{\rm A}^{3/2} - \mu_{\rm A} + \mu_{\rm B}^{2/3}\right]^{2/3} \tag{3}$$

with

$$\lambda_{\mathbf{A}}^{3/2} - \mu_{\mathbf{A}} = C_{\infty} - 2 \mu_{\mathbf{A}}$$

From Eq. (1) we obtain

$$\lambda_{\rm B} = \left[{\rm C_{\infty}} - 2 \, \mu_{\rm A} + \mu_{\rm B} \right]^{2/3} \tag{4}$$

In the case of $~\mu_{A}=\mu_{B}=0$, Eq. (4) reduces to the flat-plate solution $\lambda_{R}=\lambda_{\infty}$.

Eq. (4) contains some useful physical information. The maximum deflection δ_{max} which is connected to each Mach number M cannot be used for the leading-edge-slope of an airfoil in a cascade. Instead it has to be considerably smaller due to the double deflection in A and B. Only in case of a flat upper surface (5), the slope in B could be δ_{max} . The general case, μ_{A} and μ_{B} not zero, demands smaller deflections if detached leading edge shocks shall be avoided. This is a necessity according to our earlier assumption of weak shocks.

The calculation of the particular constant value, valid along the respective considered characteristic is done in C O N S T 1 .

4.3 The Phase Lag

Each time the unsteady properties \mathbf{u}_1 , \mathbf{v}_1 or Ψ are computed, the result is always the amplitude of the oscillating function

$$F(x,y,t) = A(x,y) \cdot e^{ikt}$$

Therefore, if we consider the two neighbouring blades (n-1) and n , where n leads the oscillation with the phase lag μ , we can express this for the time t as follows

$$F_{n} = A_{n} \cdot e^{ikt}$$

$$F_{n-1} = A_{n-1} \cdot e^{i(kt-\mu)}$$

For the single airfoil we did not need to write the exponent expression because it was always a common factor.

We reconsider this. For an example we take Eq. (2.40) which is the used unsteady shock polar. The first of those two formulas reads now, if written explicitly

$$\hat{u}_{1} \cdot e^{ikt} = (m_{1} G_{y} + i m_{2} G) \cdot e^{ikt} + (m_{3} u_{1} + m_{4} v_{1}) e^{i(kt-\mu)}$$

$$\hat{u}_{1} = m_{1} G_{y} + i m_{2} G + (m_{3} u_{1} + m_{4} v_{1}) e^{-i\mu}$$
(5)

We see that the phase lag does not cancel out. If we do this for all the equations used in W I N G, we find as a general rule: the term exp (ikt) always can be dropped. Each time when properties of the previous blade appear in the equation, they have to be multiplied by $\exp(-i\mu)$. Thus we reduce the magnitude of the (n-1)st -blade-values with the phase lag to the actual size they have

when they influence the nth -blade-properties. These are the unknown amplitudes of the oscillating functions connected to the nth -blade.

Eq. (5) gives an example how the equations change from W I N G to C A S C A D E . It is not necessary to show the whole set of finite difference formulas because they can be generated simply by using the concept outlined above.

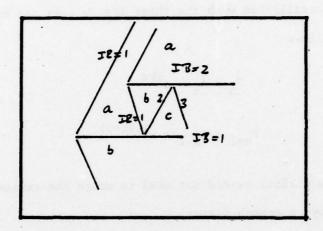


Fig. 34. Reflected Shocks in the Passage

The subroutines for the general and the boundary step can be taken as they were used in W I N G for the unsteady field. R A N D S is changed in so far, as the phase lag has to be added to the system of equations. The steady field is not affected.

For the field behind the reflected shocks in the passage one has to be careful, see Fig. 34.

It has not only to be considered that the oscillation of field b is ahead of that in field a , according to the phase lag μ . But also it is ahead of field c , which is again connected to the movement of blade l . This is done by reversing the sign of μ in RANDS for this case.

4.4 The Organization of the Program

As the computation of the reflected shocks and the fields behind them in the passages of a cascade is considered a rather complex procedure, the goal for the design of the program was to keep it as straightforward as possible. Therefore the convenient array organization from W I N G was adopted and extended. The main variable fields have the form

X(IB, IR, M, N)

The part (...,M,N) corresponds to Fig. 11 and 12. It allows the separation of fields on the upper and lower surfaces. IB indicates the blade, which is connected to the respective field and IR counts the shock in the passage, Fig. 34. This notation allows again repeated use of the same routines when only IB and IR are set correctly. Again throughout the whole program upper and lower sides are indicated by I=1 and I=2. The index I is used to control the M,N-notation and the sign like it was done in W I N G.

In CASCADE there appear three different kinds of flow fields which are shown in Fig. 34.

Type a is the field over the whole upper surface. The procedure is here the same as in WING. Computation of the b-field follows also WING, but it is stopped when the shock crosses blade 1. At this point the shock is reflected and the main difference in the subroutines is a transformed system of coordinates, because the origins of the shock and the system of coordinates have to be identical. This transformation causes considerable changes of the statements in SHOCK and FIND, compared to their counterparts in WING. However, the computational sequence is not really changed, but the extension of those subroutines enables them to identify the type of the field and to make all necessary changes for signs, coordinates and termination points.

CASCADE starts with blade 2 (IB=2) exposed to an initial field where all perturbations are zero. After evaluating the complete upper field, the b-field is computed. This is terminated when the shock hits blade 1, which is assumed to be a fixed flat plate for this first step. So here blade 1 could be considered a wind tunnel wall. When all the reflections in this first channel are done, the IB=2-field is copied to the IB=1-field. Now the former can be used again. From this point on blade 1 is an oscillating airfoil with the given shape.

CASCADE has three output files:

- File 07 for complete field output and/or the final pressure distribution
- File 06 is a documentation which allows to follow the iterations and the main steps from blade to blade. Like WING, CASCADE contains trouble shooting statements which can be activated by IKK=1 for the desired subprogram. For a general IKK=0 File 06 will be limited to a few pages. IKK=1 will print a complete tracing of the particular routine, which increases File 06 significantly.
- File 01 contains geometric data of the cascade and shock configuration.

 It is used as an input file for P L O T 2 which produces plots like

 Fig. 35 on the Tektronix terminal. P L O T 2 is described in Section 5.3.

 The input data have to be in File 03. They are shown in Section 4.6.

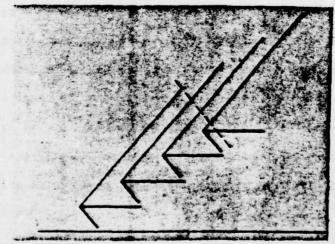


Fig. 35. Cascade A and B Shock Geometry [15]

Finally, it should be said that the organization of the storage arrays is actually a waste of memory space because fields of the b - and c - types (Fig. 34) need only small areas whereas a -fields demand large ones. This expensive way was chosen in order to make a general program clear. If one considers that the next step is the addition of the wake fields for each blade and that these have to be much larger than those in the passages, this solution gives an approach which can be extended analogous to W I N G. The desire to save storage place would certainly complicate the main organization and increase the number of subroutines significantly.

The main disadvantage, besides a slow-down of the computer is the limitation of the fields. The dashed line in Fig. 35 was not produced by P L O T 2 but it was added to show the limit of C A S C A D E . Beyond the end of the headshock for the first blade its influence has vanished. This results in completely wrong fields behind the dashed line. Hence the fields in Fig. 35 are usable only up to blade 3. The arrays in C A S C A D E used in this work have the size

X(2,3,50,20)

They can be extended to

X(2,3,KV,20)

where KV has to correspond with the first FORTRAN-STATEMENT of CASCADE.

Then all the counters and Do-loops are dimensioned properly. The number of possible blade-computation is thus given by the limitation of the computer.

4.5 Results

As a test for CASCADE three different supersonic cascades are checked for which results already exist. Cascade A and B are those introduced by Verdon (15;17) and cascade T is a model for an experimental cascade used by Fleeter (19) for which Strada (5) gives computer results for the inlet flow.

Cascade Data:

	A	В	T
M	1.345	1.281	1.550
m	0.4	0.301	0.335
τ _{up}	0.0	0.0	0.0
τ _{Lo}	0.0	0.0	0.03
ε	30.5	26.6	23.9
k	0.90226	0.75054	0.28

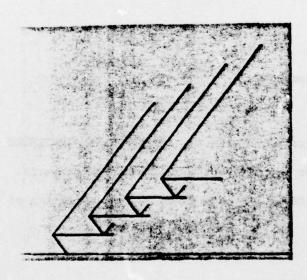


Fig. 36. Cascade B (15)

Fig. 35 shows the geometry of cascade A. Cascade B and T are shown in Fig. 36 and 37. As CASCADE is not completely programmed, it is only possible to show results of the inlet flow. We compare with theoretical aerodynamic data obtained by Verdon (15) for the linear cases, which are identical with those computed by Bell (9) and with results for the cascade T, given by Strada in (5).

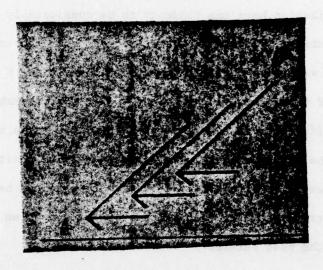


Fig. 37. Cascade T (5)

Fig. 38 and 39 show pressure difference distributions calculated by C A S C A D E for the third blade. Considering that Verdon introduced an infinite cascade and that Bell gave results for the 14th blade, the inlet flow computed with C A S C A D E looks very promising. The differences are not too significant and can be explained by the low number of blades which were examined.

As CASCADE and Verdon-results agree sufficiently, it could be expected that the linear data given by Strada and Platzer (P-) results would coincide very well with CASCADE computations. Actually, as both programs examine the second blade, the distribution of the total pressure coefficient agrees exactly for both sides of the A blade.

This is not true for the included slope effect. Fig. 40 through 43 show computations of this study in direct comparison with Strada (5). The upper surface represents the linear case and shall not be discussed because of good agreement. However, the lower surface is shaped and the agreement is good only at the leading edge. In all cases results of CASCADE for the total pressure distribution have the tendency to be considerably smaller than those presented by Strada. There is no real progress compared to the Fleeter Experiments (19) but with some imagination one could say that CASCADE meets the tendency of those data slightly better. It is remarkable that in spite of significant differences in the magnitude, the phase angle between imaginary and real part of the pressure coefficients agrees very well with Strada. The difference between the Strada results and those obtained here may be caused by the different treatment of the shock and remains a problem to be investigated more thoroughly in the near future.

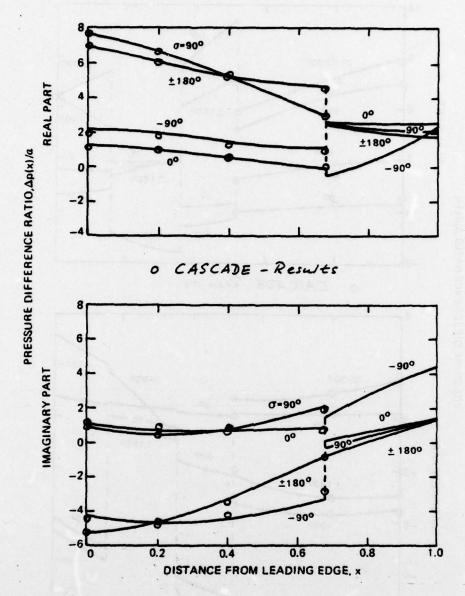


Fig. 38. From (15) Comparison of Linear Results Casc. A

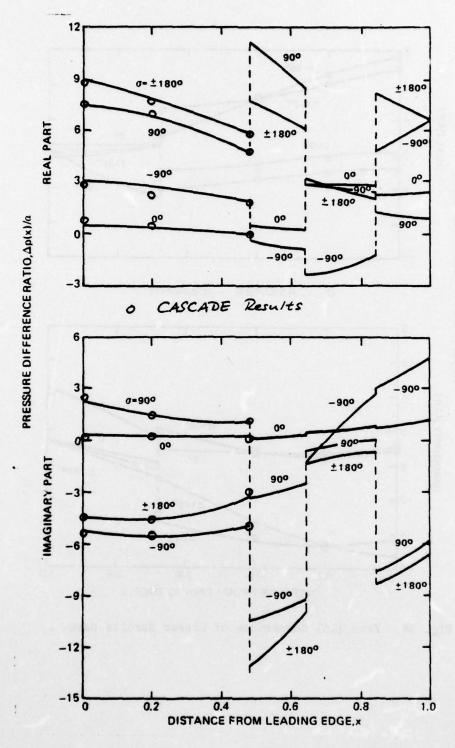


Fig. 39. From (15) Comparison of Linear Results Casc. B

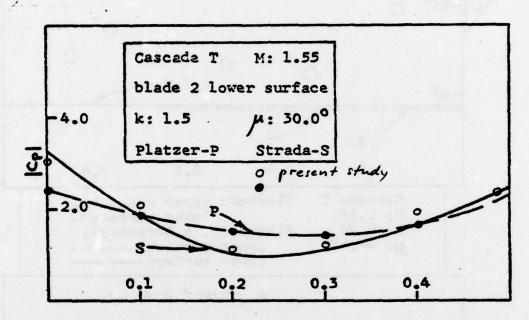


Figure 3.10.2 High-k comparison of the linear and non-linear solutions.

Fig. 40. From (15) Comparison with Strada

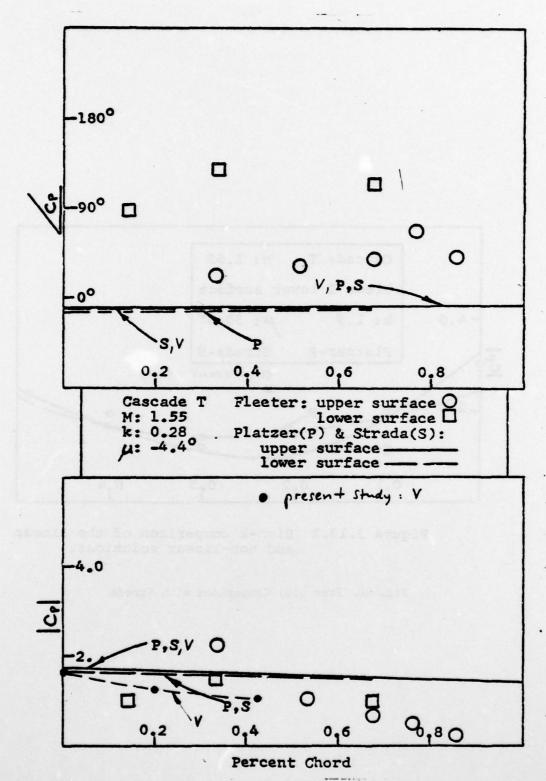


Fig. 41. From (5) Comparison with Strada

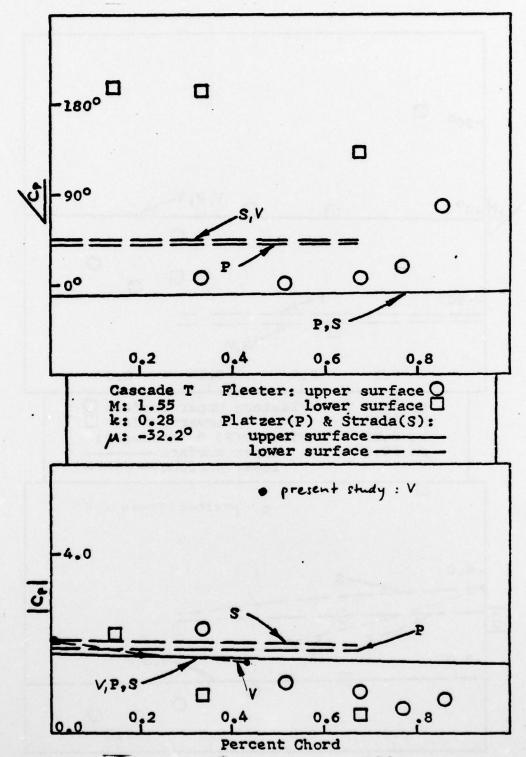


Fig. 42. From (5) Comparison with Strada

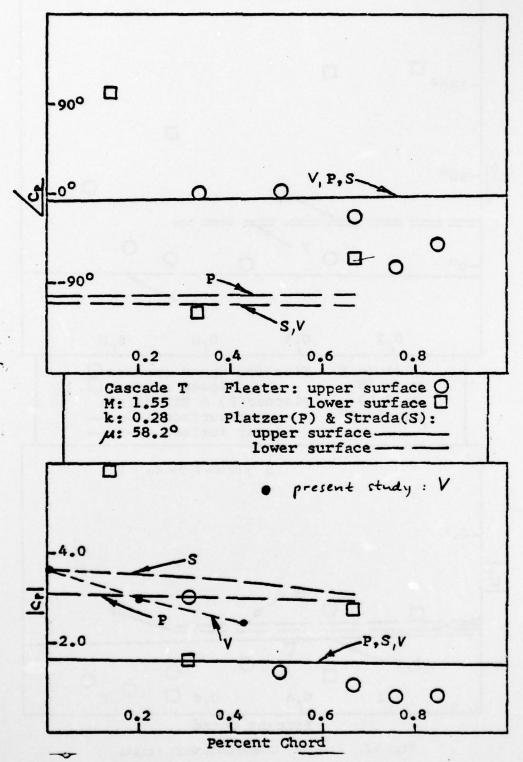
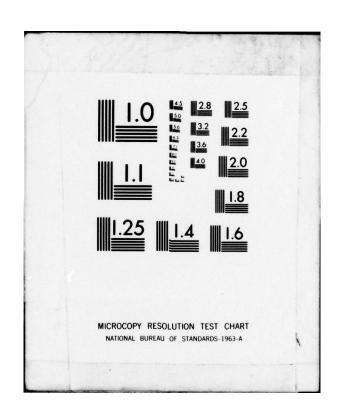


Fig. 43. From (5) Comparison with Strada

NAVAL POSTGRADUATE SCHOOL MONTEREY CA A METHOD OF CHARACTERISTICS APPROACH TO THE PROBLEM OF SUPERSON--ETC(U) OCT 78 K VOGELER NPS67-78-912 UNCLASSIFIED NL 2 OF 2 AD A076629 #15 E.E.

END DATE FILMED 12-79

AD-A076 629



4.6 Updated Listing, Data and Output of CASCADE

If one refers to the comment cards at the beginning of CASCADE, it is obvious that this program is a derivation of WING.

Č	CASCADE OF OSCILLATING AIRFOILS FOR MACH-GT-1
·	CCMFLEX*& U,V,PSI.G,CYDXU CCMPLEX*& AI.PU.EI
	CCMFLFX*16 EL,RI,ES REAL*4 TX,TY,XX,YY,XL,D1,D2,D3,XP,ALC,P,C3,ALZ COMMCN,EAL ALO,AM,ALC,C1,AK,I,IN,TX,IX,IL,MA,I,C,CT,XX,IX,IX,IX,IX,IX,IX,IX,IX,IX,IX,IX,IX,
	CCMMCN/EC/ TI, T2, T, B, CYCX, C2YCX2, DYDXU, AI, EI, II, IBACK, I E COMMCN/SP/ XS(2,5C), YS(2,5C), A(2,50,4), R(2,50)
	COMMCN/SCL/ 3L14,41,81(41,23(4) COMMCN V(2,3,50,2C),X(2,3,50,2C),P(2,50),U(2,3,5),2C), FPS1(2,3,50,2)),G(2,3,50,2D),AL(2,3,5C,2C),Y(2,3,50,2D), FG(2,50),CX(2),IN(2,5C),IN2(2,5C),PX(2,30),PS(2,30),FL(2,30)
<u>ç</u>	CATIONS:
Č	LC1 =C : COMPLETE OUTPUT =1 : CNLY PRESSURE DISTRIBUTION
C	=2 : \$TOP LC2 =C : PITCH =1 : PLUNG:
Ç	LC3 =1 : DC THE 1. FASSAGE NOT
COUCOCO	#0 : SURFACES AFE POINTWISE GIVEN MA = MAXIMUM NUMEER OF POINTS ON THE SURFACE .LT .20 MAXS = BLADE, WHERE THE PRESSURE DISTRIBUTION IS COMPUTED IBLA = FIRST BLADE OF FIELD - OUTPUT
7	KV=5(AI=CMSL X(C .1)
102	REAC(3,955) LOI, LC2, LC3, LO4, MA, MAXS, IBLA IF(LC1. EG. 2) GOTC 101
	MAX=MA MAX=MA CALL FRCFIL(LC4.T1.T2.NX)
100	REAC(3,1000) AK,AM,C.OX(1),CX(2) IF (AM.EQ.O.) GDTO 102

The data are set up in the same way. There are five additional input variables. Maxs and IBLA are explained above. IBLA is only meaningful for LO1=0. Then for example IBLA=3 starts the field output only from the third blade.

The other three variables define the Cascade:

ET = ε Stagger Angle

EM = m Distance Between the Blades

AMY = \mu Phase Lag

See Fig. 32.

The rest of the input data have the same definition as those in $\mbox{W I N G}$. Two sample data sets shall be shown:

1. Analytically Given Blades

FILE: FILE	FTC3F	001 P1		MAVAL P
1001160301	2 12	2		
1.75054 50.0	1.281	1.4	0.3925-02	
22 22 22 22 22	3.000	1.4	0.3925-02	0.3923-02

This would cause a run as follows.

lst Row: No field output, pitch motion, no computation between the first blade and the imaginary wall (first channel), 16 grid points on each surface, they are analytically given, the third blade shall be examined.

2nd Row: The blades have flat upper and thick lower sides

3rd Row: k, M, x , Dx

4th Row: μ , b, ϵ , m

5th and 6th Row: STOP

2. Pointwise Given Blades

FILE: FILE	FTC3	F001 P1		NAVAL
1000160201	- (, (3	17 1		
100.01	5.21	11.53	17.84	24.14
30.45	-C. 185 36.75 C. 056	43.05 43.05 -0.014	-0.050 49.35 0.087	-0.005 55.65
61.95	68.25	74.55 C.CE	80.88	67.21 0.016
-0.035 C.C	-0.10 -0.29	100.	10).	100. 25.21
31.53	37.35	-1.65 44.19	-2.12 50.51	-2.45 56.83
63.15 -2.60	65.47	75.7E.	62.76	88.36 -1.17
-0.70 0.28	-0.10 1.55	100.	100.	1 CO. 0 • 454 5 - C2
7.28	1.55	1.4	1.544t-J3	0.454=- 12
0.28	1.55 C.5	23.5	0.5446-03	0.4548-C2
222222222	(.00)	1.4	3.544103	0.454:-02

lst Row: No field output, pitch motion, no first passage, pointwise given blades, 16 gridpoints on the surface, 2nd blade shall be examined;

2nd Row: Identification for upper and lower side (no real importance here),
17 points given as input for each side, the geometry of the blades
shall be printed;

3rd Row: Chord is 100

46h to 19th Row: Input geometry of blades, according to subprogram PROFIL;

20th Row: k, m, , Dx

21st Row: μ, b, m, ε

Two other cases and STOP

The output of the pressure distribution for the first case is shown next, followed by a listing of C A S C A D E in Appendix B.

MOCE PHASE = -4.43 / 2. ELACE W+C/U=0.28000, M=1.550, K=1.40, B/C=0.50, T/C=0.0100FRESSURE-CCEFFICIENTS LPPER SURFACE: PCINT CPS RCPU ICPU C.C -C-169E C1 0. 3645 00 0.167 0.251 0.415 0.503 0.145E-01 0.995E-02 0.995E-02 0.418E-02 C.325E-02 0.870E-03 -0.282E-03 0.334 C.317 23 4, -0.1585 -0.1575 -0.1555 -0.1555 -0.1555 7.301 0.2525 0.2659 0.2480 0.6850 7. 0.568 9, 200 C.752 C.814 30 0.6461-02 0.8511-02 0.650 14,13 15,14 16,15 0.876 0.538 1.000 01 *********** h*C/U= J.28000. M= 1.55C. K= 1.40. B/C= 0.50. T/C= J.0300 PRESSURE-CCEFFICIENTS LOWER SURFACE: PCINT CFS RCPU ICPU 0.4686 0.3458 0.3348 0.3358 0.3178 0.2178 0.2208 0.235E (.1800 i, 4, 6. 9,10 0.242 0.1929 0.341 C.398 -0.9586 -0.9878 -0.9805 0.1578 0.442F-C.192E--)î -)î -)î 13,14 0.462 C.451 0.168E ************************ 90

5. The Programs as a Working System

5.1 TEST

TEST was written to examine the blade geometry if it is pointwise given.

CASCADE and WING are not prepared to handle shocks which are caused by a changing slope over the airfoil. Therefore it has to be checked that the surface has no turning point.

For a given set of points one has to deform the PROFIL in such a way that the curvature of it never changes the sign. This would be indicated by a changing sign for the second derivative of the surface function.

An example for TEST - input data is shown below:

100.	3.03	17 7		
-2.26	-5.185	11.53	17.84	24.14
30.45	36.75	43.05 6.674 74.55	49.35 0.087 80.38	2.07 2.797 87.21
9.059	100.00	₹.08 5.0	1. 1516	0.116
-0.26	3.29 -1.08	12.59	18.90	25.21
31.53 -2.76 63.15	37.35 -2.53 65.47	44.18 -3.0)	20.51	56.33 -2.34
-2.60 54.66	100.00	-1.94 C.C	82.06 -1.37 0.0	-1:17 0.0
-0:13	-0.10	10).	130:	100.

TEST reads from File 05 and writes the result x, y_L^* , y_u^* , y_L^* and y_u^* on File 06 and 01. File 06 can be printed and shows the original input data, the interpolated surfaces and besides the derivatives mentioned above, the coefficients of the cubic splines for each side. File 06 contains only x, y, y' and y". It is used as an input file for P L O T 1 which produces plots like Fig. 44 and Fig. 45 with it. Fig. 44 shows y' of the two PROFIL sides and Fig. 45 shows y'. P L O T 1 is explained in Section 5.2.

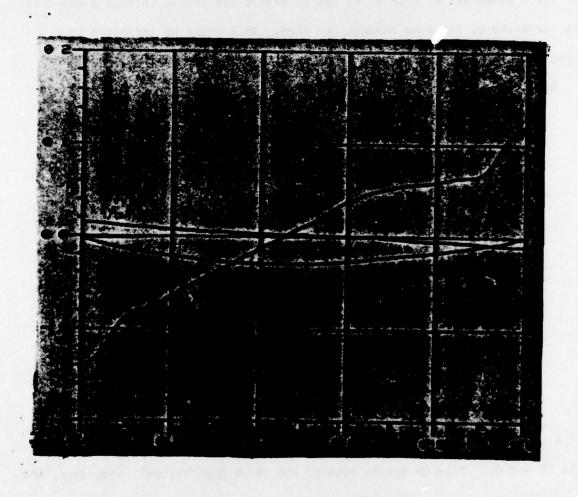


Fig. 44. Blade Geometry and y'

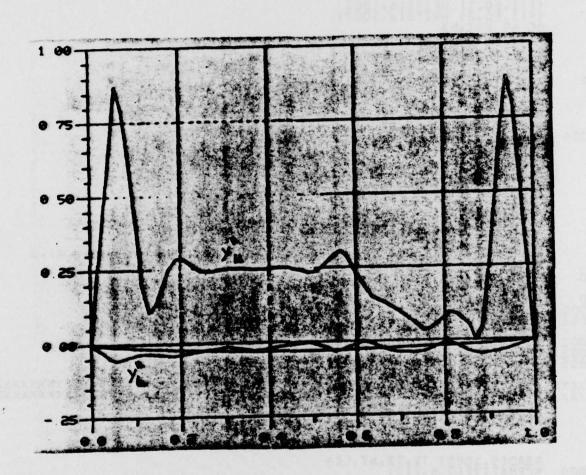


Fig. 45. Blade Geometry and y"

```
TEST LISTING :
ooo ooooo ooo
                        TEST
COMMON/SP/ XS(2,50),YS(2,50),A(50,4),R(2,50)
                                                           HE SPLINEFUNCTIONS FOR THE PROFILSURFACES ANALYTICALLY GIVEN POINTWISE GIVEN
                        CALL PRCFIL (LO4, T1, T2, N)
                        CCMPUTATION OF THE SPLINE - COEFFICIENTS
                      CCMPOTATION OF THE SPLINE - COEFFICIENTS

CC 30 I=1,2
IF(I.EQ.2) WRITE(6,1002)
IF(I.EQ.2) WRITE(6,1003)
OC 49 K=1,NX
A(K,1)=YS(I,K)
A(K,3)= R(I,K)
H=XS(I,K+1)-XS(I,K)
A(K,2)=(YS(I,K+1)-YS(I,K))/H-H*(R(I,K+1)+2.*R(I,K))/3.
A(K,4)=(R(I,K+1)-R(I,K))/(3.*H)
WRITE(6,1001)
DX=1./65.
MM=70
WRITE(1,1004) MM
FCRMAT(I3)
DC 10 M=1,70
X=DX*(M-1)
DC 5 K=2,50
J=K-1
IF(XS(I,K).GE.X) GDTD 6
CCNTINUE
H=X-XS(I,J)
DYDX=A(J,2)+2.*A(J,3)*H+3.*A(J,4)*F*F
D2YDX=2.*A(J,3)+6.*A(J,4)*F*F
D2YDX=2.*A(J,3)+6.*A(J,4)*F*F
49
1004
56
                       H=X-XS(I,J)
DYDX=A(J,2)+2.*A(J,3)*H+3.*A(J,4)*H+*H
D2YDX2=2.*A(J,3)+6.*A(J,4)*H
Y=A(J,1)+A(J,2)*H+A(J,3)*H*H+A(J,4)*H*H+
WRITE(6,1000) X,Y,DYDX,C2YDX2,A(J,1),A(J,2),A(J,3),A(J,4),J
WRITE(1,1005) X,DYDX,C2YCX2
CCNTINUE
MM=O
WRITE(1,1004) MM
FCRMAT(1x,8(2x,E12.5),I4)
FCRMAT(1x,7x,'x',13x,'y',9x,'DYDX',Ex,'D2YDX2',10x,
FCRMAT(1x,7x,'x',13x,'y',9x,'DYDX',Ex,'D2YDX2',10x,
FCRMAT(1H1,1x,'PROFIL LPPER SURFACE: ',//)
FCRMAT(1H1,1x,'PROFIL LCWER SURFACE: ',//)
10
CCC
                         PREPARATION OF THE PROFIL - SURFACES
                        READ(5,1000) T1,T2,NT,LC4
FCRMAT(2F10.5,212)
IF(L04.EC.0) READ(5,1001) SP
DC 8 J=1,2
I = 0
1000
                       IZ=0

N=NT

T=T1

IF(J.EG.2) T=T2

IF(L04.NE.0) GOTO 12

DC 9 K=1.4

M=(K-1)*5+1

READ(5,1001) XS(J,M),XS(J,M+1)

READ(5,1001) YS(J,M),XS(J,M+1)

WRITE(6,2001)XS(J,M),XS(J,M+1)

WRITE(6,2001)XS(J,M),XS(J,M+1)

WRITE(6,2001)YS(J,M),YS(J,M+1)

FCRMAT(1X,5F10.5)

FCRMAT(5F10.5)
2001
```

```
IF(YS(J,M+4).EQ.100.) GCTO 7

CCNTINUE

CCNTINUE

DC 2 LL=1,N

XS(J,LL)=XS(J,LL)/SP

YS(J,LL)=YS(J,LL)/SP

WRITE(1,1004) NT

FCRMAT(13)

DC 16 KV=1,NT

WRITE(1,1005) XS(J,KV),YS(J,KV),YS(J,KV)

CCNTINUE

IF(L04.EQ.0) GOTO 3

DX=1.0/(N-1)

I1=(-1)**(J+1)

DC 20 K=1,N

XS(J,K)=(K-1)*DX-0.25

YS(J,K)=4.*T*XS(J,K)*(1.-XS(J,K))*I1
2
1004
16
 12
NCOCOMOCO
O
                     TC ENTER THIS PART, THE SURFACES SHOULD ALREACY BE GIVEN POINTWISE
                     CONTINUE
                     INTERPOLATION THROUGH CUBIC SPLINES
                    IF(T.NE.O. .OR

DC 52 I=1,N

A(I,1)=0.

GOTO 53

CONTINUE

DC 10 I=1,N

A(I,3)=XS(J,I)

A(I,4)=YS(J,I)
                                                          .OR. LO4.EC.C) GOTO 51
52
 51
10
C
C
C
                     MATRIX OF COEFFICIENTS AND RIGHT SICES
                 K=N-2

DC 25 I=1,K

A(I,1)=A(I+1,3)-A(I,3)

A(I,3)=A(I+2,3)-A(I+1,3)

A(I,2)=2.*(A(I,1)+A(I,3))

A(I,4)=3.*(A(I+2,4)-A(I+1,4))/A(I,3)-

F3.*(A(I+1,4)-A(I,4))/A(I,1)

CCNTI NUE

A(1,1)=0.0

A(N-2,3)=0.0
 25
                     THE STEP OF GAUSS
                    CC 30 I=1,K

DO 35 M=3,4

A(I,M)=A(I,M)*A(I+1,1)/A(I,2)

A(I,1)=0.0

A(I,2)=A(I+1,1)

A(I+1,2)=A(I+1,2)-A(I,3)

A(I+1,4)=A(I+1,4)-A(I,4)

CONT INUE
 35
                      SCLUTION
                    A(1,1)=0.

A(N,1)=0.

A(N,1)=0.

A(N,2,1)=0.0

L=N-1

DC 40 I=2,L

K=N-I

M=K+1

A(M,1)=(A(K,4)-A(K,3)*A(M+1,1))/A(K,2)

CCNTINUE

DC 11 M=1,N

R(J,M)=A(M,1)

IF(IZ.EQ.1) GOTO 8

KV=N-1
 11
```

```
DC 49 K=1,KV

A(K,1)=YS(J,K)

A(K,2)=(J,K+1)-XS(J,K)

A(K,2)=(YS(J,K+1)-YS(J,K))/H-H*(R(J,K+1)+2.*R(J,K))/3.

A(K,2)=(R(J,K+1)-R(J,K))/(3.*H)

DX=1./49.

DC 4 M=1,50

X=DX*(M-1)

DC 5 K=2,50

I=K-1

IF(XS(J,K).GE.X) GOTO 6

CONTINUE

H=X-XS(J,I)

Y=A(I,1)+A(I,2)*H+A(I,3)*H*H+A(I,4)*F*F*H

R(J,M)=X

YS(J,M)=X

YS(J,M)=X

YS(J,M)=X

YS(J,M)=X

CCNTINUE

IZ=1

N=50

DO 13 I=1,50

XS(J,I)=R(J,I)

GCTO 3

CCNTINUE

WRITE(6,1003) N

DG 15 M=1,N

WRITE(6,1002) XS(1,M),YS(1,M),XS(2,M),YS(2,M)

N=N-1

1002 FCRMAT(5X,4(E12.5,2X))

RETURN

FORMAT(1HI,5X,*THE COMPLETE EXTRAPOLATEC SURFACES N=*,I3,//)

RETURN

FORMAT(1SF10.5)

RETURN
```

5.2 PLOT 1

P L O T 1 was written to produce plots from two pointwise given functions

$$Y1 = F1(x)$$

$$Y2 = F2(x)$$

The x-stations are for both functions identical. This is useful to make diagrams for real and imaginary part of pressure distributions (see Fig. 18 through 25) and to visualize the y' and y" values of the blade geometry (Fig. 44 and 45). The input is read from File 01. This was produced from the respective program whose results shall be plotted. That is either W I N G or T E S T in this work.

The next step is to look in File 01 for the highest and lowest values of the functions and to decide the scale of the diagram. The last line of File 01 is a zero. Behind this zero there has to be inserted now the coordinates of Points 1 to 4 from Fig. 46.

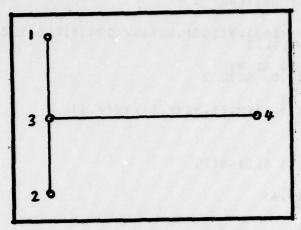


Fig. 46. Dimensioning of the Plot

This must be done for both expected diagrams. An example for the final lines of File Ol is shown below.

This file produced Fig. 44 and 45.

x	yl	у2
<u>i</u>		
0.89855	C. 16626	0.21406
0.913)4	C.C7518 C.C5115	C.69786 C.89099
0.97101	0.1344 0.11344 0.11613	0.83529 0.54123 0.25650
01.00100	0.12099	0.00000
9.	-5:13	
1: 3:	Ç.	
0. 1.	ć. J.	

One can see the inserted eight additional lines behind the zero.

P L O T 1 works only on the Tektronix-terminal. After unpacking and compiling, it can be called on this device with

\$\$ PLOT 1

After the first plot is done, the terminal makes a tone. Then PLOT1 hits a pause-statement. By striking an arbitrary key on the keyboard and CTRLS, it produces the second plot.

```
C
         PLOTI
                            LISTING
             MENSICN X(100), Y(100), X1(10,100), Y1(10,100), Y2(10,100), 4), B(4), N1(11)
             AD(1,1000) N1(M)
(N1(M).EQ.O) GOTO 2
=N1(M)
10 I=1,N2
AD(1,1001) X1(M,I),Y1(M,I),Y2(M,I)
1
10
2
                   [=1,4
,1001) A(I),B(I)
14
3
                      001) A(I),B(I)
15
5
         FORMAT(13)
FORMAT(3F10-5)
STOP
END
 _5.3 PLOT2
```

P L O T 2 is a special program which graphs the shock and cascade geometry on the Tektronix terminal. It is called with

SS PLOT 2_

Input is read from File Ol which is prepared by CASCADE in each run.

The resulting plots are for example Fig. 35, 36 and 37. PLOT 2 asks

for the size of the diagram by SCALE:

By enterine 1.5, 1., 0.5 or any other digital number from the keyboard, one may change the graph to the desired size. Each time the program has to be started again with the \$\$ - command.

```
C PLOT 2 READS THE DATA FROM FILE FT01F001
AND PROCUCES ON THE TEXTRONIX-TERMINAL PLOTS
OF THE GEOMETRY OF THE HEADSHOCKS AND THE SHOCKS

IN THE FASSAGES

DIMENSION X (50), Y (50)
WRITE (6,999)
READ (5,999) ST
CALL NPTS (2)
X(1)=20.
Y(2)=100.
Y(2)=20.
CALL CPLOT (X,Y)

CONTINUE
READ (1,1000) X(1), Y(1), IM
IF (IM.EQ.-1) GOTO 30
IK=IM
X(2)=X(1)+1.
Y(2)=Y(1)
GOTO 35
READ (1,1000) X(2), Y(2), IM
CONTINUE
CALL CPLOT (X,Y)
GOTO 35
CONTINUE
CALL CPLOT (X,Y)
GOTO 35
CONTINUE
CONTINUE
CALL CPLOT (X,Y)
GOTO 35
CONTINUE
CALL PAUSE
CALL FIN, 'GIVE SCALE : ')
FORMAT (F10.5)
FORMAT (2F10.3, I 3)
END
```

5.4 Sample Sessions

The programs are too large to fit into the two cylinder disk space which is available to the time sharing user. Therefore only so-called packed versions are present on the private disk.

To log in an additional temporty T-Disk, execute the program S T A R T by calling it from the CMS-Level. The T-Disk is then logged in and eventually files are read in which are in the virtual card reader of the system.

The execution of one of the programs is prepared by

GETTEST

GETWING

GETCAS

or

This means for TEST, WING or CASCADE to be copied to the T-Disk, unpacked, altered to a FORTRAN file and compiled. After this, the input/output files are defined correctly. Finally the execution can be initialized by

\$ FILE NAME

If a compiled version of the desired program is already on the T-Disk,

DOTEST

DOWING

or

DOCAS

will also define the input/output files. Again \$ starts the execution.

See the three sample sessions on the next pages.

CONTENTS OF THE DISK:

1				
FILENAME	FILETYPE	MODE	NO.REC.	DATE
TEST	DATA	P1	19	9/30
FILE4	DATA	F1	3	9/29
PLOT1	PACKED	P5	2	9/19
FILE	FT03F001	P1	1	9/30
FILE	FT04F001	P1	2	9/29
DOTEST	EXEC	P1	1.	9/29
START	EXEC	P1	1	9/29
DOWING	EXEC	P1	1	9/29
CLOSE	EXEC	P1	1	10/01
TEST	PACKED	P5	8	9/29
PLOT2	PACKED	P5	2	9/28
FILE3	DATA	P1	3	9/29
.WING	PACKED	P5	77	9/26
FILE	FT05F001	F1	2	9/29
CASCADE	PACKED	P5	81	9/28
GETWING	EXEC	P1	1	9/29
GETCAS	EXEC	P1	1	9/29
DOCAS	EXEC	P1	1	9/29
GETTEST	EXEC	P1	1	9/29

this is a sampel session for t e s t . condition : you have just lossed in.

```
start
17.30.06 VSET RDYMSG OFF
17.30.07 VSET BLIP /
17.30.08 CP DEFINE T2314 192 10
17.30.09 FORMAT T
** "FORMAT T" WILL ERASE ALL YOUR T-DISK (192) FILES **
**DO YOU WISH TO CONTINUE? ENTER "YES" OR "NO":
>yes
FORMATTING T-DISK (2314)...
T (192): 010 CYL
17.30.30 OFFLINE READ *
READER EMPTY OR NOT READY.
!!! E(00009) !!!
17.30.33 OFFLINE READ *
READER EMPTY OR NOT READY.
111 E(00009) 111
Ri
>settest
17.30.42 VSET BLIP /
17.30.43 COMBINE TES PACKED T5 TEST PACKED P5
17.30.45 UNPACK TES
17.30.46 ALTER TES UNPACKED T5 TEST FORTRAN T1
17.30.47 ERASE TES PACKED T5
17.30.49 F TEST
////17.31.02 FILEDEF FT05F001 DSK
17.31.03 FILEDEF FT01F001 DSK-T1
17.31.05 FILEDEF FT06F001 DSK-T1
RF
>$ test
/EXECUTION BEGINS...
```

this is a sampel session for wing .

condition : t - disk is already lossed in

17.34.53 UNPACK WIN 717.35.01 ALTER WIN UNPACKED TS WING FORTRAN T1 17.35.03 ERASE WIN PACKED TS 17.35.04 F WING 17.34.47 COMBINE WIN PACKED TS WING PACKED PS 17.37.10 FILEDEF FT04F001 DSK 17.37.12 FILEDEF FT06F001 DSK-T1 17.34.46 USET BLIP / /EXECUTION BEGINS... >\$ wing setwins

T O 0 this is a sampel session for disk. ب sou have alreads lossed in the condition :

14.16.37 FILEDEF FT03F001 DSK 14.16.38 FILEDEF FT04F001 DSK-T1 14.16.39 FILEDEF FT07F001 DSK-T1 RECFM F BLKSIZE 133 14.12.35 UNPACK CAS UNPACKED TS CASCADE FORTRAN T1 14.12.24 COMBINE CAS PACKED TS CASCADE PACKED PS 14.12.48 ERASE CAS FACKED TS Ri T=122.98/127.06 14.16.39 14.12.23 USET BLIF / setcas

>\$ cascade

/EXECUTION BEGINS...

6. Critique, Conclusion and Outlook

A systems of programs was written, which allows a systematic approach to the problem of oscillating shaped airfoils in a staggered cascade. TEST is only a geometrical examination of the PROFIL. WING is a test for the theoretical behaviour of the blade, if exposed oscillating to a supersonic stream. Finally, CASCADE is a computer model for a finite cascade of airfoils which may oscillate with a phase lag from blade to blade in a supersonic flow. CASCADE is not yet finished. Results are available up to now only for the inlet flow, as the wake could not be added because of lack of time.

All three programs do not only give numerical but also graphical results. The latter can be obtained by the two plot codes PLOT1 and PLOT2.

CASCADE - and WING -results are compared in this report with solutions of Teipel (2), Verdon (15,17,18) and Platzer and collaborators (3,4,5,6,7,9,13,14). For the linear cases the agreement is considered to be very good. The included thickness effect gives different results to those from Strada (5). This has to be investigated more closely in the near future.

It is clearly understood that C A S C A D E is a rather expensive approach to the problem. This is considered worth while due to the complexity of the problem. After the model of shaped, oscillating airfoils is better understood, it would be desirable to write a more efficient computer code.

To finish this work, there are two more steps to do:

- Adding the procedure of wake computation to C A S C A D E and then obtaining results over the whole blades.
- Optimizing the program in this form.
 This will be done in the near future.

When this is accomplished, a rather flexible logic structure for the problem of an oscillating supersonic cascade is available. This could be used to examine different systems of basic equations and their influence on the results with the final purpose to substitute the potential equations by those of Euler. Thus we could get rid of the assumption of constant entropy in the field which would be a new agep forward.

```
CCMPL EX * 8 U. V. A I. FS I. G. CYDXU.PU

CCMPL EX * 16 EL.RI.ES

CCMMON/BA/ ALO.ALO.AM.CI.AK.I.IE.IW

COMMON/BC/ T. B. DYDX.D 2 YD > 2.D YO XU. AI.II.T3

CCMMON/SF/ X$(2.50).Y$(2.50).A(50.4).R(2.50)

CCMMON/SDL/ EL(4.4).FI(4).E$(4)

CCMMON Y(2.25.25).X(2.25.25).P(2.25).U(2.25.25).P$I(2.25.25).

FG(2.25.25).AL(2.25.25).Y(2.25.25).C(2.3).DX(2).IN(2.25).

FPX(2.20).P$(2.20).PU(2.20).TET(2).TED(20)
OFTIONS:
                             OFTIONS:
LC1 =0 : CCMPLETE OUTPUT
=1 : CNLY PRESSURE CISTRIBUTION
=2 : STOP
LC2 =0 : PITCH
=1 : PLUNGE
LC3 =0 : NC WAKE
LC4 =1 : WAKE SHALL BE COMPUTED
LC4 =1 : SURFACES ARE ANALYTICALLY GIVEN
=2 : WEDGE (NO WAKE)
=0 : SURFACES ARE PCINTWISE GIVEN
MA = NUMBER OF POINTS ON THE SURFACE
                                                                                                                                                                                                                  GIVEN
                              KV=25
AI=CMPLX(0.,1.)
READ(4,599) LO1,LO2,LC3,LO4,MA
IF(LO1.EC.2) GOTO 101
IF(LO4.EC.2) LO3=0
MAX=MA
CALL PRCFIL(WE,LO4,T1,T2,NX)
READ(4,1000) AK,AM,C,B,[X(1),DX(2)
IF(AM.EC.0.) GOTO 102
AX=AM
AM=AM*AM
ALO=AM-1.
C1=C+1.
ALD=ALO**1.5
C3=C1*AM*1.5
 102
 100
                             INITIAL FIELD

XM=1 ./SCRT ($LO)
D=3.1/9.
DC 8 N=1,10
M=N+1
Y(1,M,N)=0.
Y(1,N,M)=0.
X(1,N,M)=0*(N-1)-2.
X(1,N,M)=X(1,M,N)
DC 36 M=2,KV
X(1,M,1)=D*(M-2)/2.-2.
X(1,1,M)=X(1,M,1)
DC 37 N=2,10
L=N+2
DC 37 N=2,10
L=N+2
C1,M,N)=X(1,K,N)+0.5*C
X(1,N,M)=X(1,K,N)+XM*D/2.
Y(1,N,M)=X(1,K,N)+XM*D/2.
Y(1,N,M)=Y(1,M,N)
Y(1,N,M)=Y(1,M,N)
Y(1,N,M)=Y(1,M,N)
Y(1,N,M)=Y(1,M,N)
DC 45 N=1,KV
PSI(1,M,N)=0.
45 N=1,KV
DC 45 N=1,KV
PSI(1,M,N)=0.
U(1,M,N)=0.
U(1,M,N)=0.
T+E STEACY FLOW FIELD
                               INITIAL FIELD
 8
 36
 37
                                THE STEACY FLOW FIELD
```

```
C
            COMPUTATION OF THE SPLINE - COEFFICIENTS
            Ih=0
IE=0
T2=0.
I=I+1
T=T1
IF(I.EG.2) T=T2
DC 46 K=1,NX
A(K,1)=YS(I,K)
A(K,3)= R(I,K)
H=XS(I,K+1)-XS(I,K)
A(K,2)=(YS(I,K+1)-YS(I,K))/H-H+(R(I,K+1)+2.*R(I,K))/2.
A(K,4)=(R(I,K+1)-R(I,K))/(3.*H)
I1=(-1)**(I+1)
MA=MAX
MZ=1
30
46
16
            THE STEACY BOUNDERY-PROFERTIES EEFING THE SHOCK
            CX1=1.07/(MA-1)
DC 9 J=2,KV
KT=J
XF=DX([]*(J-2)
IF(T.EQ.O.) XP=O.
IM=1
CALL SWITCH(J,IM,M,N,I)
CALL SHOCK(KV,WE,LC4,M,N,XP,C3,MA,CX1)
           ALL OTHER STEPS OF THE STEADY FLOD FIELD
7
12
            17(1.Ne.2, GD, G, L, K1=J

K1=J

K2=N+1

A[(2,K1,K2)=AL(2,M,N)

X(2,K1,K2)=(Y(2,J5,J6)-Y(2,M,N))*I1

X(2,K1,K2)=-02*X(2,M,N)+C1*X(2,J5,J6)-X(2,K1,K2)

X(2,K1,K2)=X(2,K1,K2)/(C1-02)
```

```
Y(2,K1,K2)=C2*(X(2,K1,K2)-X(2,M,N))*I1+Y(2,M,N)
IF(I.EC.2) J3=J4
DC 11 K=J3,KA
IF(K.EC.KV) GOTO 11
IM=K+1
CALL SWITCH(IM,J,M,N,I)
IM=J-1
CALL SWITCH(K,IM,K1,K2,I)
X(2,M,N)=X(2,K1,K2)
Y(2,M,N)=Y(2,K1,K2)
Y(2,M,N)=X(2,K1,K2)
CONTINUE
IF(KA.LT.KV) KA=KA+1
IF(MA.GT.IN(I,J-1)) MA=MA+1
CCNTINUE
IF(T.EQ.Q.) GOTO 47
  11
  15
                       ALTOMATIC STEP - SIZE CENTREL
                       IF(J1.GT.MA) DX(I)=DX(I)*1.05
IF(J1.LT.MA) DX(I)=DX(I)*0.96
IF(J1.NE.MA) GOTO 16
000047
                       THE UNSTEADY BOUNDERY PROPERTIES BEHIND SHOCK
                      KA=KT

IF(L01.EQ.0) WRITE(6,1007)

DC 18 J=2,KT

IM=1

CALL SWITCH(J,IM,M,N,I)

CALL RANDS(KV,M,N,DX(I),LC2,M2,N2)

X(2,1,2)=0.

Y(2,1,2)=0.

Y(2,2,1)=0.

Y(2,2,1)=0.
  18
                       ALL OTHER STEPS OF UNSTEADY FLOW FIELD
                       J5=J1-1
DO ZO J=2,J5
I = J+1
CALL SWITCH(I = J, M, N, I)
CALL RANDB(M, N, AM, AK, I,LO2)
                      CALL RANDS (M, N, AM, AK, I

L=J+2

J2=IN (I, J-1)

DO 22 K=L, J2

CALL SWITCH (K, J, M, N, I)

CALL GENU (M, N, I, AK, AM)

IF (J2. EG. KA) GOTO 20

K1=M+1

K2=J

K3=K1

IF (I. NE. 2) GOTO 24

K1=J
   22
                     IF(I.NE.2,

K1=J

K2=N+1

K3=K2

P(I,K3)=x(2,M,N)

CALL RANCS(KV,K1,K2,CX(I),LO2,M2,N2)

J3=J2+1

DC 29 K=J3,KA

IF(K.EC.KV) GOTO 20

TM=K+1
  24
                               EK+T
LL ŞWITCH(IM,J,M,N,I)
                      IN=J-1

CALL SWITCH(K,IM,K1,K2,I)

U(2,M,N)=U(2,K1,K2)

Y(2,M,N)=Y(2,K1,K2)

G(2,M,N)=G(2,K1,K2)

PSI(2,M,N)=PSI(2,K1,K2)

IF(J2,L1,KA,ANO,KA,LT,KV) KA=KA+1

CCNTINUE
   29
  20
```

```
CLTPUT STEACY AND UNSTEACY FIELD
                 IF(L01.NE.0) GOTO 38
WRITE(6,1004)
IF(I.EQ.1) WRITE(6,1002)
IF(I.EC.2) WRITE(6,1003)
J5=J1+1
D( 25 J=2,J5
K=J-1
LL=IN(I,K)
WRITE(6,1007)
WRITE(6,1008)
N=K
                  N=K
IF(I.NE.2) GOTO 26
              P=K
DC 31 N=J,LL
WRITE(6,1009) M,N,X(2,M,N),Y(2,P,N),AL(2,M,N),U(2,M,N),
FV(2,M,N),PSI(2,M,N),G(2,P,N)
GCTO 17
CCNTINUE
DC 32 M=J,LL
WRITE(6,1009) M,N,X(2,M,N),Y(2,M,N),AL(2,M,N),U(2,P,N),
FV(2,M,N),PSI(2,M,N),G(2,M,N)
WFITE(6,1005)
CCNTINUE
26
32
175
CCC38
                 CCMPUTATION OF THE PRESSURE - CCEFFICIENTS
                 K=J1+1
DC 27 J=2,K
IM=J-1
CALL SWITCH(J,IM,M,N,I)
CALL PRESS(M,N,JI)
IF(I.EC.1) IUS=K-1
IF(I.EQ.2) ILS=K-1
IF(I.EQ.1) GOTO 30
27
                 CHANGING THE FIELDS AND WAKE - COMPUTATION
                 DC 48 M=1,KV

DC 48 N=1,KV

X(1,M,N)=X(2,M,N)

Y(1,M,N)=Y(2,M,N)

U(1,M,N)=U(2,M,N)

U(1,M,N)=V(2,M,N)

G(1,M,N)=G(2,M,N)

AL(1,M,N)=AL(2,M,N)

PSI(1,M,N)=PSI(2,M,N)

CALL WAKE(KV,WE,T1,T2,LC4,IUS,ILS,LO3,NX,MA)
48
                  CCRRECTION OF THE INDEX NUMBERS
                 CC 35 I=1,2
K=IUS
IF(I.EQ.2) K=ILS
                IF(I.Eq. 2.

KA=K

DC 43 J=1,K

IF(IN(I,J).GT.K) GOTO 35

KA=KA-1

J1=IN(I,J)

CO 44 L=J1,KA

M=L+1

CY(I.L)=PX(I,M)
                 M=L+1
PX(I,L)=PX(I,M)
FL(I,L)=PU(I,M)
PS(I,L)=FS(I,M)
IF(I,EQ-1) IUS=KA
IF(I,EC-2) ILS=KA
CCNTINUE
                  OLTPUT PRESSURE DISTRIBLTION
                 WRITE(6.1004)
IF(L02.EC.0) WRITE(6,556)
IF(L02.EC.1) WRITE(6,595)
```

```
WFITE(6,1001) AK,AX,C,B,DX(1),T1
WRITE(6,1010)
WFITE(6,1011)
CC 33 I=1,2
IF(I.EQ.2) WRITE(6,1001) AK,AX,C,E,CX(2),T
IF(I.EQ.2) WRITE(6,1012)
IF(I.EQ.2) WRITE(6,1011)
K=IUS
IF(I.EQ.2) K=ILS
WRITE(1,1016) K
CC 34 J=1,K
IM=J+1
CALL SWITCH(IM,J,M,N,I)
WRITE(6,1013) M,N,PX(I,J),PS(I,J),PU(I,J)
WRITE(1,1017) PX(I,J),PL(I,J)
WRITE(1,1017) PX(I,J),PL(I,J)
                                                                                                                             WRITE(6,1001) AK,AX,C,E,CX(2),T2
WRITE(6,1012)
WRITE(6,1011)
                                              INTEGRATION OF THE MOMENTUM AND FORCES, POINTWISE GIVEN, OVER THE SURFACES OF THE AIRFOIL
                                             CALL LIFT(IUS,ILS,LO2)
GCTO 100
K=0
WRITE(1,1016) K
 101
                                  WRITE(1,1016) K

END OF MAIN PROGRAM
FCRMAT(1X, 'PLUNGE - MOCE', /)
FCRMAT(1X, 'FITCH - MOCE', /)
FCRMAT(41,13)
FCRMAT(41,13)
FCRMAT(4510.5,2510.3)
FCRMAT(4510.5,2510.3)
FCRMAT(1X, ', 1X, 'W*C/U=', F5.3, ', M=', F4.2, ', K=', F4.2, ', F*, B/C=', F4.2, ', E2.2, ', E2.2, ', E2.2, ', E3.3, ', T/C=', F6.4, //)
FCRMAT(1X, ', 1HE PROPERTIES AT THE MESHPCINTS OF THE FLOW', F' FIELD', /, 1X, ' FOR THE UPPER SURFACE:', //)
FORMAT(1X, ', THE PROPERTIES AT THE NESHPCINTS OF THE FLOW', F' FIELD', /, 1X, ' FOR THE LOWER SURFACE:', //)
FCRMAT(1X, ', 1X, ' FOR THE LOWER SURFACE:', //)
FCRMAT(1X, ', 1X, ', 10X, 'Y', 7X, 'LAMBDA', 7X, 'RU', 9X, 'IU', FORMAT(1X, ', 12, 11, 13, 13, 14, 15), EX, 'RG', 9X, 'IG', //)
FCRMAT(1X, 'FRESSURE-CISTRIBUTION UPPER SURFACE:', //)
FCRMAT(1X, 'FRESSURE-CISTRIBUTION LOWER SURFACE:', //)
FCRMAT(1X, 'FRESSURE-DISTRIBUTION LOWER SURFACE:', //)
FCRMAT(1X, 'FRESSURE-
C
995
996
999
1000
 1002
   100 3
  100 4
100 5
100 7
100 8
 1009
1010
1011
1012
1013
1014
1016
READ(4,1000) T1,T2,NT,IKF

IF(L04.EC.0) READ(4,1001) SP

IF(L04.EC.2) READ(4,1001) WE

DC 8 J=1,2

IZ=0
N=NT
T=T1

IF(J.EQ.2) T=T2
IF(L04.AE.0) GOTO 12

CC 9 K=1,4

READ(4,1001) XS(J,M),XS(J,M+1),XS(J,M+2),XS(J,M+3),XS(J,M+4)

READ(4,1001) YS(J,M),YS(J,M+1),YS(J,M+2),YS(J,M+3),YS(J,M+4)

IF(YS(J,M+4).EQ.1CO.) GCTO 7

CCNTINUE

CCNTINUE
                                                PREPARATION OF THE PROFIL - SURFACES
  9
```

```
DC 2 LL=1,N

XS(J,LL)=XS(J,LL)/SP

YS(J,LL)=YS(J,LL)/SP

IF(L04.EC.0) GOTO 3

0X=1.0/(N-1)

I1=(-1)**(J+1)

DC 20 K=1,N

XS(J,K)=(K-1)*CX-0.25

YS(J,K)=I1*4.*T*XS(J,K)*(1.-XS(J,K))
  2
  12
NOU COMO CO
                    TC ENTER THIS PART, THE SURFACES SHOULD ALREADY EEGIVEN POINTWISE EXTRAPOLATING POINTS
                     CCNTINUE
                     INTERPOLATION THROUGH CUBIC SPLINES
                    IF(T.NE.O. .OR. LO4.EG.() GCTC 51

D( 52 I=1,N

A(I,1)=0.

GCTO 53

CCNTINUE

DO 10 I=1,N

A(I,3)=XS(J,I)

A(I,4)=YS(J,I)
  52
  51
  10
                     MATRIX OF COEFFICIENTS AND RIGHT SICES
                 K=N-2

DC 25 I=1,K

A(I,1)=A(I+1,3)-A(I,3)

A(I,3)=A(I+2,3)-A(I+1,3)

A(I,2)=2.*(A(I,1)+A(I,3))

A(I,4)=3.*(A(I+2,4)-A(I+1,4))/A(I,3)-

F3.*(A(I+1,4)-A(I,4))/A(I,1)

CCNTINUE

A(1,1)=C.0

A(N-2,3)=0.0
  25
                    THE STEP OF GAUSS
                    K=N-3

OC 30 I=1,K

CC 35 M=3,4

A(I,M)=A(I,M)*A(I+1,1)/A(I,2)

A(I,1)=0.0

A(I,2)=A(I+1,1)

A(I+1,2)=A(I+1,2)-A(I,3)

A(I+1,4)=A(I+1,4)-A(I,4)
  35
  3000
                     SCLUT ION
                    A(1,1)=0.
A(N,1)=C.
A(N-2,1)=0.C
                     L=N-1
CC 40 I=2,L
K=N-I
                   CC 40 I=2,L
K=N-I
M=K+1
A(M,I) = (A(K,4)-A(K,3)*A(M+1,1))/A(K,2)
CCNTINUE
DO 11 M=1,N
R(J,M)=A(M,1)
IF(IZ.EQ.1) GOTC &
KV=N-1
CC 49 K=1,KV
A(K,1)=YS(J,K)
A(K,3)= R(J,K)
H=XS(J,K+1)-XS(J,K)
A(K,2)=(YS(J,K+1)-YS(J,K))/H-H*(R(J,K+1)+2.*R(J,K))/3.
A(K,4)=(R(J,K+1)-R(J,K))/(3.*H)
DX=1./49.
DC 4 M=1,50
  40
  11
                                                                                                                                                                                                  1 20
  49
```

```
X=DX*(M-1)

DC 5 K=2,50

I=K-1

IF(XS(J,K).GE.X) GOTO 6

CCNTINUE

H=X-XS(J,I)

Y=A(I,1)+A(I,2)*H+A(I,3)*H*H+A(I,4)*H*H*H

R(J,M)=X

YS(J,M)=Y

CONTINUE

IZ=1
 5
                      IZ=1
N=50
D( 13 I=1.50
XS(J, I)=R(J,I)
GCTO 3
CCNTINUE
IF(IKP.EQ.O) GOTO 17
DC 15 M=1.N
WRITE(6,1002) XS(1,M),YS(1,M),XS(2,M),YS(2,M)
N=N-1
FORMAT(2F10.5,2I2)
FORMAT(5F10.5)
FORMAT(5X,4(E12.5,2X))
RETURN
END
 13
 8
 15
17
1000
1001
1002
FIND LOOKS FOR THE MESHINDEX, RESPONSIBEL FOR THE POINT XS; YS
                     IF(IE.EC.1) GOTO 19
IKK=0
IF(IW.GT.5 .AND. XS.GE.1. .AND. XS.LE.1.10) IKK=1
IF(IKK.EQ.1) WRITE(6,2C11) I,IE,I1
FCRMAT(IX,'FIND ENTRY: ',3I4)
IZ=0
KA=0
L1=1
L2=1
I9=21
IF(IKK.EQ.1) WRITE(6,2003) I9,I,IE,XS,YS
DO 22 I2=L1.20
IM=I2+L2
CALL SWITCH(IM,I2,M,N,I)
L3=I2+L2
CALL SWITCH(IM,I2,M,N,I)
L3=I2+L2
IF(IKK.EQ.1) WRITE(6,2CCC) I9,M,N,X(1,M,N),Y(1,M,N)
IF(L3.GE.KV) GOTO 13
J1=M+1
IR=J1-J2
IF(I.EQ.2) KM=J2-J1
IS=24
C
 2011
 21
                      IF(I.EQ.2) KM=J2-J1
IS=24
IF(IKK.EQ.1) WRITE(6,20CC) I9,M,N,X(1,M,N),Y(1,M,N)
IF(XS.GE.X(1,M,N) .AND. XS.LT.X(1,J1,J2)) GOTO 23
IF(XS.GE.X(1,M,N) .AND. XS.LT.X(1,J1,J2)) GOTO 23
IF(I.EQ.2 .AND. KA.EQ.J2) KA=J1
IF(I.EQ.1 .AND. KA.NE.J2) GOTO 22
IF(I.EQ.2 .AND. KA.NE.J1) GOTO 22
IF(I.EQ.2 .AND. KA.NE.J1) GOTO 22
IF(I.EQ.2) J2=J2+1
IF(J.EQ.1) J1=J1+1
IF(J.EQ.KV .OR. J2.EQ.KV) GOTO 13
GCTO 24
IS=22
IF(IKK.EQ.1) WRITE(6,2CCC) I9,M,N,X(1,M,N),Y(1,M,N)
IS=23
  22
                        IS=23
IF(IKK.EQ.1) WRITE(6,2000) I9,M,N,X(1,M,N),Y(1,M,N)
  23
```

```
[F(I.EC.2) GOTO 27
[F(YS.LT.Y(I.JI,N)) GCTC 1
[F(N.EQ.1) GOTO 13
_2=L2+1
_1=N-1
GCTO 21
[9=27
[F(IKK.EQ.1) WRITE(6,2C(0) I9,M,N,X(1,M,N),Y(1,M,N)
[F(YS.GT.Y(1,M,J2)) GCTC 1
[F(M.EQ.1) GCTO 13
_2=L2+1
_1=M-1
[CTO 21
[9=1]
   27
                                                                                                          F(IKK.EQ.1) WRITE(6,2C(0) 19,M,N,X(1,M,N),Y(1,P,N)
     1
                                                                                        IF(IKK.eu...

M2=M

N2=M

N2=N

GCTO 14

IS=28

IF(IKK.EC.1) WRITE(6,200C) I9,M,N,X(1,M,N),Y(1,M,N)

IZ=IZ+1

IF(IZ.EQ.5) GOTO 29

M=M2
     28
                                                                                     M=M2
N=N2
IS=18
IF(IKK.EQ.1) WRITE(6,20CC) I9,M,N,X(1,M,N),Y(1,M,N)
IF(IZ.EC.0) GDTO 28
IF(IZ.EC.1 .AND. I.EQ.2) M=M-1
IF(IZ.EQ.1 .AND. I.EQ.2) M=M-1
IF(IZ.EQ.1 .AND. I.EQ.2) M=M+1
IF(IZ.EQ.2 .AND. I.EQ.1) M=M+1
IF(IZ.EQ.2 .AND. I.EQ.2) N=N+1
IF(IZ.EQ.2 .AND. I.EQ.2) N=N+1
IF(IZ.EQ.3 .AND. I.EQ.2) N=N+1
IF(IZ.EQ.3 .AND. I.EQ.2) M=M+1
IF(IZ.EQ.3 .AND. I.EQ.3 M=M-1
IF(IZ.EQ.3 .AND. I.EQ.3 M=M-1
IF(IZ.EQ.3 .AND. I.EQ.3 M=M+1
IF(IZ.EQ.3 .AND. 
   18
   3
     2
   29
                                                                                               IE=1
WFITE(6,1001) M,N,XS,YS,IE
GOTO 19
00014
                                                                                               CCNTRCL - STEP
                                                                                          J1=M+1
J2=N+1
K1=M+1
K2=N
K3=M
                                                                                                   K4=N+1
IF(I.NE.2) GOTO 20
     20
                                                                                                                                                                                                                                                                                AND. M.EG.14

WRITE(6,2000)

WRITE(6
   11
```

```
IF(K4.EQ.K3) GOTO 25
D3=I1*(Y(1,J1,J2)-Y(1,K2,K4))/(X(1,J1,J2)-X(1,K3,K4))
IF(K4.EQ.K3) D3=0.
D4=I1*(Y(1,J1,J2)-Y(1,K1,K2))/(X(1,J1,J2)-X(1,K1,K2))
D34=I1*(Y(1,J1,J2)-YS))/(X(1,J1,J2)-XS)
IF(D1.LT.D12 .OR. C2.ET.D12) GCTC 18
IF(D3.LT.D34 .OR. D4.GT.C34) GCTC 18
IF(XS.LT.X(1,M.N) .OR. XS.GE.X(1,J1,J2)) GCTC 18
IF(XS.LT.X(1,M.N) .OR. XS.GE.X(1,J1,J2)) GCTC 18
FCRMAT(1X, NO-FIND : ",I2,",",I2,4(2X,F8.3))
FCRMAT(1X, FIND: ",I2,",",I2,2(2X,E10.3)," IE= ",I2)
FCRMAT(1X,3I4,2X,F8.3,2X,F8.3)
RETURN
END
**TURN
END
**TURN
END
**TURN
**T
  25
STEADY BOUNDERY CONDITIONS ALONG THE AIRFOIL
                                  IF(LO4.EC.1) GOTO 3
IF(LO4.EC.2) GOTO 4
DC 5 K=2,50
J=K-1
IF(XS(I,K).GE.X) GOTO 6
CONTINUE
H=X-XS(I,J)
DYDX=A(J,2)+2.*A(J,3)*H+3.*A(J,4)*+++
D2YDX2=2.*A(J,3)+6.*A(J,4)*+
GCTO 2
DYDX=I1*4.*T*(1.-2.*X)
D2YDX2=-I1*8.*T
GCTO 2
W=WE*4.*ATAN(1.)/180.
CYCX=I1*TAN(W)
D2YDX2=0.
GOTO 2
GOTO 2
  5
  3
CCC1
                                      STEADY EQUNDERY CONDITIONS ALONG MAKE-SLIP-LINE
CCMPUTATION OF THE FIELD NEAR BEHIND THE SHOCK
                                      ITE=0
IKK=0
IF(IW.NE.0) IKK=1
IF(IKK.EC.1) WRITE(6,2000) M,N,I,IE
FORMAT(IX,'SHOCK ENTRY: ',414)
   2000
                                     TX=0.

TX=0.

IF(IW.NE.0) TX=1.

J=M-1

L=M

IF(I.NE.2) GOTO 5

J=M

K=N-1

L=N
                                     CALL BCUND (WE, LO4, I, XP, IW)
    5
```

```
AL(2,M,N)=(ALD-C3*DYDX*I1)**(2./3.)
IF(L.EC.2) GOTO 17
D1=I1/SQRT(AL(2,M,N))
X5=X(2,J,K)
Y5=Y(2,J,K)
CALL FIND(KV,IW,X5,Y5,M2,N2,IE,I)
IK=9
IF(IKK.EQ.1) WRITE(6,1000) IK,IE,AL(2,M,N),X5,Y5
                                                                          F(IKK.EQ.1) WK11ELG,1000

[L=9]

F(IE.EG.1 .OR. M2.GE.KV) GOTO 16

.1=M2+1

.2=N2

IF(I.NE.2) GOTO 7

.1=M2

L2=N2+1

x1=x(1,M2,N2)

x1=x(1,M2,N2)

x2=x(1,L1,L2)

x2=x(1,L1,L2)

x3=x(1,M2+1,N2+1)

x3=x(1,M2+1,N2+1)

x3=x(1,M2+1,N2+1)

x3=x(1,M2+1,N2+1)

x3=x(1,M2+1,N2+1)

x3=x(1,M2+1,N2+1)
7
                                                                              (3=Y(1, M2+1, N2+1)

(7=0)

(2=0)

(2=20RT(AL(1, M2, N2))+SCFT(AL(2, J, K))

(3=2.*11/D2

(4=7)

(4=14)

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 14
                                                                                4=Y(2,J,K)+(x4-x(2,J,K)+02

K=14

F(IKK.EQ.1) WRITE(6,10CQ) IK,IE, x4, y4, x5, y5

CTO 13

(4=(D1*xp-D2*x5+y5)/(D1-D2)

(4=(X4-XF)*D1

F(X4.LE.TX) IT==1

F(ITE.EC.1) GOTO 14

K=2

F(IKK.EC.1) WRITE(6,10CC) IK,IE, x4, y4, x5, y5
2
                                                              13
  19
 20
 10
 8
```

```
16
                        2=2. * SQRT(ALO)+SQRT(AL(2, J, K))+SCFT(AL(2, M, N))
2=4. * 11/D2
                    DZ=4.#11/02
IK=16
IK=16
IF(IKK.EC.1) WRITE(6,10C0) IK, IL,
IF(IKK.EQ.1) WRITE(6,1CC0) M,N,C1
IF(T.EQ.C.) GOTO 1
X(2, M,N)=(D1*XP-D2*X5+Y5)/(D1-D2)
Y(2, M,N)=(X(2, M,N)-XP)*C1
GCTO 11
X(2, M,N)=X5+DX1/2.
Y(2, M,N)=D2*DX1/2.+Y5
                                                                                                           IK, IL, >F, X5, Y5, X (2, M, N)
M, N, C1, C2, DX1
 C
C
C
C
11
                    CCMPUTATION OF THE ACCITIONAL FCINTS FOR THE UNSTEADY
CCC
                     BCUNDERYSTEP OF THE STEACY FIELD TO THE AIRFOIL
                    IKK=0

IF(LO4.EC.O) IKK=1

IF(IKK.EQ.1) WRITE(6,1002)

FCRMAT(1X, RAND-ENTRY)

A=-C1*I1

J1=M

J2=N-1

IF(I.NE.2) GOTO 5
 C
 100 2
                   J2=N-1
J2=N
J2=N
X(2,M,N)=X(2,J1,J2)
OC II Kx=[,15
CALL BOUNC(hE,LO4,I,X(2,M,N),IM)
IF(IKK.EC.1) WRITE(6,1001) M,N,KX,X(2,M,N),DYC>,D2YCX2
FCRMAT(1X,313,3E12.5)
AL(2,M,N)=(ALD+A+DYDX)++(2./3.)
IF(IW.NE.O) GOTO 2
IF(DYDX.EQ.O. AND. D2YCX2.EQ.O.) GCTO 2
F=2.*(X(2,M,N)-X(2,J1,J2))+I1/Y(2,J1,J2)
F=F-(SQRT(AL(2,J1,J2))+SCRT(AL(2,M,N)))
FS=2.*I1/Y(2,J1,J2)-A+C2YDX2/(3.*AL(2,M,N))
D=F/FS
X(2,M,N)=X(2,M,N)-C
IF(ABS(C).LE.O.0000001) GCTO 10
CCNTSINUE
WRITE(6,1000) M,N,D
FCRMAT(1X,'EDUNDERYSTEP DID NCT CCNVERGE: ',I2,',',I2,
F' D= ',E10.3)
Y(2,M,N)=O.
CALL BOUND(hE,LO4,I,X(2,M,N),IW)
AL(2,M,N)=(ALD+A*OYDX)++(2./3.)
GCTO 1
 5
 1001
 11
 1000
 10
```

```
Y(2,M,N)=0.
D=-2.*I]/(SQRT(AL(2,J1,J2))+SQRT(AL(2,M,N)))
X(2,M,N)=X(2,J1,J2)-Y(2,J1,J2)/D
RETURN
   2
   GENERAL STEP OF STEADY FIELD
IKK=0
IF(IKK.EQ.1) WRITE(6,10CC)
FORMAT(1X, GEN - ENTRY)
            AN4,S)

IK=5

IF(IKK.EQ.1) WRITE(6,1001) IK,M,N,U(1,M2,N2),V(1,M2,N2),

AL(1,M2,N2)

DV=V(1,M2,N2)

DU=U(1,M2,N2)

ALT=AL(1,M2,N2)

IF(IW.EQ.0.) TX=0.

IF(IE.EQ.0.) ALT=ALO

IF(IE.EQ.1) DV=0.

IF(IE.EQ.1) DU=0.

IF(IE.EQ.1) DU=0.

CALL BOUNDU(IW,TX,I,AK,L)

V(2,M,N)=DYDXU
```

```
D2=AM1*V(2, M, N)/A h1+DU*(AM3-AM1*Ah2/AN1)
U(2, M, N)=D2+DV*(AM4-AM1*AN4/AN1)
K1=M+1
K2=N
IF(I.NE.2) GDTD 4
K1=M
K2=N+1
D1=(AL(2, M, N)+AL(2, K1, K2))/2.
                                                                    D1={AL(2,M,N)+AL(2,K1,K2)}/2.
IK=4
IF(IKK.EQ.1) WRITE(6,1001) IK,M2,N2,C1,X1,TX,S
D1=I1/SQRT(D1)
D3=I1/SQRT(AL(2,K1,K2))
XX=-Y(2,K1,K2)/D3+X(2,K1,K2)-TX
IF(T.20,00.) XX=0.
X(2,M,N)=D1*XX/(TAN(S)+D1)+TX
Y(2,M,N)=(X(2,M,N)-TX)*TAN(S)
PSI(2,M,N)=PSI(1,M2,N2)
P(I,3)=XX
GCTO 9
J1=M-1
                                                                     GCTO 9
J1=M-1
J2=N
LP=M
IF(I.NE.2) GOTC 8
J1=M
J2=N-1
LF=N
CCNTINUE
IF(AL(2,J1,J2).LT.AL(2,M,N).OR. T.EQ.O.) GOTC 12
IF(IW.NE.O) GOTO 12
J2=N-1
6
8
                                                                     J2=N-1

IF(I.NE.2) GOTO 12

J1=M-1

D0=X(2, M,N)-P(I,LP)

DAL=AL(2,M,N)
                                                              D0=x(2, M, N) -P(I, LP)

D4=AL(2, M, N)

IE=O

CALL FIND(KV, IW, X(2, J1, J2), Y(2, J1, J2), M3, N3, IE, I)

DU=U(1, M2, N2)+U(1, M3, N3)

V=IO00000*(P(I, LP)-X(2, J1, J2))

IF(MP, LT.O) WRITE(6, 1CC2) M, N, P(I, LP), X(2, J1, J2)

IF(MP NE.O) GOTC 3

O1=0.

GCTO 2

D1=(AL(2, M, N) -AL(2, J1, J2))*D0

IK=3

IF(IKK. EG.1) WRITE(6, 1CC1) IK, IM, IE, CU, CV, CAL

C1=01/(4,*DAL*(P(I, LP)-X(2, J1, J2)))

D2=AI*AK*AM*DO/DAL

C1=01/(4,*DAL*(AK*DO))**Z/DAL

D4=1./SQRT(AL(2, M, N))

D5=-D3*C4

D6=1.*AK*AK*AM*PSI(2, J1, J2)*DO/CAL

D7=Y(2, M, N)-Y(2, J1, J2)

A(1, 1)=1.*D1+D2+D3

A(1, 2)=-(C4+D5)*I1

A(1, 3)=C.

RI(1)=U(2, J1, J2)*(2.-A(1, 1))-Y(2, J1, J2)*(D4-D5)*I1+C6

A(2, 3)=-2.*AM1/D7-AI*AM2)*G(2, J1, J2)-U(2, J1, J2)

A(2, 1)=1.

A(2, 1)=1.

A(2, 1)=0.

A(3, 2)=0.

A(3, 2)=1.

NC=3

CC=3

CC=4

CC
12
1002
3
 2
C
                                                                    NC=3
CALL SOLVE(NO)
U(2,M,N)=ES(1)
V(2,M,N)=ES(2)
G(2,M,N)=ES(3)
```

```
CU=U(2,M,N)+U(2,J1,J2)+(\(2,M,N)+V(2,J1,J2))*D4*I1
PSI(2,M,N)=EU*0.5*D0+PSI(2,J1,J2)
RETURN
FCRMAT(1x,3I4,5(2x,E11.4))
7
1001
UNSTEACY BOUNDERY - CONCITIONS ALONG THE AIRFOIL
BCUNDARY STEP OF THE LISTEADY FIELD TO THE AIRFOIL
     GENERAL STEP OF UNSTEACY FIELD
     J1=M
J2=N-1
J2=M-1
J4=N
IF(I.NE.2) GD TD 5
J1=M-1
J2=N
J2=M
J4=N-1
```

```
DC=X(2,M,N)-X(2,J1,J2)
D1=0.5*(AL(2,M,N)-AL(2,J1,J2))
D2=AI*2.*AK*AM*D0
D3=-0.5*AM*(AK*D0)**2
D4=AL(2,M,N)+AL(2,J1,J2)
D5=SQRT(AL(2,M,N))+SQRT(AL(2,J1,J2))
D6=(D1+D2+D3)/D4
A(1,1)=1.+D6
A(1,2)=2.*(1.-D3/C4)*I1/D5
RI(1)=U(2,J1,J2)*(1.-D6)+V(2,J1,J2)*I1*2.*(1.+D3/D4)/D5
RI(1)=2.*AK*AK*AM*PSI(2,J1,J2)*CG/D4+RI(1)
BC=X(2,M,N)-X(2,J3,J4)
B1=(AL(2,M,N)-AL(2,J1,J2))*BO/(2.*D0)
B2=2.*AI*AK*AM*PO
B3=-0.5*AI*AK*AM*PO
B3=-0.5*AI*AK*AM*PO
B3=-0.5*CFT(AL(2,M,N))
B5=-0.5*(B1+B2+B3)/AL(2,M,N)
B6=0.5*(B1+B2+B3)/AL(2,M,N)
A(2,1)=1.*B6
A(2,2)=-(B4+B5)*I1
RI(2)=U(2,J3,J4)*(1.-B6)+V(2,J3,J4)*I1*(-B4+B5)
RI(2)=PSI(2,J3,J4)*AM*PC*AK*AK/AL(2,M,N)+RI(2)
NG=2
  5
  C
                        NC=2
CALL SCLVE(NO)
U(2,M,N)=ES(1)
V(2,M,N)=ES(2)
D2=V(2,J1,-2)+V(2,M,N)
D2=(U(2,J1,J2)+U(2,M,N)-I1*2.*D2/C5)*D0/2.
D6=(V(2,J3,J4)+U(2,M,N))*B4
D6=(U(2,J3,J4)+U(2,M,N)+I1*D6)*B0/2.
D2=PSI(2,J1,J2)+PSI(2,J3,J4)+D2+D6
PSI(2,M,N)=C2/2.
RETURN
END
CCMPUTATION OF THE PRESSURE - COEFFICIENTS ALCNG THE AIRFOIL PU=CPU - PS=CPS
                        K=N
IF(I.EG.2) K=M
IF(K.EG.J1) GOTO 5
PS(I,K)=-2.*(AL(2,M,N)-ALC)/(C1*AM)
PU(I,K)=-2.*(U(2,M,N)+AI*AK*PSI(2,M,N))
PX(I,K)=X(2,M,N)
GOTO 6
K1=K-1
K2=K-2
D2=PX(I,K1)-PX(I,K2)
D1=(1.-PX(I,K1))/C2
PS(I,K)=PS(I,K1)+(FS(I,K1)-PS(I,K2))*D1
PU(I,K)=PU(I,K1)+(PU(I,K1)-PU(I,K2))*D1
PU(I,K)=I.
RETURN
END
  5
 6
```

```
COCO
                                                        COMPUTATION OF THE LIFT- AND MOMENTUM - COEFFICIENTS
FOR BOTH OF THE SURFACES AND FOR THE COMPLETE AIRFOIL
LIFT: CL,CLS ; MOMENTUP: CM,CMS
                                                TOR BOTT OF THE SERFACES AND FOR THE COMPLETE MARKET OF COMPLETE MARKE
   5
    7
    8
    1000
    100 1
1002
1003
  CCMPUTATION OF THE FIELD BEHIND THE AIRFOIL
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                     91
                                                          IKK=0
IF(LO3.EC.0) GOTO 1
OXS=0.5*(DX(1)+DX(2))
EX=E1*AM2*1.5
E=E1-1.
```

```
E5=E-1.
Ih=0
IE=0
                 COMPUTATION OF THE SPLINE COEFFICIENTS
                DC 2 I=1,2

T=TA

If(I.EG.2) T=TB

I=(-1)**(I+1)

DC 46 K=1,NX

A(K,1)=YS(I,K)

A(K,3)= R(I,K)

H=XS(I,K+1)-XS(I,K)

IS=46

IF(IKK.EQ.1) WRITE(6,2CCC) I9,K,I,JX,H,XS(I,K)

A(K,2)=(YS(I,K+1)-YS(I,K))/H-H*(R(I,K+1)+2.*R(I,K))/3.

A(K,4)=(R(I,K+1)-R(I,K))/(3.*H)
                 THE STEADY WAKE
                 XXS=1.

YYS=0.

CALL FIND(KV,IW,XXS,YYS,M,N,IE,I)

I 9=2

IF(IKK.EC.1) WRITE(6,2000) I9,IW,M,N,X(1,M,N)

TE=2.*I1*(ALD-AL(1,M,N)**1.5)/3.

TE=TE/(AL(1,M,N)+E*AM2+1.)

TET(I) =ATAN(TE)

I 5=2

IF(IKK.EC.1) WRITE(6,2000) I9,IW,M,N,DYDX

M=2
2
              F1
M=2
N=1
XF1=1.
XP2=1.
3=0
7=4
                   3=0

9=4

F(IKK.EQ.1) WRITE(6,20CC) I9,IW,M,N,TET(1),TET(2)

)X1=DXS

)X2=DXS

)XU=1./(MA-1)

:XL=DXU

)C 6 I4=1,30

(P3=XP1+DX1

(F4=XP2+CX2

IF(IW.EQ.0) GOTC 36
                    F(IW.LE.2) T4=TED(1)
F(IW.LE.2) GOTC 28
I=IW-1
C 29 I2=1, II
F(X(2, I2+1, I2).GT.XP3) GOTO 30
                    F(X(2,12+1,12).GT.XP3) GUTU 30

S=29

F(IKK.EQ.1) WR ITE(6,2000) I 9,I W, P, N, XP

4=TED(I2-1)

S=30

F(IKK.EQ.1) WRITE(6,2000) I 9,I W, P, N, T4

=1

=1
29
                                                        WR ITE(6,2000) I 9, I W, F, N, XP3
30
28
                     9=28
F(IKK.EQ.1) WRITE(6,2000) 19,1%, M, XP3, XP4
                    TA
1=1
ALL SHOCK (KV, WE, LO4, P, N, XP3, EX, P2, OXU)
F(IE.EG.1) GOTO 17
=2
1=-1
=TB
F(IW.LE.2) GOTO 35
O 33 I2=2; II
F(X(2, I2, I2+1).GT.XP4) GOTO 34
                  IS=33
IF(IKK.EC.1)
T4=TED(I2-1)
IS=35
33
                                                         WRITE(6,20CC) 19, 1W, P, N, XP4
34
```

```
IF(IKK.EC.1) WRITE(6,2000) I9,IW,P,P,T4
IF(IKK.EC.1) WRITE(6,2000) I9,IW,P,P,EX
CALL SHCCK(KV,WE,LO4,P,M,XP4,EX,MA,CXL)
IF(IKC.EQ.1) WRITE(6,2000) I9,IW,P,P,X(2,1,P)
IF(IW.GT.1) GOTO 17
IF(IW.GT.1) GOTO 12
Ih=IW+1
M=M+1
GCTO 4
I9=12
IF(IKK.EQ.1) WRITE(6,2000) I9,IW,P,N,X(2,M,1)
IF(IW.EC.2) GOTO 7
IX=IW-1
OC 5 N=2,IX
I=1
I]=1
12
                                            1=1
ALL GEN(M,N,I,IW)
=2
                                           1=-1
ALL GEN(N,M,I,IW)
9=5
F(IKK.EQ.1) WR ITE(6,2000) I9,IW,M,N,X(2,M,N)
00 8 I2=1,2
M=N+1
ALL SWITCH(IM,N,L1,L2,I2)
5
                                             P=N+1
ALL SWITCH(IM, N, L1, L2, I2)
S=11
F(IKK.EC.1) WRITE(6, 2000) I9, IW, L1, L2, AL(2, L1, L2)
1=(-1)**(I2+1)
E=2.*I1*(ALD-AL(2, L1, L2)**1.5)/3.
E=TE/(AL(2, L1, L2)+E*AM2+1.)
ET(12) = ATAN(TE)
8
                                           ET(12) = AT AN(TE)
9=36
F(IK. EQ. 1) WRITE(6, 2 CCC) I9, IW, M, N, TET(1), TET(2)
1=TET(1)
2=TET(2)
3=0.5*(T1+T2)
C 47 IA=1,2
M=N+1
ALL SWITCH(IM, N, L1, L2, IA)
F(N, GT, 1) GOTO 53
1=(-1)**(IA+1)
(XS=1.**(YS=0.**)
(YS=0.**(YS=0.**)
                                     XXS=1.
YYS=0.
IE=0
CALL FIND(KV, IW, XXS, YYS, M2, N2, IE, IA)
IF(N.EQ-1) ALT=AL(1, M2, N2)
IF(N.GT-1) ALT=AL(2, L1, L2)
XM=(1.-(ALT-AL0)*E/E1)**(E5/E)
XM=2.*((1.+E5*AM2/2.)/XM-1.)/E5
IG=47
IF(IKK.EQ-1) WRITE(6, 20CC) I9, M, N, IA, XM, ALT
IF(IA.EQ-1) XM1=XM
IF(IA.EQ-2) XM2=XM
PII=((1.+E5*XM1/2.)/(1.+E5*XM2/2.))**(E/E5)
CALL KONST(Q(1, 1), C(1, 2), Q(1, 3), XM1, E1)
CALL KONST(Q(2, 1), Q(2, 2), Q(2, 3), XN2, E1)
DC 9 K=1, 20
T4=T3-T1
T5=T3-T2
F=E*XM1*(Q(1, 1)*T4+Q(1, 2)*T4+T4+C(1, 3)*T4**3)/2.+1.
FS=E*XM2*(-Q(2, 1)+2.*C(2, 2)*T5*T5-Q(2, 3)*T5*T5)/2.*PII
FS=E*XM2*(Q(2, 1)-2.*C(2, 2)*T5+3.*C(2, 3)*T5*T5)/2.*PII
FS=FFFFS+E*XM1*(Q(1, 1)+2.*C(1, 2)*T4+3.*C(1, 3)*T4*T4)/2.
D=F/FS
UF(ABS(0).LE.0.000001) GOTO 10
CCNTINUE
WRITE(6, 1000) M,N,T3,C
IS=10
IF(IKK.EQ.1) WR ITE(6, 2000) I9,IW, N,N,T3
IF(IW.EQ.0) IW=1
IF(IW.EQ.1) GOTC 4
 53
47
9
 10
```

```
N=N+1
IS=7
7
                        if(ikk.Eq.1) write(6,2000) 19,1w,+,N,X(2,M,N)
T4=TED(im)
                        T=TA
                       CĂLÊ RAND(WE, LO4, M, N, E1)
I 1=-1
I =2
                        Î =2
T=TB
                       CALL RAND(he,LO4,N,M,E1)
D=X(2,M,N)-X(2,N,M)
IF(ABS(D).LE.O.0001) GCTC 14
IF(X(2,M,N).LT.X(2,N,M) .OR. I3.EC.1) GCTG 15
16
                       13=2
19=16
1F(IKK.EQ.1) WRITE(6,2C(C) 19,1W,M,N,X(2,M,N),D
DX1=DX1*((X(2,N,M)-1.)/(X(2,M,N)-1.))**(1.+0.08*(N-2))
GCTO 6
15=15
IF(IKK.EC.1) WRITE(6,2CCC) 19,1W,M,N,X(2,N,M),D
IF(13.EQ.2) GCTO 16
15
                        DX2=DX2*((X(2,M,N)-1.)/(X(2,N,M)-1.))**(1.+0.08*(N-2))
CONTINUE
WRITE(6,100G) M,N,X(2,M,N),X(2,N,M),C
6
                      WRITE(6, 1000) M,N,A(2, N, N, N, N, S(2, M, N) I S=14

I F(IKK.EQ.1) WRITE(6, 2CCQ) I 9, I h, M, N, X(2, M, N) X(2, M, N) = (X (2, M, N) + X (2, N, M))/2.

X(2, N, M) = X(2, M, N) X F1 = X P2

XP2 = XP4

I W = I W + 1

M = M + 1

GCTO 4
14
CCC 17
                        THE UNSTEADY WAKE
                       THE INITIAL STEP

IW=IW-1

I=1

M=2

N=1

I1=1

CALL CCEF1 (KV,I1
                   CALL CCEF1 (KV,I1,M,N,M2,N2,AM1,AMZ,AM3,AM4,AN1,AN2,AN3,
FAN4,S1)
IF(IKK,EC.1)WRITE (6,1009)M2,N2,AM1,AMZ,AM3,AM4,AN1,AN2,
FAN3,AN4,S1
I=2
                  FAN3,AN4,S1
I=2
I1=-1
CALL CCEF1 (KV,I1,N,M,M3,N3,AM5,AM6,AM7,AM8,AN5,AN6,AN7,FAN8,S2)
I F(IKK.EG.1) WRITE (6,1005)M3,N3,AM5,AM6,AM7,AM8,AN5,AN6,FAN7,AN8,S2
ALU=U(1,M2,N2)*(AM3-AN1*AN3/AN1)+V(1,M2,N2)*(AM4-A*1*AN4/AN1)
ALL=U(1,M3,N3)*(AM7-AM5*AN7/AN5)+V(1,M3,N3)*(AM8-AM5*AN8/AN5)
PSI(2,M,N)=PSI(1,M2,N2)
PSI(2,N,M)=PSI(1,M3,N3)
T4=TED(1)
CL=(PSI(2,N,M)-PSI(2,M,N))*(1.-AM1*TAN(T4)/AN1)*AI*AK
CL=ALL-AUU+CL
V(2,N,M)=CL/(AM1/AN1-A*5/AN5)
U(2,N,M)=CL/(AM1/AN1-A*5/AN5)
U(2,N,M)=CL/(AM1/AN1+AK*(PSI(2,N,M)-PSI(2,M,N))
V(2,M,N)=U(2,N,M)+AI*AK*(PSI(2,N,M)-PSI(2,M,N))
G(2,M,N)=O.
G(2,N,M)=O.
G(2,N,M)=O.
G(2,N,M)=O.
CC 24 I=1,2
I1=(-1)**(I+1)
IM=M+I
CALL SWITCH(M,N,L1,L2,I)
CALL SWITCH(M,N,L1,L2,I)
CALL SWITCH(I*,N,K1,K2,I)
S=S1
```

```
IF(I.NE.2) GOTO 27
S=S2
IF(IKK.EQ.1) WRITE(6,1009) L1,L2,X(2,L1,L2),Y(2,L1,L2),
FX(2,K1,K2),Y(2,K1,K2),AL(2,K1,K2),T4,S
D1=1./SQRT(AL(2,L1,L2))
D1=-I1*D1
D=:I1/SCRT(AL(2,K1,K2))
XX=-Y(2,K1,K2)/D3+X(2,K1,K2)
IF(Y4.EC.C.) XX=1
X(2,L1,L2)=(D1*XX-TAN(S))/(D1-TAN(S))
Y(2,L1,L2)=(X(2,L1,L2)-1.)*TAN(S)
P(I,3)=XX
X(2,L1,L2)=1
Y(2,L1,L2)=0.
P(I,3)=1.
27
                                   FURTHER STEPS
                                  LC2=0
DC 22 I=1,2
IE=0
I1=(-1)**(I+1)
T=TA
IF(I.EQ.2) T=TB
DO 22 J1=3,IW
IM=1
                                DO 22 J1=3, IW

DO 22 J1=3, IW

I M=1

CALL SWITCH(J1, IM, M, N, I)

CALL RANDS(KV, M, N, DX(I), LO2, M2, N2)

X(2, 2, 1)=1.

Y(2, 2, 1)=0.

X(2, 1, 2)=1.

Y(2, 1, 2)=0.

M2=3

N2=3

N2=2

CALL SLIP(M2, N2)

DC 21 J1=4, IW

K=J1-2

DC 23 J2=2, K

I1=1

IM=1

CALL GENU(J1, J2, IM, AK, AM2)

I1=-1

IM=2

CALL GENU(J2, J1, IM, AK, AM2)

K=J1-1

CALL SLIP(J1, K)
22
23
21
C
C
                                  CALL SLIP(J1,K)
                           OCTPUT CF THE WAKE FIELD

DO 51 I=1,20
WRITE(6,1012) TED(I)
J1=IW
DC 18 I=1,2
WRITE(6,1004)
IF(I.EC.1) WRITE(6,1010)
IF(I.EC.2) WRITE(6,1011)
DC 25 J=2,J1
WRITE(6,1007)
WRITE(6,1008)
K=J-1
N=J-1
WRITE(6,1009) M,N,X(2,M,N),Y(2,M,N),AL(2,M,N),U(2,M,N)
F,V(2,M,N),PSI(2,M,N),G(2,M,N)
F,V(2,M,N),PSI(2,M,N),G(2,M,N)
CCNTINUE
DO 32 M=J,J1
WRITE(6,1009) M,N,X(2,M,N),Y(2,M,N),AL(2,M,N),U(2,M,N)
F,V(2,M,N),PSI(2,M,N),G(2,M,N)
CCNTINUE
WRITE(6,1004)
FCRMAT(1X,** WAKE BAD: *,2I4,2(2X,F8.3),2X,E12.5)
                                   OLTPUT OF THE WAKE FIELD
51
31
 26
 32
 25
18
100c
```

```
2000 FCRMAT(1X,414,2(2X,F8.3))
1004 FCRMAT(1X,/)
1005 FCRMAT(1X,/)
1008 FCRMAT(1X,/)
1008 FCRMAT(1X,/)
1010 FCRMAT(1X,/)
1010 FCRMAT(1X,/)
1010 FCRMAT(1X,/)
1010 FCRMAT(1X,')
1011 FCRMAT(1X,')
1012 FCRMAT(1X,')
1014 FCRMAT(1X,')
1015 FCRMAT(1X,')
1016 FCRMAT(1X,')
1017 FCRMAT(1X,')
1018 FCRMAT(1X,')
1019 FCRMAT(1X,')
1010 FCRMAT(1X,')
1011 FCRMAT(1X,')
1011 FCRMAT(1X,')
1012 FCRMAT(1X,')
1012 FCRMAT(1X,')
1013 FCRMAT(1X,')
1014 FCRMAT(1X,')
1015 FCRMAT(1X,')
1016 FCRMAT(1X,')
1017 FCRMAT(1X,')
1018 FCRMAT(1X,')
1019 FCRMAT(1X,')
1010 FCRMAT(1X,')
1010 FCRMAT(1X,')
1010 FCRMAT(1X,')
1010 FCRMAT(1X,')
1010 FCRMAT(1X,')
1010 FCRMAT(1X,')
1011 FCRMAT(1X,')
1012 FCRMAT(1X,')
1012 FCRMAT(1X,')
1013 FCRMAT(1X,')
1014 FCRMAT(1X,')
1015 FCRMAT(1X,')
1016 FCRMAT(1X,')
1017 FCRMAT(1X,')
1018 FCRMAT(1X,')
1019 FCRMAT(1X,')
1019 FCRMAT(1X,')
1019 FCRMAT(1X,')
1019 FCRMAT(1X,')
1019 FCRMAT(1X,')
1010 FCRMAT(1X,')
                                                                                                           SLIP DOES THE SIMULTANECUS STEPS FROM BOTH SICES OF THE WAKE TO THE SLIP - LINE
                                                                                                  W2=M3-1
N2=N3-1
M4=N3
N13=N3
N13=N3-1
N14=N4
DC3=X(2,M3,N3)-X(2,M13,N13)
D13=SQRT(AL(2,M3,N3)+AL(2,M13,N13)
D13=SQRT(AL(2,M3,N3))+SCRT(AL(2,M13,N13))
D23=AL(2,M3,N3)-AL(2,M13,N13)
A(2,1)=A(2,1)-AM*(AK*003)**2/OAL3*O.5
A(2,1)=A(2,1)-AM*(AK*003)**2/(CAL2*C13)
R(2)=(2,-A(2,1))**U(2,M12,N13)+(CAL2*C13)
R(2)=(2,-A(2,1))**U(2,M12,N13)+(4,/D13-A(2,2))**V(2,M12,N13)
R(2)=R(2)=R(2,1))**U(2,M12,N13)+(4,/D13-A(2,2))**V(2,M12,N13)
A(2,3)=O.
A(2,3)=O.
A(2,4)=C.
DC2=X(2,M3,N3)-X(2,M2,N2)
A(4,1)=-TAN(TED(N3))
A(4,1)=-A(4,1)
A(4,2)=1.
R(4)=O.
DC4=X(2,M4,N4)-X(2,M14,N14)
D14=SCRT(AL(2,M4,N4))+SCRT(AL(2,M14,N14))
A(1,1)=1.+AI*AK*DC3/2.
```

```
A(1,2)=-AI*AK*D03/D13

A(1,3)=-1.-AI*AK*D04/2.

A(1,4)=-AI*AK*D04/D14

RI(1)=PSI(2,M13,N13)-PSI(2,M14,N14)

RI(1)=RI(1)-D04*(U(2,M14,N14)+2.*V(2,M14,N14)/D14)/2.

RI(1)=-(RI(1)+D03*(U(2,M13,N13)-2.*V(2,M13,N13)/C13)/2.)

RI(1)=RI(1)*AI*AK

DA(4=AL(2,M4,N4)+AL(2,M14,N14)

A(3,1)=0.

A(3,2)=0.

A(3,3)=1.+0.5*(AL(2,M4,N4)-AL(2,M14,N14))/DAL4

A(3,3)=A(3,3)+2.*AI*AK*AM*D04/DAL4-AM*(AK*C04)**2/DAL4*C.5

A(3,4)=-(2./D14+AM*(AK*C04)**2/(CAL4*D14))

RI(3)=U(2,M14,N14)*(2.-A(3,3))+V(2,M14,N14)*(-4./C14-A(3,4))

RI(3)=RI(3)+2.*AK*AK*AM*FSI(2,M14,N14)*C04/CAL4
C
                            CCEF1 DETERMINES THE CCEFFICIENTS FOR THE LINSTEADY TEIPEL - CHADWICK SHOCKPOLAR
                          TEIPEL - CHADWICK SHOCKPOLAR

IE=0

CALL FIND(KV,IW,X(2,M,N),Y(2,M,N),M2,N2,IE,I)

CX=C-2.

ALT=AL(1,M2,N2)

IF(IE.EQ.1) ALT=ALO

ANY=(ALT-AM+1.)/(C*AM)+1.

ANY2=ANY*ANY

L1=M+1

L2=N

IF(I.NE.2) GOTO 1

L1=M

L2=N+1

IF(L1.GT.KV .OR. L2.GT.KV) IE=1

IF(IE.EC.1) GOTO 2

S=(Y(2,L1,L2)-Y(2,M,N))/(X(2,L1,L2)-X(2,M,N)))

IF(IE.EC.1) S=2.*I1/(SQRT(ALT)+SQRT(AL(2,M,N)))

S=ATAN(S)

CY=SIN(S)**2

AN=1./(AM*CY*ANY2)

A1=2.*CY*ANY*SIN(2.*S)/C

A2=2.*AK*CY*(1.+AN)/C

A2=CX*CY*(1.-2.*AN/CX)/C+CZ

A4=SIN(2.*S)*(1.+AN)/C

B1=-2.*CY*ANY*(COS(2.*S)+AN)/C

B2=-2.*AK*SIN(S)*COS(S)*(1.+AN)/C

B3=A4

B4=CY+CZ*CX*(1.-2.*AN/CX)/C
   2
```

```
CCMPLEX*16 G,C,X
CCMMON/SCL/ G(4,4),C(4),X(4)
       CCCC
                                                                    SOLVE GIVES THE SOLUTION FOR A COMPLEX SYSTEM OF LINEAR EQUATIONS G*X=C
C SOLVE GIVES THE SOLUTION FOR A COMPLEX SYSTEM OF

LINEAR ECUATIONS G***C

IKK=0

IF(IKK,NE.1) GOTO 2

WRITE (6,1001)

WRITE (6,1000) (G(M,L),L=1,N),C(M)

MA=N-1

DC 10

M=1,MN

CM,J=G(M,J)/G(M,M)

IF(M,So.1) GOTO 14

MS=M-1

DC 13 L=1,MA

C(NTINUE

C(1)=C(1)/G(1,1)

C(1)=C(1)/G(1,1)

C(1)=C(1)-G(1,M)*C(M)/G(I,I)

DC 12 M=1,MA

C(N)=C(N)

DC 11 J=NM

X(N)=C(N)

C(N)=C(N)

C(N)=
                                                                M=K
N=J
RETURN
END
         1
```

```
CCMPLEX*8 L,V,PSI,G,CYCXU
CCMPLEX*8 AI,PU,EI
CCMPLEX*16 EL,RI,ES
REAL*4 TX,TY,XX,YY,X,Y,AL,D1,D2,D3,XP,ALC,P,C3,ALZ
CCMMON/BA/ ALO,AM,ALD,C1,AK,I,IM,TX,TY,IM,MA,IC,ET,XX,YY
COMMON/BC/ T1,T2,T,B,DYCX,D2YDX2,CYCXU,AI,EI,II,IBACK,IE
COMMON/SP/ XS(2,50),YS(2,50),A(2,50,4),R(2,50)
COMMON/SCL/ EL(4,4),RI(4),ES(4)
CCMMON V(2,3,50,20),X(2,3,50,20),P(2,50),U(2,3,50,20),
FPSI(2,3,5C,20),G(2,3,50,20),AL(2,3,50,20),Y(2,3,50,20),
FQ(2,50),CX(2),IN(2,50),IN2(2,50),FX(2,30),PS(2,3C),PL(2,30)
                          OPTIONS:
LC1 = 0 : COMPLETE OUTPUT
=1 : ONLY PRESSURE CISTRIBUTION
=2 : STOP
LC2 = 0 : PITCH
=1 : PLUNGE
LC3 = 1 : CO THE 1. PASSAGE NOT
LC4 = 0 : DO THE 1. PASSAGE NOT
LC4 = 1 : SURFACES ARE ANALYTICALLY GIVEN
=0 : SURFACES ARE PCINTWISE GIVEN
MA = MAXIMUM NUMBER OF FCINTS CN THE SURFACE .LT.20
MAXS = BLADE, WHERE THE FRESSURE DISTRIBUTION IS COMPUTED
IBLA = FIRST BLADE OF FIELD - OUTFUT
                          KV=50

AI=CMPLX(0.,1.)

READ(3,555) LOI,LC2,LC3,LO4,MA,MAXS,IBLA

IF(LOI.EQ.2) GOTO 101

MAXS=MAXS+1

MAX=MA

CALL PRCFIL(LO4,T1,T2,NX)

READ(3,1000) AK,AM,C,OX(1),DX(2)

IF(AM.EC.0) AMY,B,ET,EM

PI=4.*ATAN(1.)

ET=ET*PI/180.

TX1=1.*EM/TAN(ET)

AMY=AM*AM

AV=AM*AM

ALO=AM-1.
  102
  100
                          AP=AM*AM

ALO=AM-1.

C1=C+1.

ALD=ALO**1.5

C3=C1*AM*1.5

IM=1

NT=NX+1

DQ 74 I=1,2

DC 74 J=1,NT

WRITE(1,1026) XS(I,J),YS(I,J),NT
  74
CCC
                             INITIAL FIELD
                           XM=1./SGRT(ALO)
D=4./9.
DC 46 N=1,1C
M=N+1
Y(1,1,M,N)=0.
X(1,1,M,N)=0*(N-1)-2.
DC 47 M=3.KV
X(1,1,M,1)=D*(M-2)/2.-2.
Y(1,1,M,1)=XM*(2.+X(1,1,M,1))
DC 48 N=2,20
DC 48 M=L,KV
K=M-1
   46
   47
                             K=M-1
X(1,1,M,N)=X(1,1,K,N)+G.5+D
```

```
Y(1,1,M,N)=Y(1,1,K,N)+XF*D/2.

DC 44 M=1,KV

DC 44 N=1,20

AL(1,1,M,N)=ALO

PSI(1,1,M,N)=0.

G(1,1,M,N)=0.

U(1,1,M,N)=0.

V(1,1,M,N)=0.
48
                  CCMPUTATION OF THE SPLINE - COEFFICIENTS
                 DC 49 I=1,2

DG 49 K=1,NX

A(I,K,1)=YS(I,K)

A(I,K,3)= R(I,K)

H=XS(I,K+1)-XS(I,K)

A(I,K,2)=(YS(I,K+1)-YS(I,K))/H-F*(R(I,K+1)+2.*R(I,K))/3.

A(I,K,4)=(R(I,K+1)-R(I,K))/(3.*F)
49
                   EI=COS(AMY)+AI*SIN(AMY)
TX=EM/TAN(ET)
TY=EM
55
                 TY=EM
IR=1
IB=2
M1=0
N1=0
WRITE(6,1C21) IM
XFL=IM*TX
YFL=IM*TY
WRITE(1,1026)XPL,YPL,IM
I=0
IE=0
CALL FINC(KV,IM,IB,IR,TX,TY,MI
I=I+1
IF(I.EQ.1) GOTO 67
T2Y=2.*EM
DC 56 J=2,KV
IF(Y(2,1,J,1).GT.T2Y) GCTO 57
CCNTINUE
MAX2=J
                                FINC (KV, IM, IB, IR, TX, TY, M1, N1)
30
567
CCC 676
                  THE STEADY FLOW FIELD
                  I 2=0
I 2=12+1
IF(I 2-GT-30) GOTO 101
HRITE(6,1024) I 2,MA,J1,CX(I)
FCRMAT(IX,I3,'- ITERATION ALONG THE SURFACE',214,2X,E9-3)
MA=MAX
XX=EM*(IE-1)/TAN(ET)
YY=EM*(IB-1)
IF(I.EQ.2..CR. IR.GT-1) PA=MAX2
IF(IR.GT-1) I=IB
I 1=(-1)**(I+1)
MZ=1
IEACK=0
IC=1
KC=0
IF((I*IM).EQ.1 .ANC. IR.GT-1) IQ=C
T=T1
T=T1
1024
                  IF((I*IM).EQ.1 .ANC. I

T=T1

IF(I.EQ.2) T=T2

MM=KV

IF(I.EQ.2. OR. IR.GT.1

DC 5 K=1,20

DC 5 J=K,MM

CALL SWITCH(J,K,M,N,I)

AL(IB,IR,M,N)=0.

M2=M1

N2=N1
                                                   OR. IR.GT.1) MM=20
5
                   THE STEACY BOUNDARY-PROPERTIES BEHING THE SHOCK
                  IE=O
KW=MM+1
```

```
DC 9 J=2,KW
KT=J-1
XF=DX(I)*(J-2)+TX-XX
IF(T.EQ.O. .OR. IQ.EQ.O) XP=TX-XX
IF(IQ.NE.O. .AND. XP.GT.1. .AND. I.EC.1) GGTO 8
IF(J.GT.MM) GOTO 8

LI=1
CALL SWITCH(J,LI,M,N,I)
CALL SHOCK(MM,LO4,EM.IR,IB,M,N,XP,M2,N2,C3)
IF(IR.GT.1 .AND. IB.EC.1 .AND.Y(IE,IR,M,N).EQ.EM) GGTO 63
IF(IB.EQ.2 .AND. IR.EQ.1) GOTO 9
IF(IB.EQ.2 .AND. Y(IB,IR,M,N).EQ.C.) GOTC 63
IF(X(IB,IR,M,N).LE.1.) GOTO 9
ALO=-Y(IB,IF,M,N)/(1.-X(IB,IR,M,N))
ALO=1./(ALO*ALO)
IF(ALO.GT.ALO .ANC. IM.NE.1) GOTO 8
IF(X(IB,IR,M,N).GT.1.) GCTO 8
CCNTINUE
CCNTINUE
CCNTINUE
IF(X(IB,IR,2,1).GE.1.) GCTO 38
IF(I.EQ.2 .AND. X(IB,IR,M,N).GE.1.) KC=2
63
                               ALL OTHER STEPS OF THE STEADY FLOW FIELD
                             KA=KT
DC 1 J=2,KA
J1=J
IF(J.EQ.KA .AND. I.EQ.2 .AND. IR.EC.1) GCTO 15
IF(J.EQ.KA .AND. IR.GT.1) GCTO 15
                             IF(J.EQ.KA .AND. IR.GT.1) GOTO 15

LI=J+1

CALL SWITCH(LI,J,M,N,I)

CALL RANC(LO4,IB, IR,M,N,C3)

IF(X(IB,IR,F,N).GT.TX1) MZ=0

IF(IR.GT.1.AND.I.EQ.1.AND.X(IB,IR,F,N).GT.1..AND.IC.NE.C) MZ=0

IF(MZ.EC.O .AND. IR.GT.1) KC=1

KE=KA+1

DC 3 J2=J,K8

IF(J2.EC.KB) GOTO 3

LI=J-1

CALL SWITCH(J2,LI,K1,K2,I)

IF(AL(IE,IR,K1,K2).GT.AL(IB,IR,F,N)) GOTO 6

CCNTINUE

J2=J2-1

IN(I,J-1)=J2

L=J+2

DC 7

CALL SWITCH(K,J,M,N,I)

CALL GEN(IB,IR,M,N,I)

IF(MZ.EG.O) GOTO 15

IF(J2.EQ.KA) GOTO 1
7
                               CHARACTERISTICS OVERTAKE THE SHOCK
                               J2=M+1
J4=N-1
J5=M
J6=N-1
IF(I.NE.2) GOTO 12
J2=M-1
J4=M+1
                             J4=N+1

J5=M-1

J6=N

D1=Y(IB, IR, J3, J4)-Y(IB, IR, J5, J6)

D1=D1/(X(IB, IR, J3, J4)-X(IB, IR, J5, J6))

D2=I1/SQRT(AL(IB, IR, M, N))

K1=M+1

K2=J

IF(I.NE.2) GOTO 14

K1=J

K2=N+1

AL(IB, IR, K1, K2)=AL(IB, IR, M, N)
12
```

```
D3=-D2*X(IB,IR,M,N)+D1*X(IB,IR,J5,J6)
D3=D3-Y(IB,IR,J5,J6)+Y(IB,IR,M,N)
X(IB,IR,K1,K2)=D3/(D1-C2)
Y(IB,IR,K1,K2)=D2*(X(IE,IR,K1,K2)-X(IB,IR,M,N))+Y(IB,IR,M,N)
IF(I.EC.2) J3=J4
DC 11 K=J3,KA
IF(K.EQ.MM) GOTO 11
LI=K+1
CALL SWITCH(LI.J.M.)
                               LI=J-1

CALL SWITCH(K,LI,K1,K2,I)

X(IB,IR,M,N)=X(IB,IR,K1,K2)

Y(IB,IR,M,N)=Y(IB,IR,K1,K2)

AL(IB,IR,M,N)=AL(IB,IR,K1,K2)

CCNTINUE

IF(KA-LT-MM) KA=KA+1

IF(MA-GT-IN(I,J-1)) MA=PA+1

CCNTINUE
  11
100015
                                  AUTOMATIC STEP-SIZE CONTROL
                                 J2=J1
IF(KC.EC.1) J2=KT
IF(T.EQ.0...OR.IC.EC.0) GOTO 28
IF(J2.EQ.MA) GOTO 28
IF(J2.ET.MA) GOTO 68
DIF=(J2*1.)/(MA*1.)
IF(DIF.GE.0.85) GOTO 65
GCTO 72
DIF=(J2*1.)/(MA*1.)
IF(DIF.LE.1.15) GOTO 65
DX(I)=DX(I)*DIF
GCTO 72
IF(J2.GT.MA) DX(I)=DX(I)*1.05
IF(J2.LT.MA) DX(I)=DX(I)*0.96
CCNTINUE
  68
  72
   69
                                   CCNTINUE
IF(J2.NE.MA) GOTO 16
   73
  28
                                 CCNTINUE
IF(IR.EG.1 .AND. I.EC.1) GOTO 59
IF(KC.NE.0) GOTO 59
                               ACDITIONAL FOINTS FOR THE FIELDS IN THE PASSAGE

DC 33 LY=2,20
LI=LY-1
CALL SWITCH(LY,LI,M,N,I)
IF(AL(IB,IR,M,N).EQ.O.) GOTO 35
WRITE(6,1009) M,N,X(IB,IR,M,N),Y(IB,IR,M,N),AL(IE,IR,M,N)
CONTINUE
KC=LY
LI=KD-1
LZ=C
CALL SWITCH(LI,LZ,M,N,I)
LE=I
CALL SWITCH(KD,LE,IIB,IZ,I)
X(IB,IR,LI,L,2)=X(IB,IR,M,N)
LZ=KD-2
CALL SWITCH(LI,LZ,M,N,I)
LZ=KD-2
CALL SWITCH(LI,LZ,L3,L4,I)
Y1=I1*(X(IB,IR,LI,L2)-X(IB,IR,M,N)I/SQRT(AL(IB,IR,M,N))+EM*(IE-2)
X1=X(IB,IR,LI,L2)=X(IB,IR,M,N)/(X(IB,IF,L2,L4)-X(IB,IR,M,N))
D1=D1*(Y(IB,IR,L3,L4)-Y(IB,IR,M,N))/(X(IB,IF,L2,L4)-X(IB,IR,M,N))
D1=D1*(Y(IB,IR,L3,L4)-Y(IB,IR,M,N))/(X(IB,IF,L2,L4)-X(IB,IR,M,N))
D2=D1*(X(IB,IC,M,N)+X(IB,IR,M,N))/(X(IB,IF,L2,L4)-X(IB,IR,M,N))
D2=D1*(X(IB,IC,M,N)+X(IB,IR,L1,L2))
X(IB,IR,M,N)=X(IB,IR,L1,L2)
Y(IB,IR,M,N)=X(IB,IR,L1,L2)
Y(IB,IR,M,N)=EM*(IB-1)
                                  ACDITIONAL FOINTS FOR THE FIELDS IN THE FASSAGE
   43
```

```
,IR,M,N), Y(18, IR, M, N), X(18, IR, L1, L2),
C
65
000005
                     THE UNSTEADY FLOW FIELD THE UNSTEADY BOUNDARY PROPERTIES BEHING THE SHOCK
                     KA=KT
LZ=2
LI=1
                    LI=1
CALL SWITCH (LZ, LI, I7, I8, I)
XFL=X(IB, IR, I7, I8)+(IM-1)*EM/TAN(ET)
YFL=Y(IB, IR, I7, I8)+(IM-1)*EM
WRITE(I, 1026) XPL, YPL, IP
CALL SWITCH(KT, LI, I18, I)
XFL=X(IB, IR, I7, I8)+(IM-1)*EM/TAN(ET)
YPL=Y(IB, IR, I7, I8)+(IM-1)*EM
WRITE(1, 1026) XPL, YPL, IP
WRITE(6, 1025) IR
FCRMAT(1X, I3, ** UNSTEACY FIELD*)
IE=0
DC 18 J=2, KA
CALL SWITCH(J, LI, M, N, I)
IF(X(IB, IR, M, N, I), LT.00.) CCTO 18
CALL RANCS(MM, IB, IR, M, N, LOZ, AMY, MZ, N2)
CCNTINUE
IF(I.EC.2) X(IB, IR, 1, 2)=TX
IF(I.EC.2) X(IB, IR, 1, 2)=TX
IF(I.EC.2) X(IB, IR, 1, 2)=TX
IF(I.EC.2) X(IB, IR, 1, 2)=TY
 1025
 18
 CCC
                                   OTHER STEPS OF THE UNSTEADY FIELD
                       IF(J1.EC.2) GOTO 38
                     IF(J1.EC.2) GOTO 38

J=J1-1
DC 20 J=2,J5

LI=J+1
CALL SWITCH(LI,J,M,N,I)
CALL RANCB(IB,IR,M,N,LC2)
IF(J.EG.J5 .AND. I.EQ.2 .AND. IR.EG.1) GCTC 2C

IF(J.EG.J5 .AND. IR.GT.1 .AND. MZ.NE.0) GOTO 2C

L=J+2
J2=IN(I,J-1)
DC 22 K=L,J2
CALL SWITCH(K,J,M,N,I)
CALL GENU(IM,IB,IR,M,N,I,AK,AM)
IF(J2.EC.KA) GOTO 20
K1=M+1
K2=J
 22
                         2=J

2=K1

5(I.NE.2) GOTO 24

1=J

2=N+1
                          (I,K3)=X(IB,IR,M,N)
ALL RANDS(MM,IB,IR,K1,K2,L02,AMY,M2,N2)
 24
                           3=J2+1

5 29 K=J3, KA

F(K.EQ.MM) GOTO 20

I=K+1
                                           WITCH(LI,J,M,N,I)
                     LI=J-1

CALL SWITCH(K,LI,K1,K2,I)

U(IB,IR,M,N)=U(IB,IR,K1,K2)

V(IB,IR,P,N)=V(IB,IR,K1,K2)

G(IB,IR,M,N)=G(IB,IR,K1,K2)

PSI(IB,IR,P,N)=PSI(IB,IR,K1,K2)
 29
```

```
IF(J2.LT.KA .AND. KA.LT.MM) KA=KA+1
                       CUTPUT STEADY AND UNSTEACY FIELD
                     IF(IM.LT.IBLA) GOTO 38
IF(LOI.NE.O) GOTO 38
WFITE(7,1004)
IF(I.EQ.I.AND. IR.EQ.I) WRITE(7,1002) IM
IF(I.EC.2.OR. IR.GT.I) WRITE(7,1003) IR, IM
WRITE(7,1001) AK,AX,C,B,CX(I),T
DC 25 J=2,J1
LL=IN(I,K)+1
LL=IN(I,K)
WRITE(7,1007)
WRITE(7,1007)
WRITE(7,1007)
WRITE(7,1008)
N=K
C
                   IF(I.NE.2) GOTO 26

M=K

OC 31 N=J,LL

WRITE(7.1005) M.N.X(IB.IR.M.N),Y(IB.IR.M.N),G(IB.IR.M.N)

-.U(IB.IR.M.N),V(IB.IR.M.N),PSI(IB.IR.M.N),G(IB.IR.M.N)

CCNTINUE

OC 32 M=J,LL

WRITE(7.1005) M.N.X(IB.IR.M.N),Y(IB.IR.M.N),G(IB.IR.M.N)

F.U(IB.IR.M.N),V(IB.IR.M.N),PSI(IB.IR.M.N),G(IB.IR.M.N)

CCNTINUE

WRITE(7.1005)

CCNTINUE

WFITE(7.1004)

CCNTINUE

IF(I.EQ.1) IUS=K-1

IF(I.EQ.1) IUS=K-1

IF(I.EQ.1) GOTO 30

IF(IM.EC.1 .AND. LO3.EC.C) GOTO 62
                          F(I.NE.2) GOTO 26
 31
 26
 32
 17
 38
00000
                      CHANGING THE COUNTERS FOR THE COMPUTATION OF THE FIELDS BEHIND THE REFLECTED SPOCKS
                       IF=IR+1
IF(IR-GE-4) GOTO
IF(KC-NE-0) GOTO
WFITE(6,1023) IR
                      LI=1
CALL SWITCH(KT,LI,M,N,I)
CALL SWITCH(KT,LI,MI,NI,I)
TX=X(IB,IR-1,M,N)
TY=Y(IB,IR-1,M,N)
IF(IR.EG.2 .UR. IR.EG.4) IB=1
IF(IR.EG.3 .UR. IR.EG.5) IB=2
GCTO 67
CCNTINUE
IF(KC.EG.0) GOTO 77
CCMPUTATION OF THE BLACE - WAKE IN
SLBROUTINE *WAKE*
 62
0000077
                      CALL WAKE
NEXT BLADE
IM=IM+1
IF(IM-LE.MAXS) GOTO 58
                       IR=IR-I
CALL PRESS(AMY, IR, EM, LG2, AX, C)
IM=0
WRITE(1, 1026) XPL, YPL, IM
GCTO 100
C
C
C
58
                       TRANSFORMATION TO PROFIL 1
                      XX=EM/TAN(ET)
DC 70 IL=1,3
MF=20
```

```
IF(IL.EC.1) MM=KV

OC 51 N=1,20

X1.1L,P.N=X(2,IL,P.N)-EM
P$(1,IL,M.N)=P$(2,IL,P.N)

U1.1L,M.N)=P$(2,IL,P.N)

V(1,IL,M.N)=P$(2,IL,P.N)

V(1,IL,M.N)=P$(1(2,IL,P.N)

OC TINUE

101 IM=1

102 FCRMAT(4I1,3I2)

103 FCRMAT(4I1,3I2)

100 FCRMAT(11,7,4,2,IL,P.N)

100 FCRMAT(11,7,4,2,IL,P.N)

100 FCRMAT(11,7,4,2,IL,P.N)

100 FCRMAT(11,7,4,2,IL,P.N)

100 FCRMAT(11,FCR THE UPPER SURFACE OF THE FLOW FIELC*,

100 FCRMAT(11,FCR THE UPPER SURFACE OF THE FLOW FIELC*,

100 FCRMAT(11,FCR THE UPPER SURFACE OF THE FLOW FIELC*,

100 FCRMAT(11,FCR THE UPPER SURFACE OF THE FLOW FIELC*,

100 FCRMAT(11,FCR THE UPPER SURFACE OF THE FLOW FIELC*,

100 FCRMAT(11,FCR THE UPPER SURFACE OF THE FLOW FIELC*,

100 FCRMAT(11,FCR THE UPPER SURFACE OF THE FLOW FIELC*,

100 FCRMAT(11,FCR THE UPPER TIES AT THE MESTICIN'S CF THE FLOW FIELC*,

100 FCRMAT(11,FCR THE UPPER TIES AT THE MESTICIN'S CF THE FLOW FIELC*,

100 FCRMAT(11,FCR THE UPPER TIES AT THE MESTICIN'S CF THE FLOW FIELC*,

100 FCRMAT(12,FCR THE UPPER TIES AT THE MESTICIN'S CF THE FLOW FIELC*,

100 FCRMAT(12,FCR THE UPPER TIES AT THE MESTICIN'S CF THE FLOW FIELC*,

100 FCRMAT(12,FCR THE UPPER TIES AT THE MESTICIN'S CF THE FLOW FIELC*,

100 FCRMAT(12,FCR THE UPPER TIES AT THE MESTICIN'S CF THE FLOW FIELC*,

100 FCRMAT(12,FCR THE UPPER TIES AT THE MESTICIN'S CF THE FLOW FIELC*,

100 FCRMAT(12,FCR THE UPPER TIES AT THE MESTICIN'S CF THE FLOW FIELC*,

100 FCRMAT(12,FCR THE UPPER TIES AT THE MESTICIN'S CF THE FLOW FIELC*,

100 FCRMAT(12,FCR THE UPPER TIES AT THE MESTICIN'S CF THE FLOW FIELC*,

100 FCRMAT(12,FCR THE UPPER TIES AT THE MESTICIN'S CF THE FLOW FIELC*,

100 FCRMAT(12,FCR THE UPPER TIES AT THE MESTICIN'S CF THE FLOW FIELC*,

100 FCRMAT(12,FCR THE UPPER TIES AT THE MESTICIN'S CF THE FLOW FIELC*,

100 FCRMAT(12,FCR THE UPPER TIES AT THE MESTICIN'S CF THE FLOW FIELC*,

100 FCRMAT(12,FCR THE UPPER TIES AT THE MESTICIN'S CF THE FLOW FIELC*,

100 FCRMAT(12,FCR THE UPPER TIES AT THE MESTICIN'S CF THE FLOW FIELC*,

100 FCRMAT(12,FCR THE UPPER TIES AT THE MESTICIN'S CF THE FLOW FIELC*,

100 FCRMA
        CCC
                                                                                                                                                                                                                                                                                                                                              - SURFACES
                                                                        FFEPARATION OF THE
                                                                  READ(3,1000) T1,T2,NT,IKP
IF(L04.E(.0) READ(3,1001) SP

OC 8 J=1,2
IZ=0
N=NT
T=T1
IF(J.EQ.2) T=T2
IF(L04.NE.0) GOTO 12
DC 9 K=1,4
M=(K-1)*5+1
READ(3,1001) XS(J,M),XS(J,M+1),XS(J,M+2),XS(J,M+2),XS(J,M+4)
READ(3,1001) YS(J,M),YS(J,M+1),YS(J,M+2),YS(J,M+2),YS(J,M+4)
IF(YS(J,M+4).EQ.100.) GCTO 7
CCNTINUE
DC 2 LL=1,N
XS(J,LL)=XS(J,LL)/SP
YS(J,LL)=XS(J,LL)/SP
         57
        2
        12
NOUN MUCU
                                                                     TC ENTER THIS PART, THE SURFACES SHOULD ALREADY BE GIVEN PCINTWISE CONTINUE
                                                                        INTERPOLATION THROUGH CLBIC SPLINES
                                                                      IF (T.NE.0.
DC 52 I=1,N
A(J,I,I)=0.
                                                                                                                                                                                          .OR. LG4.E6.01 GOTO 51
        52
```

```
51
                                             =XS(J; I)
10
C
C
                          TRIX OF COEFFICIENTS AND RIGHT-HAND SIDES
                K=N-2

CC 25 I=1,K

A(J,I,1)=A(J,I+1,3)-A(J,I,3)

A(J,I,3)=A(J,I+2,3)-A(J,I+1,3)

A(J,I,2)=2.*(A(J,I,1)+A(J,I,3)

A(J,I,4)=3.*(A(J,I+2,4)-A(J,I,4))/A(J,I,4)

F3.*(A(J,I+1,4)-A(J,I,4))/A(J,I,1,1)

CCNTINUE

A(J,I,1)=0.0

A(J,N-2,3)=0.0
 25
CCC
                    THE STEP OF GAUSS
                         N-3
30 I=1,K
35 M=3,4
J,I,M)=A(J,I,M)+A(J,I+1,1)/A
J,I,1)=0.0
J,I,2)=A(J,I+1,1)
J,I+1,2)=A(J,I+1,2)-A(J,I,3)
J,I+1,4)=A(J,I+1,4)-A(J,I,4)
 35
30
                    SCLUTION
                       (J,1,1)=0.
(J,N,1)=0.
(J,N-2,1)=0.0
                  DC 40 1-2, c

K=N-1

M=K+1

A(J,M,1)=(A(J,K,4)-A(J,K,3)*A(J,M+1,1))/A(J,K,2)

CCNTINUE

DC 11 M=1,N

R(J,M)=A(J,M,1)

IF(IZ.EQ-1) GOTO 8

KV=N-1

DC 49 K=1,KV
 40
53
 11
                                               ,KV
YS(J,K)
R(J,K)
1)-XS(J,K)
(YS(J,K+1)-YS(J,K))/H-+*(F(J,K+1)+2.*R(J,K))/3.
(F(J,K+1)-R(J,K))/(3.*+)
 49
                             1./49.
1./49.
4 M=1.50
X*(M-1)
5 K=2.50
                    DC 5 K=2,50
I=K-1
IF(XS(J,K).GE.X) GOTO 6
CONTINUE
H=X-XS(J,I)
Y=A(J,I,1)+A(J,I,2)*H+A(J,I,3)*H*++A(J,I,4)*H*H*+
R(J,M)=X
YS(J,M)=Y
CONTINUE
TTTT
             YS(JiNUE

IZ=1

N=50

DC 13 I=1,50

XS(J,I)=R(J,I)

GCTO 3

CONTINUE

IF(IKP.EQ.2) GOTO 17

DC 15 M=1.N

WRITE(6,1002) XS(1,M),YS(1,M),XS(2,M),YS(2,M)

N=N-1

FCRMAT(2F10.5,212)

136
 13
 8
 15
17
1000
```

```
FIND LOOKS FOR THE MESHINDEX, RESPONSIBEL FOR THE POINT XS:YS
                                                                    NX7=10000*XS
NY7=10000*YS
                                                               IC=I

IF(IR-EC-1) GGTO 26

IC=IR-1

IF(IB-EG-2) IC=1

IF(IB-EG-1) IC=2

I1=(-1)**(IC+1)

IF(IE-EG-1) GGTO 19

IKK=0

IF(I.EQ.2) IKK=1

IF(IKK-EG-1) WRITE(6,1004)

FCRMAT(IX, 'FIND-ENTRY')

KA=0

L1=1
   26
   C
    1004
                                                             FCRMAT(1x, *FIND-ENTRY*)

KA=0
L1=1
L2=1
IK=21
IK=21
IF(IKK.EQ.1) WRITE(6,1003) IK, M,N,NX7,NY7,NX9,NY9
DC 22 I2=L1,20
CALL SWITCH(L3,I2,M,N,IC)
IK=22
IF(IKK.EQ.1) WRITE(6,1003) IK, M,N,NX8,NY8,NX9,NY9
IF(IC.EQ.1 -AND. L3.GE.KV) GOTO 13
J1=M+1
J2=N+1
KM=J1-J2
IK=24
IF(IK.EQ.1) WRITE(6,1003) IK,M,N,NX8,NY8,NX9,NY9
IF(IC.EC.2) KM=J2-J1
NX8=10000*X(IC,ID,M,N)
NX9=10000*X(IC,ID,M,N)
NX9=10000*X(IC,ID,M,N)
NY9=10000*X(IC,ID,M,N)
NY9=10000*X(IC,ID,M,N)
NY9=10000*X(IC,ID,M,N)
IF(NX7.GE.NX8.AND. NX7.LT.NX9) GCTC 23
IF(IC.EC.2 -AND. KA.EC.J2) KA=J1
IF(IC.EC.2 -AND. KA.NE.J1) GOTO 22
IF(IC.EQ.2 -AND. KA.NE.J1) GOTO 23
IF(IC.EQ.2 -AND. KA.NE.J1) GOTO 23
IF(IC.EQ.2 -AND. KA.NE.J1) GOTO 13
IF(IC.EQ.2 -AND. CA.NE.J2 -AND. LA.NE.J1) IF(IC.EQ.2 -AND. KA.NE.J1) IF(IC.EQ.2 -AND. KA.NE.J1) GOTO 22
IF(IC.EQ.2 -AND. KA.NE.J1) GOTO 13
IF(IC.EQ.2 -AND. CA.NE.J1) IF(IC.EQ.2 -AND. KA.NE.J1) IF(IC.EQ.
    21
   24
  22
```

```
GCTO 21
I K=28
IF(IKK.EC.1) WRITE(6,10C3) IK, M, N, NX8, NYE, NX9, NY9
NY8=10000*Y(IC, ID, M, J2)
IF(NY7.GT.NY8) GOTO 1
IF(M.EQ.1) GOTO 13
L2=L2+1
L1=M-1
GCTO 21
IK=1
IF(IKK.EQ.1) WRITE(6,1002) IK, M, N, NX8, NY8, NX9, NY9
IF(M.GE.KV) GOTO 13
J=M+1
K=N+1
NX8=10000*X(IC, ID, M, N)
NY8=10000*X(IC, ID, M, N)
IF(NX7.EQ.NX8.AND.NY7.EQ.NY8) GCTO 14
IF(IR.GT.1) GOTO 29
 28
1
                                                          GRIENTATION IN THE FIELD GVER THE UPPER PLADE
IF(NX7.NE.NX8) GGTD 5
IF(NX7.GT.NY8 .AND. N.EC.1) GGTD 13
IF(NY7.GT.NY8) N=N-1
IF(NY7.LT.NY8) M=M-1
GCTD 1
IF(ST.KY) GGTD 7
IF(IKK.EQ.1) WRITE(6,1CC2) IK,M,N,NX8,NY8,NX9,NY9
I=M+2
IF(IKK.EQ.1) WRITE(6,1CC2) IK,M,N,NX8,NY8,NX9,NY9
I=M+2
IF(IKK.EQ.1) WRITE(6,1CC2) IK,M,N,NX8,NY8,NX9,NY9
IF(NX8.EQ.NX9 .AND. NY8.EQ.NY9) GCTC 6
GGTD 7
IF(IKK.EQ.1) WRITE(6,1CO2) IK,M,N,NX8,NY8,NX9,NY9
IF(NX8.EQ.NX9 .AND. NY8.EQ.NY9) GCTC 6
IF(IKK.EQ.1) WRITE(6,1CO2) IK,M,N,NX8,NY8,NX9,NY9
IF(IC,ID,M,N)-Y(IC,ID,M,N)/(X,IC,ID,M,K)-X(IC,ID,M,N))
D2=(Y(IC,ID,M,N)-Y(IC,ID,M,N)/(X,IC,ID,M,K)-X(IC,ID,M,K))
D3=(Y(IC,ID,M,K)-Y(IC,ID,M,K))/(X,IC,ID,M,K)-X(IC,ID,M,K))
IF(IC,ID,M,K)-Y(IC,ID,M,K))/(X,IC,ID,M,K)-X(IC,ID,M,K))
IF(IC,ID,M,N)-Y(IC,ID,M,N)/(X,IC,ID,M,K)-X(IC,ID,M,N))
IF(IC,ID,M,N)-Y(IC,ID,M,N)/(X,IC,ID,M,K)-X(IC,ID,M,N))
D5=(Y(IC,ID,M,K)-Y(IC,ID,M,N))/(X,IC,ID,M,K)-X(IC,ID,M,N))
IF(IC,ID,M,N)-Y(IC,ID,M,N)/(X,IC,ID,M,K)-X(IC,ID,M,N))
D5=(Y(IC,ID,M,K)-Y(IC,IC,M,N))/(X,IC,ID,M,K)-X(IC,ID,M,N))
D5=(Y(IC,ID,M,K)-Y(IC,IC,M,N))/(X,IC,ID,M,K)-X(IC,ID,M,N))
D5=(Y(IC,ID,M,K)-Y(IC,IC,M,N))/(X,IC,ID,M,K)-X(IC,ID,M,N))
D5=(Y(IC,ID,M,K)-Y(IC,IC,M,N))/(X,IC,ID,M,K)-X(IC,ID,M,N))
D5=(Y(IC,ID,M,K)-Y(IC,IC,M,N))/(X,IC,ID,M,K)-X(IC,ID,M,N))
D5=(Y(IC,ID,M,K)-Y(IC,IC,M,N))/(X,IC,ID,M,K)-X(IC,ID,M,N))
D5=(Y(IC,ID,M,N)-Y(IC,IC,M,N))/(X,IC,ID,M,K)-X(IC,ID,M,N))
D5=(Y(IC,ID,M,N)-Y(IC,IC,M,N))/(X,IC,ID,M,K)-X(IC,ID,M,N))
D5=(Y(IC,ID,M,N)-Y(IC,IC,M,N))/(X,IC,ID,M,K)-X(IC,ID,M,N))
D5=(Y(IC,ID,M,N)-Y(IC,IC,M,N))/(X,IC,ID,M,K)-X(IC,ID,M,N))
D5=(Y(IC,ID,M,N)-Y(IC,IC,M,N))/(X,IC,ID,M,K)-X(IC,ID,M,N))
D5=(Y(IC,ID,M,N)-Y(IC,IC,M,N))/(X,IC,ID,M,K)-X(IC,ID,M,N))
D5=(Y(IC,ID,M,N)-Y(IC,ID,M,N))/(X,IC,ID,M,K)-X(IC,ID,M,N))
D5=(Y(IC,ID,M,N)-Y(IC,ID,M,N))/(X,IC,ID,M,N))
D5=(Y(IC,ID,M,N)-Y(IC,ID,M,N))/(X,IC,ID,M,N))
D5=(Y(IC,ID,M,N)-Y(IC,ID,M,N))/(X,IC,ID,M,N))
D5=(Y(IC,ID,M,N)-Y(IC,ID,M,N))/(X,IC,ID,M,N))
D5=(Y(IC,ID,M,N)-Y(IC,ID,M,N))/(X,IC,ID,M,N))
D5=(Y(IC,ID,M,N)-Y(IC,ID,M,N)-X(IC,ID,M,N))
D5=(Y(IC,ID,M,N)-X(IC,ID,M,N)-X
                                                                    CRIENTATION IN THE FIELD OVER THE UPPER PLADE
  5
   7
                                                                      N=N-1
M=M-1
GCTO
                                                                       GCTO 3
[F(D12.GT.D1 .AND. N.EC.1) GOTC 13
   2
                                                                 Ik=2
IF(IKK.EC.1) WRITE(6,1003) IK,M,N,Nxe,Nye,Nx9,Ny9
IF(D12.GT.D1) N=N-1
IF(D12.LT.D2) M=M-1
IBACK=IBACK+1
J=M+1
IF(AL(IC,ID,M,N).NE.AL(IC,ID,J,N)) GGTG 13
GCTG 1
IK=8
NX8=10000*X(IC,ID,J,K)
NY8=100CC*Y(IC,ID,J,K)
NY8=100CC*Y(IC,ID,J,K)
IF(IKK.EG.1) WRITE(6,1003) IK,M,N,NX8,NYE,NX9,NY9
   3
```

```
IF(NX7.EQ.NX8.AND. NY7.EQ.NY8) GCTC 12
IF(NX7.NE.NX8) GOTO 15
IF(NY7.GT.NY8) M=M+1
IF(NY7.LT.NY8) N=N+1
GCTO 1
IK=12
IF(IKK.EQ.1) WRITE(6.10C3) IK,M,N,NX8,NY8,NX9,NY9
12
                                 iF(ikk.Eq.1) write(6.10c3) ik, m, n, n, xe, n, ye, n, x9, n, ys
M=J
N=K
GCTO 14
ik=15
iF(ikk.Eq.1) write(6,10c3) ik, m, n, n, xe, n, ye, n, x9, n, ys
iF(ixt.eq.n) goto 16
iF(k.Eq.m) goto 10
D3= (Y(ic,iD,J,k)-Y(ic,ic,m,k))/(x(ic,iD,J,k)-x(ic,iD,m,k))
D4= (Y(ic,iD,J,k)-Y(ic,ic,J,n))/(x(ic,iD,J,k)-x(ic,iD,J,n))
D34= (Y(ic,iD,J,k)-Y(ic,iC,J,n))/(x(ic,iD,J,k)-x(ic,iD,J,n))
D35= (Y(ic,iD,J,k)-Y(ic,iC,J,n))/(x(ic,iD,J,k)-x(ic,iD,J,n))
D35= (Y(ic,iD,J,k)-Y(ic,iC,J,n))/(x(ic,iD,J,k)-x(ic,iD,J,n))
D35= (Y(ic,iD,J,k)-Y(ic,iC,J,n))/(x(ic,iD,J,k)-x(ic,iD,J,n))
D35= (Y(ic,iD,J,k)-Y(ic,iC,J,n))/(x(ic,iD,J,k)-x(ic,iD,J,n))
D35= (Y(ic,iD,J,k)-Y(ic,iC,J,n))/(x(ic,iD,J,k)-x(ic,iD,J,n))
D36= (Y(ic,iD,J,k)-Y(ic,iC,J,n))/(x(ic,iD,J,k)-x(ic,iD,J,n))
D36= (Y(ic,iD,J,k)-Y(ic,iC,J,n))/(x(ic,iD,J,k)-x(ic,iD,J,n))
D36= (Y(ic,iD,J,k)-Y(ic,iC,J,n))/(x(ic,iD,J,k)-x(ic,iD,J,n))
D36= (Y(ic,iD,J,k)-Y(ic,iC,J,n))/(x(ic,iD,J,k)-x(ic,iD,J,n))
D4= (Y(ic,iD,J,k)-Y(ic,iD,J,k)-x(ic,iD,J,k)-x(ic,iD,J,n))
D4= (Y(ic,iD,J,k)-Y(ic,iD,J,k)-x(ic,iD,J,k)-x(ic,iD,J,n))
D4= (Y(ic,iD,J,k)-Y(ic,iD,J,k)-x(ic,iD,J,k)-x(ic,iD,J,k)-x(ic,iD,J,k)-x(ic,iD,J,k)
D4= (Y(ic,iD,J,k)-Y(ic,iD,J,k)-x(ic,iD,J,k)-x(ic,iD,J,k)-x(ic,iD,J,k)-x(ic,iD,J,k)-x(ic,iD,J,k)-x(ic,iD,J,k)-x(ic,iD,J,k)-x(ic,iD,J,k)-x(ic,iD,J,k)-x(ic,iD,J,k)-x(ic,iD,J,k)-x(ic,iD,J,k)-x(ic,iD,J,k)-x(ic,iD,J,k)-x(ic,iD,J,k)-x(ic,iD,J,k)-x(ic,iD,J,k)-x(ic,iD,J,k)-x(ic,iD,J,k)-x(ic,iD,J,k)-x(ic,iD,J,k)-x(ic,iD,J,k)-x(ic,iD,J,k)-x(ic,iD,J,k)-x(ic,iD,J,k)-x(ic,iD,J,k)-x(ic,iD,J,k)-x(ic,iD,J,k)-x(ic,iD,J,k)-x(ic,iD,J,k)-x(ic,iD,J,k)-
15
10
17
                                         16
                                     GCTO 1
M=M+1
9
                                     IF(IKK.EC.1) WRITE(6,1003) IK,M,N,NXE,NYE,NX9,NYS
                                   N=N+1
GCTO 1
IE=1
13
                                     GOTO 19
CC 29
                                     ORIENTATION IN THE PASSAGE
                                    IK=29
IF(IKK.EG.1) WRITE(6,1003) IK,M,N,NX8,NY8,NX9,NY9
L1=M+1
L2=N
L3=M
                                         4=N+1
F(IC.NE.2) GOTO 30
1=M
2=N+1
30
                                             X8=10000*X(IC, ID, L5, L6)
X9=10000*X(IC, ID, L1, L2)
Y8=10000*Y(IC, ID, L5, L6)
Y9=10000*Y(IC, ID, L1, L2)
                                   33
```

```
GCTO 1
(3=(Y(IC,IC,L5,L6)-Y(IC,ID,L1,L2))/(X(IC,ID,L5,L6)-X(IC,ID,L1,L2))
(3=D3*I1
(K=32
(F(IKK.EQ.1) WRITE(6,1003) IK,M,N,NX8,NY8,NX9,NY9
(F(L3.EC.L4) GOTO 34
(4=(Y(IC,ID,L5,L6)-Y(IC,ID,L3,L4))/(X(IC,IC,L5,L6)-X(IC,ID,L3,L4))
(4=D4*I1
(F(L3.EC.L4) D4=0.
(34=D4*I1
(F(L3.EC.L4) D4=0.
(34=I1*(YS-Y(IC,ID,L5,L6))/(XS-X(IC,ID,L5,L6))
(F(D34.GT.C3.AND.D34.LT.D4) GOTC 14
(Y8=100CC*Y(IC,ID,L5,L6)
(F(NY7.GT.NY8) M=M+1)
(F(NY7.LT.NY8) N=N+1)
(F(NY7.LT.NY8) N=N+1)
(FL5)
32
34
                                       31
                                         M=L3
N=L4
IK=35
IF(IKK.EQ.1) WRITE(6,1003) IK,M,N,NX8,NY8,NX9,NYS
36
                                          IK=14
IF(IKK.EQ.1) WRITE(6,1003) IK, M, N, NX8, NY8, NX9, NY9
14
                                         CENTROL - STEP
                                         K1=M+1
K2=N
K3=M
K4=N+1
IF(IC.NE.2) GOTO 35
35
                                                                                 000*X(IC,ID,M,N)
000*X(IC,ID,M,N)
000*Y(IC,ID,M,N)
000*Y(IC,ID,M,N)
000*Y(IC,ID,J1,J2)
000*Y(IC,ID,M,N)
000*Y(IC,ID,M,N)
000*Y(IC,ID,M,N)
000*Y(IC,ID,M,N)
000*Y(IC,ID,M,N)
0000*Y(IC,ID,M,N)
0000*Y(IC
                                                                                            ,1000) M,N,NX7,KY7,NX8,NY8,D1,D2,C3,C4,C12,C34
18
20
                                                                                                                                   D, K1, K2)-Y(IC, ID, M, N))/(X(IC, IC, K1, K2)-X(IC, ID, M, N))
GDTO 11
D, K3, K4)-Y(IC, ID, M, N))/(X(IC, ID, K3, K4)-X(IC, ID, M, N))
11
 25
                                                                                                                                                                                                                                    ID, K1, K2))
ID, K1, K2))
(X (IC, IC, J)
GOTC 18
X9) GCTC 18
X9, GCTC 1
I2, 4(1x, I8
                                                                                                                                                                                                                                                                                                    J1, J21-XS1
                                                                                                                                                                                                                                                                              TČ 18
X, I 8), 1C(1X, F8.3))
```

```
CCMPUTATION OF THE FIELD NEAR BEHIND THE SHOCK
          IKK=0
IF(I.FQ.2) IKK=1
IF(IKK.EQ.1) WRITE(6,1003)
 C
         IF(IKK.EQ.1) WRITE(6,1003)
MAX=1
IC=1
IC=1
IF(IR.EC.1) GOTO 19
IC=IR-1
IF(IB.EQ.1) IC=2
IF(IB.EC.2) IC=1
Dx1=DX(I)
IF(TEQ.00..OR. IG.EQ.0) GOTO 17
GCTO 21
D=EM*2.*SQRT(ALO)
IF(IB.EQ.2..AND. T.EQ.0..AND. I.EC.1) C=1.05
DX1=D/(MA-1)
IK=21
IF(IKK.EQ.1) WRITE(6,1000) M,N,IE,IK,D,DX1,XO,TX,TY
J=M-1
K=N
 19
 17
 21
          K=N
          I F(I.NE.2) GOTO 5
          J=M
         J=F
K=N-1
L=N
Call ccnst1(LG4,IR,IE,ALZ,C3,X0)
 5
```

```
CALL BCUND(LO4, IB, IR, IN, I, XO)
AL(IB, IR, M, N) = (ALZ-C3*CYCX*I1)**(2./3.)
IF(L.NE.2) GOTO 3
WRITE(6,1001) M,N,AL(IC, ID, M2, N2),M2,N2
GOTO 6
                                         WRITE (6,1001) M,N,AL(IC,ID,MZ,NZ),PZ,NZ

GOTO 6

IK=3

IF(IKK.EG.1) WRITE(6,1000) M2,N2,IR,IK,CYCX,X0,AL(IE,IR,M,N)

XP=XV+XX
YP=YY

X5=X(IB,IR,J,K)

IK=9

IF(IKK.EG.1) WRITE(6,10CC) J,K,II,IK,DI,XP,YP,X5,Y5

IF(IE.EC.1 • OR. M2.GE.KV) GOTO 16

L=M2+1
L2=N2+1
L2=N2+1
L4=N2+1
IF(I.NE.2) GOTO 7
L1=N2+1
CCNTINUE
IF(IR.EG.1) GOTO 22
L1=M2+1
L2=N2+1
L7=N2+1
L7
3
9
22
                                                Y3=Y(IC,ID,L3,L4)
IZ=0
IX=X0
IK=22
IF(IKK.EC.1) WRITE(6,10C0) M2,N2,IZ,IK,X5,Y5,X1,Y1,X2,Y2
D2=SQRT(AL(IC,ID,M2,N2))+SQRT(AL(IB,IR,J,K))
D2=2.*I1/D2
IF(T.NE.O. .AND. IQ.NE.O) GOTO 2
X4=(L-2)*DX1/2.+TX
Y4=D2*(X4-TX)+TY
14
                                                   GCTO 1
X4=(D1*XP-D2*X5+Y5-YP)/(C1-D2)
Y4=(X4-XP)*D1+YP
2
                                                        4=(X4-XP)=01+7P

k=1

F(IKK.EG.1) WRITE(6,1000) IC,ID,IE,IK,X4,Y4,XF,XX,C2

F(I.EQ.1 -ANC. IR.GT.1) GOTO 20

F(Y4-LE.0.COO1 -AND. Y4-GE.0.) GCTC 8

F(Y4-GT.0. -AND. IZ-EQ.C) GOTO 13

F(T-NE.0.) GOTO 4
1
                                                 IF(T.NE.O.) GOTO 4
Y4=0.
X4=-EM/D2+TX
GCTO 8
CONTINUE
IF(Y4.LE.EM .AND. Y4.GE.(0.999*EM)) GOTO 8
IF(Y4.LT.EM .AND. IZ.EC.O) GOTO 13
IF(T.NE.O. .AND. IQ.NE.O) GOTO 4
Y4=EM
X4=TX+EM/D2
GCTO 8
20
                                                          CTO 8

S=X(IB,IR,J,K)

S=Y(IB,IR,J,K)

6=(EM*(2-IB)-Y5)/02+X5

T=(XP-X4)*EM/(I1*Y4+EM*(IB-1))+X6-XX
                                                IK=4

IF(IKK.EQ.1) WRITE(6,1000) IC.IC, IE, IK, X6, XT, D1, X5, Y5

CALL CONSTI(LO4, IR, IB, ALZ, C3, XT)

CALL BOUND(LO4, IB, IR, [P, I, XT)

AL(IB, IR, M, N) = (ALZ-C3*DYCX*I1) **(2./3.)

D1=I1/SQRT(AL(IB, IR, M, N))

IZ=IZ+1
```

```
XP=XT+XX
IF(IZ.LT.30) GOTO 14
IF(IKK.EG.1) WRITE(6,1000) M,N,IZ,IE,X4,Y4,X5,Y5
                                                   IF([XK.eG.1] WRITE(6,1000) M,N,IZ,IE,X4,Y4,X5,Y5

X4=X6

Y4=EM*(2-IB)
AL(IB,IR,M,N)=AL(IB,IR,J,K)
GCTO 8
CALL FIND(KV,IM,IB,IR,X4,Y4,M1,N1)
IK=13
IF(IKK.eG.1) WRITE(6,1000) M1,N1,IE,IK,X4,Y4
-:x(IC,ID,M1,N1);Y(IC,ID,P1,N1)
IF(IE,EG.1) WRITE(6,1002) M,N,M1,N1,X4,Y4,X(IC,ID,M1,N1)
IF(IC,ID,M1,N1)
IF(M1.eG.M2 .AND. N1.EC.N2 .AND. IE.EG.0) GCTC 8
IF(T.NE.O. .AND. IQ.NE.C) GOTO 15
IF(IE.EC.1) GOTO 16
M2=M1
M2=M1
M2=M1
M2=N1
GCTO 8
X6=(Y2-Y5)/D2+X5
IF(X6.GE.X2) GOTO 10
D3=(Y2-Y1)/(X2-X1)
X5=(Y1-Y5+D2*X5-D3*X1)/(C2-D3)
Y5=D3*(X5-X1)+Y1
IK=15
IF(IKK.EC.1) WRITE(6.1000) M1.N1.IE.IK.Y5.Y5.Y6.D3
13
                                                          X5=(Y1-Y5+D2*X5-D3*X1)/(C2-D3)
Y5=03*(X5-X1)+Y1
IK=15
IF(IKK.EQ.1) WRITE(6,1000) M1,N1,IE,IK,X5,Y5,X6,D3
IF(IC.EC.2) GDTO 23
N2=N2-1
IF(IR.GT.1) GDTO 26
IF(IR.GT.1) GDTO 16
IF(IR.GT.1) GDTO 16
IF(IR.GT.1) GDTO 16
IF(IR.GT.1) GDTO 9
N2=N2-1
IGCTO 9
N2=N2-1
IF(IR.GT.1) GDTO 24
IF(IR.GT.1) GDTO 25
IF(IR.
 26
 23
 10
 25
 24
                                                                Z=MZ-T
GTO 9
(IB, IR, M, N) = X4
(IB, IR, M, N) = Y4
(IB, IR, M, N) = Y4
F(Y4.LE.O.0001 .AND. Y4.GE.O. .AND. IB.
 8
                                                   IF(Y4.LE.EM .AND. Y4.GE.(U.9797-EH)
IK=8
IF(IKK.EQ.1) WRITE(6,1000) M,N,IE,IK,X4,Y4,X1,Y1
GCTO 11
X(IB,IR,M,N)=TX
Y(IB,IR,M,N)=TY
GCTO 12
IE=1
D2=SQRT(AL(IB,IR,J,K))+SQRT(AL(IB,IR,M,N))
D2=4.*I1/(2.*SQRT(AL0)+[2)
IF(T.NE.O. .AND. IQ.NE.0) GCTO 18
X(IB,IR,M,N)=(L-2)*CX1/2.+TX
Y(IB,IR,M,N)=02*(X(IB,IF,M,N)-TX)+TY
GCTO 11
   6
 16
                                                                     (iB, iR, M, N)=(D1*xP-D2*x5+Y5-YP)/(C1-D2)
(iB, iR, M, N)=(X(iB, iR, M, N)-XP)*D1+YF
 18
                                                       CEMPUTATION OF THE ADEITIONAL PEINTS FOR THE UNSTEADY
```

```
C
11
      BCUNDERY STEP OF THE STEADY FIELD TO THE BCCY
     IK=0
IF(IR.GT.1) IK=1
IF(IK.EC.1) WRITE(6,1003)
A=-C1*I1
J1=M
J2=N-1
IF(I.NE.2) GOTO 5
J1=M-1
J2=N
X1=X(IB,IR,J1,J2)-XX
Y1=Y(IB,IR,J1,J2)-YY
X2=X1
     1002
11
10
```

```
GENERAL STEP OF STEADY FIELD
REAL#4 X
COMMON/BC/ T1,T2,T,B,DYCX,D2YDX2,EYCXU,AI,EI,I1
    UNSTEADY BOUNDERY - CENEITIONS ALENE BOCY
COMPUTATION OF UNSTEADY BOUNDERY-FROPERTIES ALONG SHOCK
    IKK=0
IF(IM.EQ.4 .AND. IR.GT.1)
IF(IKK.EC.1) WRITE(6.1001)
FORMAT(1X, RANDS-ENTRY)
C
                     IKK=1
1001
```

```
Ik=0
IF(IKK.EQ.1) WRITE(6,1CC2) M,N,IK,X(IB,IR,M,N),Y(IB,IR,M,N)
FCRMAT(314,5(2X,E11.4))
1002
          C=1

C=1

F(IR.EQ.1) GOTO 1

C=IR-1

F(IB.EQ.1) IC=2

F(IB.EQ.2) IC=1

I2=1./Ei

F(IC.EQ.2) EI2=EI
1
           F(I.NE.2) GOTO 5
5
         TE=O
CALL FIND(KV,IM,IB,IR,X(IB,IR,M,N),Y(IB,IR,M,N),M2,N2)
ALT=AL(IC,ID,M2,N2)
DV=V(IC,ID,M2,N2)
CU=U(IC,ID,M2,N2)
IF(IE.EQ.1) ALT=ALO
IF(IE.EQ.1) DV=O.
IF(IE.EC.1) DU=O.
IF(IE.EC.1) DU=O.
        J2=N
LF=M
IF(I.NE.2) GOTO 8
J1=M
```

```
J2=N-1

LF=N

CONTINUE

IF(AL(IB,IR,J1,J2).LT.AL(IB,IR,M,N) .OR. T.EQ.O.) GCTC 12

IF(IQ.EC.O) GOTO 12

J2=N-1

IF(I.NE.2) GOTO 12

J1=M-1
  8
                                     J2=N-1
IF(I.NE.2) GOTO 12
J1=M-1
AMN=AM*CY
AA1=-2.*(C-1.)*SIN(2.*S)*AMN*ANY2/C
AA2=-4.*(C-1.)*AK*ANY*AMN/C
AA5=(C-1.)*AM*(CX-2.*AMN*ANY2)/C
AA3=4.*(C-1.)*ANY*AMN/C+AA5
AA4=-4.*(C-1.)*ANY*AMN*CCS(S)/(C*SIN(S))
AA5=AK*AA5
IF(J.EC.2) GOTO 7
DO=X(IB,IR,M,N)-P(I,LP)
WRITE(6.4711) M,N,LP,X(IB,IR,M,N),X(IB,IR,J1,J2),P(I,LP)
FORMAT(1X,314,3(2X,E12.5))
DAL=AL(IB,IR,M,N)
IE=O
CALL FIND(KV,IM,IB,IR,X(IB,IR,J1,J2),Y(IE,IR,J1,J2),M3,N3)
DU=U(IC,ID,M3,N3)+U(IC,IC,M2,N2)
DV=V(IC,ID,M3,N3)+V(IC,IC,M2,N2)
IF(IE.EG.1) DU=O
IF(IE.EG.1) DU=O
IF(IE.EG.1) DU=O
IF(IE.EG.1) DV=O
MP=1000C000-*(P(I,LP)-X(IB,IR,J1,J2))
IF(MP.LE.O. AND. T.NE.C.) WRITE(6,1003) M,N,P(I,LP),X(IB,IF)
FCRMAT(1X,*P-X.LE.O. *,214,2(2X,E12.5))
IF(MP.NE.C) GOTO 3
IF(T.NE.O. AND. IQ.NE.C) GOTO 2
DO1=O.
GCTO 2
1000000
                                                                                                                                                                                                                                                                                                   M,N,P(I,LP),X(IB, IR, J1, J2
  1003
                                        D1=0.

GCTO 2

D1=(AL(IB, IR, M, N) - AL(IE, IR, J1, J2))*CC

D1=D1/(4.*DAL*(P(I, LP) - X(IB, IR, J1, J2)))
  3
                                       D1=D1/(4**DAL*(P(I,LP)-X(IB,IR,J1,J2)))
IK=3
IF(IKK.EC.1) WRITE(6,1002) M2,N2,IK,D0,D1,CU,DV,DAL
D2=AI*AK*AM*D0/DAL
D3=-AM*(AK*D0)**2/(4.*CAL)
D4=1./SCRT(AL(IB,IR,P,N))
D5=-D3*D4
D6=AK*AK*AM*PSI(IB,IR,J1,J2)*D0/DAL
D7=Y(IB,IR,P,N)-Y(IB,IR,J1,J2)
A(1,1)=1.+D1+D2+D3
A(1,2)=-(C4+D5)*I1
A(1,3)=C
RI(1)=U(IB,IR,J1,J2)*(1.-D1-D2-D3)-V(IB,IR,J1,J2)*(C4-D5)*I1+C6
A(2,1)=1.
    2
                                       RI(1)=U(IB, IR, J1, J2)*(1.-D1-D2-D3)-V(IB, IR, J1, J2)*(C4-DA(2,1)=1.

A(2,2)=C.

A(2,3)=-2.*AM1/D7-AI*AM2

RI(2)=-(2.*AM1/D7-AI*AM2)*G(IB, IR, J1, J2)-U(IB, IR, J1, J2)

RI(2)=RI(2)+(AM3*DU+AM4*CV)*EI2

A(3,1)=0.

A(3,2)=1.

A(3,3)=-2.*AN1/D7-AI*AN2

RI(3)=-(2.*AN1/D7-AI*AN2)*G(IB, IR, J1, J2)-V(IB, IR, J1, J2)

RI(3)=RI(3)+(AN3*DU+AM4*DV)*EI2
                                   NC=3
CALL SOL VE (NO)
U(IB, IR, M, N)=ES(1)
V(IB, IR, M, N)=ES(2)
G(IB, IR, M, N)=ES(3)
PII=V(IB, IR, M, N)-AI *AN2 *G(IB, IR, M, N)
PII=PII - IAN3 * U(IC, ID, M2, N2) - AN4 * V(IC, ID, M2, N2)) * EI2
PII=AAI * PII / AN1 + AI * AAA 2 * G(IB, IR, M, N) + AA3 * U(IC, ID, M2, N2) * EI2
PII=PII + (AA4 * V(IC, ID, M2, N2) + AA5 * AI * PSI(IC, IC, M2, N2)) * EI2
PII=AI * G(IB, IR, M, N) = AI * (U(IB, IR, M, N) + PII / ((C-1.) * AM)) / AT
IF(AK. EC.O.) AT=1.
PSI(IB, IR, M, N) = PSI(IC, IC, M2, N2) * EI2
GCTO 21
PII=(V(IB, IR, M, N) + V(IB, IR, J1, J2)) * C4 * I1
     20
```

```
GENERAL STEP OF UNSTEACY FIELD
                   IKK=0

IF(IR.EQ.2) IKK=1

IF(IKK.EC.1) WRITE(6,1002)

J1=M

J2=N-1

J3=M-1

J4=N

IF(I.NE.2) GOTO 5
                  IF(I.NE.2) GUTU 5

J1=M-1

J2=N

J3=M

J4=N-1

DC=X(IB,IR,M,N)-X(IB,IR,J1,J2)

D1=O.5*(AL(IB,IR,M,N)-AL(IB,IR,J1,J2))

D2=AI*2.*AK*AM*DO)**2

D4=AL(IB,IR,M,N)+AL(IB,IR,J1,J2)

IF(IKK.EG.1) WRITE(6,10C1) M,N,AL(IE,IR,M,N),D4,CO

D6=(D1+D2+I3)/D4

A(1,1)=1.*D6

A(1,1)=2.*(1.*D3/D4)*I1/D5

RI(1)=2.*AK*AK*AM*PSI(IE,IR,J1,J2)*D0/D4+RI(1)

BC=X(IB,IR,M,N)-X(IB,IR,J3,J4)

B1=(AL(IB,IR,M,N)-AL(IE,IR,J1,J2))*E0/(2.*D0)

B2=2.*AL*AK*AM*BO)

B2=2.*AL*AK*AM*BO)

B3=-O.5*AM*(AK*BO)**2

B4=1./SQRT(AL(IB,IR,M,N))

B5=-O.5*BB3*B4/AL(IB,IR,M,N)

A(2,1)=1.*B6

A(2,2)=-(B4+B5)*I1

RI(2)=U(IB,IR,J3,J4)*(1.*B6)+V(IB,IR,J3,J4)*II*(-B4+B5)

RI(2)=PSI(IB,IR,J3,J4)*AM*BO*AK*AK/AL(IB,IR,M,N)+RI(2)

NG=2
 5
C
```

```
REAL*4 TX, TY, XX, YY, X, Y, AL, DO, D1, D4, D5, ALD, D3, F, XT CCMMON/BA/ ALO, AM, ALD, C1, AK, I, IP, TX, TY, IW, PA, IC, ET, XX, YY CCMMON/BC/ T1, T2, T, B, DYCX, D2YDX2, CYDXU, AI, EI, I1 CCMMON V(2,3,50,20), X(2,3,50,20), F(2,50), U(2,3,50,20), FPSI(2,3,50,20), G(2,3,50,20), AL(2,2,50,20), Y(2,3,50,20), FQ(2,50), DX(2), IN(2,50), IN2(2,50)
COMPUTATION OF UNSTEACY BOUNDERY-FROPERTIES ON BODY
          DETERMINATION OF THE CONSTANT VALUE ALONG THE CHARACTERISTIC
         C
 1000
 3
```

```
CCMPUTATION AND OUTPUT OF THE PRESSURE - CCEFFICIENTS ALONG THE CHOSEN BLADE PU=CPU , PS=CPS
                          PU=CPU , PS=CPS

YN=-EM
YM=AMY*180 ./(4.*ATAN(1.))

IM2=IM-2
DC 7 I=1,2
CT 0C 1 K=1,IR
K3=2
IF(I.EG.1 .AND. K.EQ.1) K3=1
K3=K+K3
DC 3 J1=1,20
LI=J1+1
CALL SWITCH(LI.J1,M.N.I)
IF(K3.GT.IR) GOTO 10
LI=1
LZ=2
CALL SWITCH(LZ,LI,MI,NI,I)
IF(X(1,K,M.N).GT.X(1,K3,MI,NI)) GCTC 6
GCTO 3
D1=REAL(U(1,K,M.N))**2+AIMAG(U(1,K,M.N))**2
IF(D1.EG.1) GOTO 6
IF(J1.EG.1) GOTO 1
J1=J1-1
K1=K1+J1
DO 4 J2=1,J1
LI=J2+1
CALL SWITCH(LI.J2,M.N.I)
K3=J2+K1-J1
PS(I.K3)=-2.*(AL(1,K.M.N)-ALO)/(C1*AM)
PU(I.K3)=-2.*(AL(1,K.M.N)+AI*AK*PSI(1,K.M.N))
PX(I.K3)=X(1,K.M.N)
CCNTINUE
IN2(I,1)=K1
OUTPUT
    10
    36
                            OUTPUT
                            WRITE(7,1015)
IF(L02.EC.0) WRITE(7,996)
IF(L02.EC.1) WRITE(7,995)
WRITE(7,1016) YM,IM2
WRITE(7,1001) AK,AX,C,E,T1
WRITE(7,1011)
DC 8 I=12
IF(I.EQ.2) WRITE(7,1011)
AIF(I.EQ.2) WRITE(7,1012)
IF(I.EQ.2) WRITE(7,1011)
K=IN2(I,1)
DC 9 J=1,K
                                                                          WRITE(7,1001) AK,AX,C,E,T2
WRITE(7,1012)
WRITE(7,1011)
```

```
LI=J+1
CALL SWITCH(LI,J,M,N,I)
WRITE(7,1013) M,N,PX(I,J),PS(I,J),PU(I,J)
HRITE(7,1014)
SHATTE(7,1014)
SHATTE SHATT
    C LINEAR EQUATIONS G*X=C

IKK=0
IF(IKK, NE, 1) GOTO 2

WRITE(6,1001)

MITE(6,1000) (G(M,L),L=1,N),C(M)

MN=N=1
DC 10 M=1,M

C 10 J=K,N

G(M,J)=E(M,J)/G(M,M)

IF(M,EQ.1) GOTO 14

MA=N=1

13 G(M,J)=G(M,J)-G(M,L)*G(L,J)/G(M,M)

14 CCNTINUE

CC 15 C 15 L=1,M

CC 11 = C(1)/G(1,1)

CC 11 = C(1)/G(1,1)

CC 12 = C(1)/G(1,1)

CC 13 L=1,MA

CC 11 = C(1)/G(1,1)

CC 11 = C(1)/G(1,1)

CC 11 = C(1)/G(1,1)

CC 11 = C(N)

NA=N-1

12 C(N)

NA=N-1

13 DC 12 M=1,MA

C(N)=C(N)

NA=N-1

14 C(N)

C(1)=C(1)/G(1,M)*C(M)/G(1,1)

C(1)=C(1)/G(1,M)*C(M)/G(1,M)

C(1)=C(1)/G(1,M)

C(1)=C(1)/G(1,M)

C(1)=C(1)/G(1,M)

C(1)=C(1)/G(
```

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